

THE SEXTIC PERIOD POLYNOMIAL

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In this paper we show that the method of calculating the Gaussian period polynomial which originated with Gauss can be replaced by a more general method based on formulas for Lagrange resolvents. The period polynomial of cyclic sextic fields of arbitrary conductor is determined by way of example.

1. INTRODUCTION

Suppose $p = ef + 1$ is prime. Define the e cyclotomic classes

$$C_j = \{g^{ek+j} \bmod p, \quad j = 0, \dots, e - 1, \quad k = 0, \dots, f - 1\},$$

where g is any primitive root modulo p . The Gaussian periods η_j are defined by

$$(1.1) \quad \eta_j = \sum_{\nu \in C_j} \zeta_p^\nu, \quad \zeta_p = \exp(2\pi i/p).$$

The principal class C_0 contains the e -th power residues and the other classes are its cosets. The η_j are Galois conjugates and the *period polynomial* $\Psi_e(X)$ is their common minimal polynomial over \mathbb{Q} . Gauss introduced the *cyclotomic numbers* (h, k) determined, for a given g , by

$$(h, k) = \#\{\nu \in (\mathbb{Z}/p\mathbb{Z})^* : \nu \in C_h, \nu + 1 \in C_k\}.$$

It follows that

$$(1.2) \quad \eta_0 \eta_h = \sum_{k=0}^{e-1} (h, k) \eta_k + f \delta(h, \ell)$$

where δ is Kronecker's delta and $\ell = 0$ or $e/2$ according as f is even or odd. The coefficients of $\Psi_3(X)$ in terms of p and the coefficients of the quadratic form $4p = A^2 + 27B^2$ were determined by Gauss in *Disquisitiones Arithmeticae*: enough relations exist to determine all (h, k) in terms of p , A , and B . The period polynomial's

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coefficients are then calculated as the symmetric functions of the η s. (See Gauss, Bachmann [1], or Mathews [13].) The same general method and (1.2), with the appropriate quadratic form, was used to solve the cases $e = 4$ (Gauss [5], 1825), $e = 5$ (Emma Lehmer [10], 1951), $e = 6$ (D. H. and Emma Lehmer [9], 1984), and $e = 8$ (Evans [4], 1983).

Throughout ζ_n is the root of unity $\exp(2\pi i/n)$ and $\mu(\cdot)$ is the Möbius function. Let χ be a primitive character of degree e and modulus p . Recall the Lagrange resolvent (sometimes called a Gauss sum) defined by $\tau(\chi) = \sum_{j=0}^{p-1} \chi(j)\zeta_p^j$. Provided that the periods are defined with primitive root g such that $\chi(g) = \zeta_e$,

$$(1.3) \quad \tau(\chi^j) = \sum_{k=0}^{e-1} \chi^j(k)\eta_k.$$

The inverse of (1.3) is

$$(1.4) \quad \eta_j = e^{-1} \sum_{h=0}^{e-1} \zeta_p^{-hj} \tau(\chi^h).$$

Since (1.4) does not depend on the existence of a primitive root, it defines η_j for any character χ of arbitrary modulus.

Field-theoretically, embed an abelian field K of degree $[K : \mathbb{Q}] = e$, Galois group $\mathcal{G} = \text{Gal}(K/\mathbb{Q})$, and conductor m in $\mathbb{Q}[\zeta_m]$. Let $\widehat{\mathcal{G}}$ be the group of Dirichlet characters modulo m which annihilate $\text{Gal}(\mathbb{Q}[\zeta_m]/K) \subset \text{Gal}(\mathbb{Q}[\zeta_m]/\mathbb{Q}) \cong (\mathbb{Z}/m\mathbb{Z})^*$. We say that K belongs to $\widehat{\mathcal{G}}$ and $\widehat{\mathcal{G}}$ is associated to K . Then $\widehat{\mathcal{G}}$ is the dual of \mathcal{G} and $\widehat{\mathcal{G}} \cong \mathcal{G}$ [15, Chapter 3]. The Gaussian period is defined in this most general case by

$$(1.5) \quad \eta_j = e^{-1} \sum_{\chi \in \widehat{\mathcal{G}}} \zeta_p^{-hj} \tau(\chi).$$

This reduces to (1.4) when $\widehat{\mathcal{G}}$ is cyclic and to (1.1) when m is prime. It is easy to see that $\eta_0 = \text{Tr}_K^{\mathbb{Q}[\zeta_m]} \zeta_m$. The class \mathcal{C}_0 becomes the kernel of $\widehat{\mathcal{G}}$ in $(\mathbb{Z}/m\mathbb{Z})^*$. For all $j \in \mathcal{C}_0$, the map $\zeta_m \mapsto \zeta_m^j$ is an automorphism of $\mathbb{Q}[\zeta_m]$ which fixes K . The period polynomial in this general case was determined in an ad hoc way for cyclic cubic fields by Châtelet [3], and for cyclic quartic fields by Nakahara [14] and (independently) by the author [8]. The computation of period polynomials can be made systematic through well-known formulas for arithmetic of Lagrange resolvents.

LEMMA 1.1. (A) For χ_m and χ_n of conductors m, n respectively with $\text{gcd}(m, n) = 1$,

$$\tau(\chi_m \chi_n) = \chi_n(m) \chi_m(n) \tau(\chi_m) \tau(\chi_n).$$

(B) If the conductor of χ is m and c is odd, then

$$\sum_{j=0}^{cm-1} \chi(j)\zeta_{cm} = \mu(c)\chi(c)\tau(\chi).$$

If c is even the sum vanishes.

(C) If \mathcal{X} is the set of characters of prime conductor p , $\ell \mid p - 1$, and $\chi^\ell \neq 1$, then

$$\prod_{\substack{\psi \in \mathcal{X} \\ \psi^\ell = 1}} \tau(\chi\psi) = \bar{\chi}^\ell(\ell)\tau(\chi^\ell) \prod_{\substack{\psi \in \mathcal{X} \\ \psi^\ell = 1}} \tau(\psi).$$

PROOF: The first two formulas are routine; the last is the theorem of Hasse and Davenport [6, 20.4.IX]. □

REMARKS. (1) The Hasse-Davenport Theorem has been used extensively in cyclotomy of prime modulus and composite *degree*; see, for example, Buck, Smith, Spearman and Williams [2].

(2) The Gaussian periods appear as summands of the ‘Basiszahl’ of an abelian number field in the elegant paper of Lettl [11].

2. SEXTIC DIRICHLET CHARACTERS

NOTATION. From now on ψ , ξ , and χ will be primitive quadratic, cubic, and sextic characters, respectively. A subscript, for example χ_m , will indicate the conductor. Powers such as χ_m^2 denote a possibly imprimitive character.

We may associate to every cubic character ξ_m an integer m in $\mathbb{Q}[\zeta_3]$ as follows. There is a prime-power decomposition of

$$\xi_m = \xi_{p_1} \cdots \xi_{p_\nu}$$

where each p_j is either 9 or a prime congruent to 1 modulo 3. We may assume that if $3 \mid m$, the divisors are ordered such that $p_1 = 9$. ξ_{p_j} is a complex cubic residue symbol modulo p_j , where p_j is a prime of $\mathbb{Q}[\zeta_3]$ lying over p_j (over 3, if $p_1 = 9$). Set

$$(2.1) \quad m = \begin{cases} \prod_{j=1}^{\nu} p_j & 3 \nmid m \\ 3p_1 \prod_{j=2}^{\nu} p_j & \text{otherwise.} \end{cases}$$

Since there are two conjugate primes lying over each p_j , it is clear that there are 2^ν different cubic characters, where m is divisible by ν distinct primes, but ξ and $\bar{\xi}$ generate the same group and are associated to the same field. Therefore there are $2^{\nu-1}$ distinct cubic fields of conductor m .

LEMMA 2.1. *If ξ_m is a cubic character and ψ_m is a quadratic character, both primitive with conductor m , not necessarily prime, then*

$$\tau(\xi_m \psi_m) = \xi_m(2) \tau(\psi_m) \frac{\tau(\bar{\xi}_m)^2}{m}.$$

PROOF: The conductor of a primitive quadratic (respectively cubic) character is square-free except for powers of 2 (respectively 3), so m is square-free. If m is prime, the lemma is just the Hasse-Davenport Theorem (Lemma 1.1C). Proceeding by induction on the number of primes dividing m , write $m = pq$, p prime, and factor $\xi_m = \xi_p \xi_q$ (similarly ψ_m).

$$\begin{aligned} \tau(\xi_m \psi_m) &= \xi_q(p) \psi_q(p) \xi_p(q) \psi_p(q) \tau(\xi_p \psi_p) \tau(\xi_q \psi_q), \quad \text{by Lemma 1.1A and } \gcd(p, q) = 1 \\ &= \frac{\xi_p(2)}{p} \frac{\xi_q(2)}{q} \xi_q(p) \xi_p(q) \tau(\bar{\xi}_p)^2 \tau(\bar{\xi}_q)^2 \psi_q(p) \psi_p(q) \tau(\psi_p) \tau(\psi_q), \\ &\hspace{15em} \text{by inductive hypothesis} \\ &= \xi_m(2) \tau(\psi_m) \frac{\tau(\bar{\xi}_m)^2}{m}. \end{aligned}$$

□

LEMMA 2.2. *The complex integer m defined by (2.1) is equal to the Jacobi sum $J(\xi_m, \xi_m) = \tau(\xi_m)^2 / \tau(\bar{\xi}_m)$. The sextic resolvent $\tau(\xi_m \psi_m) = \xi_m(2) \tau(\psi_m) \tau(\xi_m) \bar{m} / m$.*

PROOF: The first clause is well-known, for example, Hasse [7, Section 2(1)]. The second now follows using the previous lemma. □

3. THE PERIOD POLYNOMIAL

The sextic period polynomial will be determined from these lemmas. Let K be a cyclic sextic field of conductor m . K is the compositum of a quadratic field K_2 of conductor m_2 and a cyclic (hence real) cubic field K_3 of conductor m_3 . Set $g = \gcd(m_2, m_3)$. In this section we shall assume that $3 \nmid g$; we shall treat the other case afterwards. Let $n_i = m_i/g$, $i = 2, 3$. Then $m = gn_2n_3$ and $3 \nmid g$ implies g , n_2 , and n_3 are pairwise co-prime. The sextic character χ_m can be factored into a product of cubic and quadratic characters:

$$\chi_m = \xi_{m_3} \psi_{m_2} = \xi_g \xi_{n_3} \psi_g \psi_{n_2}.$$

We define g and m as the Jacobi sums of ξ_g and ξ_{m_3} . (If $g = 1$, we define $g = 1$.) Let $n = m/g$, which is $J(\xi_{n_3}, \xi_{n_3})$ if $n_3 > 1$.

We are now equipped to determine the resolvents necessary to use (1.4). Clearly $\tau(\chi^6) = \mu(m)$. Since K_3 is real, $\psi_{m_2}(-1) = \chi_m(-1)$. From Lemma 1.1B and the result of Gauss on $\tau(\psi)$,

$$\tau(\chi_m^3) = \mu(n_3)\psi_{m_2}(n_3)\sqrt{m_2^*}, \quad m_2^* = \chi_m(-1)m_2.$$

From Lemmas 1.1B and 2.2,

$$\tau(\chi_m^2) = \mu(n_2)\xi_{m_3}(n_2)\tau(\bar{\xi}_{m_3}).$$

The sextic resolvent can be found in terms of the quadratic and cubic ones by Lemmas 1.1 and 2.2.

(3.1)

$$\begin{aligned} \tau(\chi_m) &= \xi_{n_3}\psi_{n_2}(g)\xi_g\psi_g(n_3n_2)\tau(\xi_{n_3}\psi_{n_2})\tau(\xi_g\psi_g) \\ &= \xi_{n_3}\psi_{n_2}(g)\xi_g\psi_g(n_3n_2)\xi_{n_3}(n_2)\psi_{n_2}(n_3)\tau(\xi_{n_3})\tau(\psi_{n_2})\xi_g(2)\tau(\psi_g)\tau(\xi_g)\bar{g}/g \\ &= \xi_{m_3}(n_2)\psi_{m_2}(n_3)\xi_g(2)\tau(\psi_{m_2})\tau(\xi_{m_3})\bar{g}/g \end{aligned}$$

It is easy to see that $\tau(\chi_m^4) = \overline{\tau(\chi_m^2)}$ and $\tau(\chi_m^5) = \chi_m(-1)\overline{\tau(\chi_m)}$.

The symbolic coefficients are simpler if we work with the *reduced Gaussian period* $\lambda_j = e\eta_j - \mu(m)$. The reduced period polynomial $\Lambda(X)$ is given by

$$\Lambda(X) = \text{Irr}_{\mathbb{Q}}^K \lambda_0 = \text{Irr}_{K_2}(\lambda_0)\text{Irr}_{K_2}(\lambda_1).$$

The normal and reduced period polynomials of degree e are related by $\Psi_e(X) = e^{-e}\Lambda_e(eX - \mu(m))$.

PROPOSITION 3.1. *The minimal polynomial $\text{Irr}_{K_2}(\lambda_0)$ of λ_0 over K_2 is*

$$(X - \lambda_0)(X - \lambda_2)(X - \lambda_4) = X^3 + c_2X^2 + c_1X + c_0$$

where a calculation shows

$$\begin{aligned} c_2 &= -3\mu(n_3)\psi_{m_2}(n_3)\sqrt{m_2^*} \\ c_1 &= -3\mu(n_2)^2m_3 - 3\chi_m(-1)(m - \mu(n_3)^2m_2) \\ &\quad - 6\mu(n_2)\psi_{m_2}(n_3)n_3\text{Re}(\xi_g(2)\bar{g})\sqrt{m_2^*} \\ c_0 &= -2\mu(n_2)m_3\text{Re}(m) - 6\mu(n_2)m\chi_m(-1)\text{Re}(\xi_g(2)[\bar{n}g - \mu(n_3)\bar{g}]) \\ &\quad \{-3\mu(n_2)^2m_3(2\text{Re}(\xi_g(2)n) - \mu(n_3)) \\ &\quad - n_2\chi_m(-1)(2n_3\text{Re}(\bar{n}g^2) + \mu(n_3)(g - 3m_3))\}\psi_{m_2}(n_3)\sqrt{m_2^*}. \end{aligned}$$

PROOF: The coefficients were computed from (1.4) and the values of $\tau(\chi^j)$ using the Maple symbolic algebra system, and then simplified by Lemma 2.2. □

To rewrite Proposition 3.1 with rational integer coefficients, create new variables with the assignments

$$(3.2) \quad \xi_g(2) = \frac{z_0 + \sqrt{-3}z_1}{2}, \quad g = \frac{A + 3\sqrt{-3}B}{2}, \quad m = \frac{L + 3\sqrt{-3}M}{2}, \quad n = \frac{R + 3\sqrt{-3}S}{2}$$

so that

$$(3.3) \quad R = \frac{AL + 27BM}{2g}, \quad S = \frac{AM - BL}{2g}$$

Maple’s Gröbner basis reduction package normalised c_0 and c_1 with respect to (3.3) and the obvious relations on the conductors. We have that

$$\begin{aligned} c_1 &= -3\mu(n_2)^2 m_3 - 3\chi_m(-1) \left(m - \mu(n_3)^2 m_2 \right) \\ &\quad - \frac{3}{2} \mu(n_2) \psi_{m_2}(n_3) n_3 (z_0 A + 9z_1 B) \sqrt{m_2^*} \\ c_0 &= -\mu(n_2) m_3 L \\ &\quad - \frac{3}{4} \mu(n_2) m \chi_m(-1) [z_0 (AR + 27BS) + 9z_1 (AS - BR) - 2\mu(n_3) (z_0 A + 9z_1 B)] \\ &\quad - \left\{ \frac{3}{2} \mu(n_2)^2 m_3 (z_0 R - 9z_1 S - 2\mu(n_3)) + \frac{1}{2} n_2 \chi_m(-1) [n_3 (2gR \right. \\ &\quad \left. + 27ABS - 27B^2R) + 2\mu(n_3) (g - 3m_3)] \right\} \psi_{m_2}(n_3) \sqrt{m_2^*}. \end{aligned}$$

Define rational numbers c'_0 and c''_0 by $c_0 = c'_0 + c''_0 \sqrt{m_2^*}$; similarly c'_2, c'_1 and c''_1 . The conjugate polynomial $\text{Irr}_{K_2}(\lambda_1)$ is

$$X^3 - c_2 X^2 + (c'_1 - c''_1 \sqrt{m_2^*})X + (c'_0 - c''_0 \sqrt{m_2^*}).$$

Writing the reduced period polynomial

$$\Lambda(X) = X^6 + 0X^5 + d_4 X^4 + d_3 X^3 + d_2 X^2 + d_1 X + d_0$$

we have our main result.

THEOREM 1. *The coefficients d_ν are given by*

$$\begin{aligned} d_4 &= -6\mu(n_2)^2 m_3 - 3 \left(2m + \mu(n_3)^2 m_2 \right) \chi_m(-1) \\ d_3 &= \mu(n_2) \left\{ -2m_3 L \right. \\ &\quad \left. + \left(-6\mu(n_3) \left(z_0 A + 9z_1 B \right) - \frac{3}{2} [z_0 (AR + 27BS) + 9z_1 (AS - BR)] \right) \chi_m(-1) m \right\} \\ d_2 &= c_1'^2 - (c_1''^2 + 2c_0'' c_2'') \chi_m(-1) m_2 \\ d_1 &= 2(c_0' c_1' - c_0'' c_1'') \chi_m(-1) m_2 \\ d_0 &= c_0'^2 - c_0''^2 \chi_m(-1) m_2. \end{aligned}$$

The expanded expressions for the three trailing coefficients are very large.

- REMARKS. (1) d_4 depends on m_3 but not on m , that is, the choice of K_3 is irrelevant.
 (2) The coefficients are all divisible by g . This is to be expected, since λ_j is a sum of terms divisible by a prime over g .
 (3) When $4 \mid m_2$ the polynomial is even. (This can also be seen directly from the definition.)
 (4) The factor $\psi_{m_3}(n_3)$ does not appear explicitly.

COROLLARY 3.2. *If $m = m_2 = m_3 = g$, then*

$$\begin{aligned}
 c_2 &= -3 \sqrt{\chi_m(-1)m} \\
 c_1 &= -3m - \frac{9}{2} (z_0L + 3z_1M) \sqrt{\chi_m(-1)m} \\
 c_0 &= -mL + 27 \chi_m(-1) z_1 Mm + ([3 - 3z_0 + 4\chi_m(-1)]m - \chi_m(-1)L^2) \sqrt{\chi_m(-1)m} \\
 d_4 &= -3m(3\chi_m(-1) + 2) \\
 d_3 &= [-9(z_0L + 3z_1M)\chi_m(-1) - 2L]m \\
 d_2 &= [9(2 - 2z_0 - 3z_1^2)\chi_m(-1) + 33]m^2 \\
 &\quad - \frac{3}{4} [3(z_0^2L^2 + 18z_0z_1LM - 3z_1^2L^2)\chi_m(-1) + 8L^2]m \\
 d_1 &= -9(9z_1M + z_0^2L + 9z_0z_1M - z_0L)\chi_m(-1)m^2 \\
 &\quad + 3(2Lm + 4z_0Lm + 36z_1Mm - 9z_1L^2M - z_0L^3)m \\
 d_0 &= (27\chi_m(-1)z_1M - L)^2m^2 - ([3 - 3z_0 + 4\chi_m(-1)]m - \chi_m(-1)L^2)^2\chi_m(-1)m.
 \end{aligned}$$

REMARK. If moreover m is prime we recover formulas (8)–(12) of [9].

EXAMPLE. To illustrate the theorem, we show in Table 1 all ten sextic period polynomials of conductor 91. The two sextic subfields of $\mathbb{Q}[\zeta_{91}]$ of smaller conductor are omitted. Although the reduced polynomial is simpler symbolically, the period polynomial in the table has smaller coefficients.

TABLE 1. Sextic period polynomials of conductor 91

m_2	m_3	g	m	$\xi_g(2)$	$\Psi_6(X)$
7	13	1	$\frac{-5+3\sqrt{-3}}{2}$	1	$X^3 - X^5 + 14X^4 + 13X^3 + 58X^2 + 16X + 8$
{1, 8, 18, 25, 44, 51, 53, 57, 60, 64, 79, 86}					
7	91	$\frac{1+3\sqrt{-3}}{2}$	$-8 - 3\sqrt{-3}$	ζ_3	$X^6 - X^5 + X^4 + 13X^3 + 162X^2 - 400X + 736$
{1, 2, 4, 8, 16, 23, 32, 37, 46, 57, 64, 74}					
7	91	$\frac{1+3\sqrt{-3}}{2}$	$\frac{11+9\sqrt{-3}}{2}$	$\bar{\zeta}_3$	$X^6 - X^5 + X^4 + 13X^3 + 71X^2 + 419X + 827$
{1, 8, 9, 11, 30, 57, 58, 64, 67, 72, 81, 88}					
13	7	1	$\frac{1+3\sqrt{-3}}{2}$	1	$X^6 - X^5 - 17X^4 + 4X^3 + 57X^2 - 18X - 27$
{1, 22, 27, 29, 36, 43, 48, 55, 62, 64, 69, 90}					
13	91	$\frac{-5+3\sqrt{-3}}{2}$	$-8 - 3\sqrt{-3}$	$\bar{\zeta}_3$	$X^6 - X^5 - 31X^4 + 4X^3 + 253X^2 + 101X - 391$
{1, 4, 16, 17, 23, 27, 64, 68, 74, 75, 87, 90}					
13	91	$\frac{-5+3\sqrt{-3}}{2}$	$\frac{11+9\sqrt{-3}}{2}$	$\bar{\zeta}_3$	$X^6 - X^5 - 31X^4 + 4x^3 + 162X^2 - 81X - 27$
{1, 3, 9, 10, 27, 30, 61, 64, 81, 82, 88, 90}					
91	7	$\frac{1+3\sqrt{-3}}{2}$	$\frac{1+3\sqrt{-3}}{2}$	ζ_3	$X^6 - X^5 + 22X^4 - 22X^3 + 148X^2 - 148X + 337$
{1, 6, 20, 22, 29, 34, 36, 41, 43, 64, 76, 83}					
91	13	$\frac{-5+3\sqrt{-3}}{2}$	$\frac{-5+3\sqrt{-3}}{2}$	$\bar{\zeta}_3$	$X^6 - X^5 + 21X^4 - 22X^3 + 58X^2 + 23X + 155$
{1, 5, 25, 31, 34, 47, 51, 53, 64, 73, 79, 83}					
91	91	$-8 - 3\sqrt{-3}$	$-8 - 3\sqrt{-3}$	1	$X^6 - X^5 + 8X^4 - 113X^3 + 435X^2 - 666X + 428$
{1, 4, 16, 23, 34, 45, 54, 59, 64, 74, 83, 89}					
91	91	$\frac{11+9\sqrt{-3}}{2}$	$\frac{11+9\sqrt{-3}}{2}$	ζ_3	$X^6 - X^5 + 8X^4 - 22X^3 - 20X^2 + 426X + 1611$
{1, 9, 19, 24, 30, 33, 34, 64, 80, 81, 83, 88}					

REMARK. For units in sextic fields see Mäki [12].

4. THE SPECIAL CASE $3 \mid \gcd(m_2, m_3)$

In this case, we write $g = \gcd(m_2, m_3)/3$, $n_3 = m_3/9g$, $n_2 = m_2/3g$, so that 9,

$g, n_2,$ and n_3 are pairwise co-prime. We have

$$\chi_m = \xi_9 \xi_g \xi_{n_3} \psi_3 \psi_g \psi_{n_2}.$$

We can relate the Gauss sums for χ_m to the character $\chi'_m = \psi_3 \chi_m$.

LEMMA 4.1. *If ξ_9 is the cubic residue symbol $(\frac{\cdot}{9})$, then $\tau(\xi_9 \psi_3) = \tau(\xi_9)$.*

Since the character belonging to a field is determined only up to complex conjugation, we may assume the condition of the lemma without loss of generality.

LEMMA 4.2. *With χ_9 normalised, $\tau(\chi_m^6) = \tau(\chi'_m{}^6) = \tau(\chi_m^3) = \tau(\chi'_m{}^3) = 0$, $\tau(\chi_m^2) = \tau(\chi'_m{}^2)$, $\tau(\chi_m^4) = \tau(\chi'_m{}^4) = \overline{\tau(\chi_m^2)}$, $\tau(\chi_m) = \psi_3(m/9)\tau(\chi'_m)$ and $\tau(\chi_m^5) = -\psi_3(m/9)\tau(\chi'_m{}^5)$.*

PROOF: We have $\tau(\chi_m^6) = \tau(\chi'_m{}^6) = \mu(9(m/9)) = 0$. For the cube, use Lemma 1.1B. From the definition $\chi_m^2 = \chi'_m{}^2$. Since exactly one of χ_m, χ'_m is an even character, we need prove only the relation on $\tau(\chi_m)$. By Lemma 1.1,

$$\begin{aligned} \tau(\chi_m) &= \chi_9(gn_2n_3)\chi_g\psi_{n_2}\xi_{n_3}(9)\tau(\chi_9)\tau(\chi_9\psi_{n_2}\xi_{n_3}) \\ &= \chi_9(gn_2n_3)\chi_g\xi_{n_3}(9)\tau(\xi_9)\xi_{gn_3}(n_2)\psi_{gn_2}(n_3)\xi_g(2)\tau(\psi_{gn_2})\tau(\xi_{gn_3})g/g \end{aligned}$$

using Lemmas 2.2, 4.1, and $\psi_g(9) = 1$. Expanding $\chi_9(gn_2n_3)$ and recombining terms,

$$\tau(\chi_m) = \xi_{m_3}(n_2)\psi_3(gn_2n_3)\psi_{gn_2}(n_3)\xi_g(2)\tau(\xi_{m_3})\tau(\psi_{gn_2})g/g$$

Since gn_2 is the conductor of the quadratic field associated to χ' , comparing this expression to (3.1) gives the lemma. □

The coefficients can now be computed.

PROPOSITION 4.3. *The polynomial $\text{Irr}_{K_2}(\lambda_0)$ is given by $X^3 + 0X^2 + c_1X + c_0$ where*

$$\begin{aligned} c_1 &= -3 \left(\mu(n_2)^2 m_3 + \chi_m(-1)m \right) \\ &\quad - 6\sqrt{-1}\mu(n_2)\psi_{m_2}(n_3)\psi_3(m/9)n_3 \text{Im}(\xi_g(2)\bar{g})\sqrt{\chi'_m(-1)gn_2} \\ c_0 &= -2\mu(n_2)m_3 \text{Re}(m) + 6\mu(n_2)\chi_m(-1)m \text{Re}(\xi_g(2)\bar{ng}) \\ &\quad - 6\sqrt{-1}\mu(n_2)^2\psi_{m_2}(n_3)\psi_3(m/9)m_3 \text{Im}(\xi_g(2)n)\sqrt{\chi'_m(-1)gn_2} \\ &\quad - 18\sqrt{-1}\psi_{m_2}(n_3)\psi_3(m/9)\chi_m(-1)n_2n_3 \text{Im}(\bar{ng}^2)\sqrt{\chi'_m(-1)gn_2}. \end{aligned}$$

The radicand $\chi'_m(-1)gn_2 = -m_2/3$ appears because $\sqrt{3}$ is hidden in the imaginary parts. We have $\sqrt{m_2^*} = -\chi_m(-1)\sqrt{-3}\sqrt{\chi'_m(-1)gn_2}$. Making the assignments (3.2) and (3.3) we obtain:

THEOREM 2. For $3 \mid \gcd(m_2, m_3)$ the reduced period polynomial $\Lambda(X) = X^6 + \sum_{\nu=0}^4 d_\nu X^\nu$ is given by

$$\begin{aligned}
 c_1 &= -3 \left[\left(\mu(n_2)^2 m_3 + \chi_m(-1)m \right) \right. \\
 &\quad \left. - \frac{3}{2} \mu(n_2) \psi_{m_2}(n_3) \psi_3(m/9) \chi_m(-1) n_3 (3 z_0 B - z_1 A) \sqrt{m_2^*} \right] \\
 c_0 &= -\mu(n_2) m_3 L + \frac{3}{4} \mu(n_2) \chi_m(-1) m [z_0 (AR + 27 BS) + 9 z_1 (AS - BR)] \\
 &\quad + \frac{27}{2} \psi_{m_2}(n_3) \psi_3(m/9) (2 gS + ABR - A^2 S) \sqrt{m_2^*} \\
 &\quad + \frac{3}{2} \mu(n_2)^2 \psi_{m_2}(n_3) \psi_3(m/9) \chi_m(-1) m_3 (3 z_0 S + z_1 R) \sqrt{m_2^*} \\
 d_4 &= -6 \left[\left(\mu(n_2)^2 m_3 + \chi_m(-1)m \right) \right] \\
 d_3 &= -2\mu(n_2) m_3 L + \frac{3}{2} \mu(n_2) \chi_m(-1) m [z_0 (AR + 27 BS) + 9 z_1 (AS - BR)] \\
 d_2 &= c_1'^2 - c_1''^2 \chi_m(-1) m_2 \\
 d_1 &= 2 (c_0' c_1' - c_0'' c_1'' \chi_m(-1) m_2) \\
 d_0 &= c_0'^2 - c_0''^2 \chi_m(-1) m_2
 \end{aligned}$$

where $c_j = c_j' + c_j'' \sqrt{m_2^*}$.

As in the general case we suppress writing the trailing terms in full.

REMARK. The remarks after Theorem 1 hold. In addition, if $n_2 = 1$ and $m_2 \equiv 3 \pmod{4}$, d_4 vanishes and d_2 is independent of K_3 . d_2 is also independent of K_3 whenever g is prime or trivial, since we can fix g .

EXAMPLE. In Table 2 we give the period polynomials of the six sextic fields of conductor 63. When $m_2 = 7$, Theorem 1 holds; for $m_2 = 21$, we use Theorem 2.

TABLE 1. Sextic period polynomials of conductor 63

m_2	m_3	g C_0	m	$\xi_g(2)$ $\Psi_6(X)$
7	9	1	$\frac{-3+3\sqrt{-3}}{2}$	1 $X^6 + 9X^4 + 5X^3 + 36X^2 + 12X + 8$
7	63	$\frac{1+3\sqrt{-3}}{2}$	$\frac{-15-3\sqrt{-3}}{2}$	ζ_3 $X^6 + 14X^3 + 63X^2 - 168X + 161$
7	63	$\frac{1-3\sqrt{-3}}{2}$	$6 - 3\sqrt{-3}$	$\bar{\zeta}_3$ $X^6 + 14X^3 + 63X^2 + 210X + 224$
21	9	1	$\frac{-3+3\sqrt{-3}}{2}$	1 $X^6 - 12X^4 + 5X^3 + 36X^2 - 30X + 1$
21	63	$\frac{1+3\sqrt{-3}}{2}$	$\frac{-15-3\sqrt{-3}}{2}$	ζ_3 $X^6 - 21X^4 + 14X^3 + 63X^2 - 21X - 35$
21	63	$\frac{1+3\sqrt{-3}}{2}$	$6 - 3\sqrt{-3}$	ζ_3 $X^6 - 21X^4 + 14X^3 + 63X^2 - 84X + 28$

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