

# Regular solutions to the dissipative Aw–Rascle system

# Nilasis Chaudhuri

Department of Mathematics, Informatics and Mechanics, University of Warsaw, ul. Banacha 2, Warszawa, Poland (nchaudhuri@mimuw.edu.pl)

# Tomasz Piasecki

Department of Mathematics, Informatics and Mechanics, University of Warsaw, ul. Banacha 2, Warszawa, Poland (tpiasecki@mimuw.edu.pl)

# Ewelina Zatorska

Mathematics Institute, University of Warwick, Zeeman Building, Coventry CV4 7AL, United Kingdom (ewelina.zatorska@warwick.ac.uk) (corresponding author)

(Received 15 November 2024; revised 9 May 2025; accepted 9 May 2025)

In this article, we prove the local-in-time existence of regular solutions to dissipative Aw–Rascle system with the offset equal to gradient of some increasing and regular function of density. It is a mixed degenerate parabolic-hyperbolic hydrodynamic model, and we extend the techniques previously developed for compressible Navier–Stokes equations to show the well-posedness of the system in the  $L_2 - L_2$  setting. We also discuss relevant existence results for offset involving singular or non-local functions of density.

Keywords: dissipative Aw–Rascle system; regular solutions; hydrodynamic model; local well-posedness; degenerate parabolic-hyperbolic system

2020 Mathematics Subject Classification: 35M10; 35q55

# 1. Introduction

We investigate the, so-called, dissipative Aw-Rascle system

$$\int \varrho_t + \operatorname{div}\left(\varrho\mathbf{u}\right) = 0 \tag{1.1a}$$

$$(\boldsymbol{\varrho}\mathbf{w})_t + \operatorname{div}\left(\boldsymbol{\varrho}\mathbf{w}\otimes\mathbf{u}\right) = 0 \tag{1.1b}$$

on  $\mathbb{T}^3 \times (0, T)$ , where  $\mathbb{T}^3$  is a three-dimensional torus. The unknown of the system are the density  $\rho(t, x)$  and the desired velocity of motion  $\mathbf{w}(t, x)$ . The actual velocity of motion  $\mathbf{u}$  is given by the relation:

© The Author(s), 2025. Published by Cambridge University Press on behalf of The Royal Society of Edinburgh. This is an Open Access article, distributed under the terms of the Creative Commons Attribution licence (http://creativecommons.org/licenses/by/4.0), which permits unrestricted reuse, distribution and reproduction, provided the original article is properly cited.

$$\mathbf{u} = \mathbf{w} - \nabla p(\varrho), \tag{1.2}$$

where  $\nabla p(\varrho)$  is the velocity offset, with a given offset function  $p(\cdot) \in C^5(\mathbb{R}_+)$ . System (1.1) is supplemented with the initial data

 $\varrho(0, x) = \varrho_0(x), \quad \mathbf{u}(0, x) = \mathbf{u}_0(x).$ (1.3)

The purpose of the article is to prove local-in-time existence of regular solutions to system (1.1) under the following assumptions on the data

$$(\varrho_0, \mathbf{u}_0) \in H^4(\mathbb{T}^3) \times H^3(\mathbb{T}^3).$$
(1.4)

We may also assume less regularity of  $\rho_0$  at the price of additional assumption on well-prepared data, more precisely

$$\varrho_0 \in H^3(\mathbb{T}^3), \ \mathbf{u}_0 + \nabla p(\varrho_0) \in H^3(\mathbb{T}^3).$$

$$(1.5)$$

Our goal is to prove local in time existence of regular solutions to (1.1) Note that, using (1.2), we can rewrite (1.1) as

$$\begin{cases} \varrho_t + \operatorname{div}(\varrho \mathbf{w}) - \operatorname{div}(\varrho \nabla p(\varrho)) = 0, \\ (\varrho \mathbf{w})_t + \operatorname{div}(\varrho \mathbf{w} \otimes \mathbf{w}) = \operatorname{div}(\varrho \mathbf{w} \otimes \nabla p(\varrho)), \end{cases}$$
(1.6*a*) (1.6*b*)

which is the equivalent formulation as long as the solution remains sufficiently regular. Moreover, assuming  $\rho > 0$  we can further transform this system to obtain

$$\begin{cases} \varrho_t + \operatorname{div}\left(\varrho\mathbf{w}\right) = \operatorname{div}\left(\varrho p'(\varrho)\nabla\varrho\right), & (1.7a)\\ \mathbf{w}_t + \mathbf{u}\cdot\nabla\mathbf{w} = 0, & (1.7b) \end{cases}$$

subject to the initial data

$$(\varrho, \mathbf{w})|_{t=0} = (\varrho_0, \mathbf{u}_0 + \nabla p(\varrho_0)). \tag{1.8}$$

System (1.1) with closure relation (1.2) was recently considered by Acaves et al. [1] in the context of pedestrian flow. Their offset function was actually singular with respect to (w.r.t.) density

$$p(\varrho) = \epsilon \left(\frac{1}{\varrho} - \frac{1}{\varrho_{\max}}\right)^{-\beta}.$$
(1.9)

This offset function acts as a barrier to ensure that the density remains below its maximum  $\rho_{\text{max}}$ , which models the formation of congestion within the crowd. For an up-to-date overview of the literature on the macroscopic models of crowds, we refer to the recent overview articles [5–7].

The dissipative Aw–Rascle system is a model inspired by the one-dimensional Aw–Rascle road traffic model. For derivation of this model and its qualitative analysis, we refer to [3, 4]. The offset function (1.9) was actually originally proposed for

that model in the work of Berthelin et al. [8], as a remedy to the lack of uniform bound for the density.

The classical Aw-Rascle model for traffic differs from the dissipative Aw-Rascle system (1.1), not only in its spatial dimension but also because it uses a scalar offset, i.e.,  $u = w - p(\rho)$ . Incorporating the offset in the form of gradient (1.2) resolves the dimension discrepancy in the closure relation between the velocities **u** and **w**. However, the whole system changes its character from hyperbolic to mixed hyperbolic-parabolic type due to additional dissipation effect in the continuity equation (1.7a). While for description of the multi-lane traffic, first order systems seem to be a more suitable [2, 20, 33], and it was demonstrated in [1] that system (1.1) and (1.2) correctly capture the fundamental diagram for the pedestrian flow.

The mathematical properties of the dissipative Aw-Rascle system have been explored for the first time in [10]. The authors demonstrated the existence and weak-strong uniqueness of Young-measure solutions to the system (1.1) and (1.2) with  $p(\varrho) = \varrho^{\gamma}, \gamma > 0$ . Their result states that the measure-valued solution coincides with the strong solution emanating from the same initial data, as long as the latter exists. However, the existence of regular solutions was assumed rather than proven, which motivates the current article. Our aim is to address this gap. Initially, we will focus on a generalization of the offset function  $p(\varrho) = \varrho^{\gamma}$  considered in [10], followed by an analysis of the well-posedness for two other forms: the singular offset function (1.9) and a non-local offset function defined as  $p(\varrho) = K(x) * \varrho$ . These variations are inspired not only by the aforementioned pedestrian flow model [1] but also by models that address lubrication effects [25] and collective behaviors [22, 31], as discussed in [11–13, 26] and related literature.

Lastly, it is important to mention that the well-posedness of the system (1.1) and (1.2) has been previously examined in the framework of weak solutions. Using the method of convex integration, it was shown in [9] that any initial data  $(\rho_0, \mathbf{u}_0) \in C^2(\mathbb{T}^3) \times C^3(\mathbb{T}^3)$  can connect to any terminal data  $(\rho_T, \mathbf{u}_T) \in C^2(\mathbb{T}^3) \times C^3(\mathbb{T}^3)$ consistent with mass and momentum conservation, via a weak solution belonging the class

$$\varrho \in C^2([0,T] \times \mathbb{T}^3), \quad \mathbf{u} \in L^\infty((0,T) \times \mathbb{T}^3).$$

The corresponding ill-posedness result clearly shows that the existence of so-called wild solutions extends beyond the hyperbolic systems of conservation laws and, in particular, to those experiencing dissipation that degenerates in vacuum.

In this article, we extend the energy estimates approach developed for the compressible Navier–Stokes equations (see [34], [35], [14], [15], [23], among others) to systems of mixed hyperbolic-parabolic type, which exhibit dissipation in the continuity equation but lack it in the momentum equation. We prove the local existence of regular solutions to system (1.1) by applying the method of successive approximations. We restrict ourselves to direct energy approach in  $L_2$  framework. Alternative  $L_p$  approach, which has been developed for the compressible Navier–Stokes system in [28]–[29] and, with entirely different techniques based on  $\mathcal{R}$ -bounded solution operators in [17], would be another possibility in context of regular solutions to (1.1), but we leave this direction for future investigation. The main difficulty here is to derive  $L_p$  estimates for a linear transport equation. The approach is based on an explicit solution formula obtained by the method of characteristics. Partial results of this type have been used in the theory of compressible Navier–Stokes equations (see among others [28], [34], [35], [23]), but a consistent  $L_p$  theory for transport equations is still missing. Here we address this issue proving quite a general result (lemma 2.2), which may be of independent interest. The dissipativity in (1.6b) gives parabolic estimates, but a delicate part is to ensure positivity of the solution at each step of the iteration. This issue is addressed in lemma 2.3.

## 1.1. Notation

Before stating our main result, we shall introduce the notation used in the article.

- Throughout the article, by  $E(\cdot)$  we denote a positive, continuous function such that E(0) = 0 and  $\phi(\cdot)$  denotes a continuous, positive function.
- We use standard notation  $L_p(\mathbb{T}^3)$  and  $W_p^1(\mathbb{T}^3)$  for Lebesgue and Sobolev spaces on the torus, respectively, and  $H^k(\mathbb{T}^3) := W_2^k(\mathbb{T}^3)$ . Next,  $L_p(0,T;X)$ , where X is a Banach space, denotes a Bochner space.
- For T > 0 and  $k \in \mathbb{N}$  let us also denote

$$\begin{aligned}
\mathcal{X}_{k}(T) &:= L_{2}(0,T; H^{k}(\mathbb{T}^{3})) \cap L_{\infty}(0,T; H^{k-1}(\mathbb{T}^{3})), \\
\mathcal{Y}_{k}(T) &:= \{f \in L_{\infty}(0,T; H^{k}(\mathbb{T}^{3})) : f_{t} \in L_{\infty}(0,T; H^{k-1}(\mathbb{T}^{3})), \\
\mathcal{V}_{k}(T) &:= \{g \in L_{2}(0,T; H^{k+1}(\mathbb{T}^{3})) \cap L_{\infty}(0,T; H^{k}(\mathbb{T}^{3})) : \\
g_{t} \in L_{2}(0,T; H^{k-1}(\mathbb{T}^{3}))\}
\end{aligned}$$
(1.10)

with norms defined in a natural way as appropriate sums of norms.

Since all spaces are considered on the torus, we shall sometimes skip indication of the domain in the definition of space and write  $L_p$  instead of  $L_p(\mathbb{T}^3)$ , etc. We are now in a position to state our main result.

THEOREM 1.1. Assume the initial data satisfies  $\varrho_0 > 0$ , and either (1.4) or (1.5). Assume, moreover, that the pressure is an increasing function of the density of class C<sup>5</sup>. Then there exists T > 0 such that system (1.7) admits a unique solution  $(\varrho, \mathbf{w}) \in \mathcal{V}_3(T) \times \mathcal{Y}_3(T)$  with the estimate

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\mathbf{w}\|_{\mathcal{Y}_{3}(T)} \le C(\|\varrho_{0}\|_{H^{4}(\mathbb{T}^{3})}, \|\mathbf{u}_{0}\|_{H^{3}(\mathbb{T}^{3})})$$

in case of (1.4) or

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\boldsymbol{w}\|_{\mathcal{V}_{3}(T)} \le C(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}, \|\boldsymbol{u}_{0} + \nabla p(\varrho_{0})\|_{H^{3}(\mathbb{T}^{3})})$$

in case of (1.5).

The strategy of the proof involves two main steps:

- construction of solutions to a suitable approximation of system (1.7),
- proof of convergence.

We aim to approximate solutions (1.7) by solutions to the iterative scheme defined as

$$\begin{cases} \varrho_t^{n+1} + \operatorname{div}\left(\varrho^{n+1}\mathbf{w}^{n+1}\right) = \operatorname{div}\left(\varrho^n p'(\varrho^n)\nabla\varrho^{n+1}\right),\\ \mathbf{w}_t^{n+1} + \mathbf{u}^n \cdot \nabla \mathbf{w}^{n+1} = 0,\\ (\varrho^{n+1}, \mathbf{w}^{n+1})|_{t=0} = (\varrho_0, \mathbf{u}_0 + \nabla p(\varrho_0)). \end{cases}$$
(1.11)

At each step of the iteration, having  $(\rho^n, \mathbf{w}^n)$  we set  $\mathbf{u}^n = \mathbf{w}^n - \nabla p(\rho^n)$  and solve the second equation of (1.11) for  $\mathbf{w}^{n+1}$ . Next we use  $\mathbf{w}^{n+1}$  to determine  $\rho^{n+1}$  from the first equation. Therefore, each step of iteration is decoupled to solving separate linear problems

$$\varrho_t + \operatorname{div}\left(\varrho \bar{\mathbf{w}}\right) = \operatorname{div}\left(\bar{\varrho}p'(\bar{\varrho})\nabla\varrho\right) \tag{1.12}$$

with given  $(\bar{\mathbf{w}}, \bar{\varrho})$  and

$$\mathbf{w}_t + \bar{\mathbf{u}} \cdot \nabla \mathbf{w} = 0 \tag{1.13}$$

with given  $\bar{\mathbf{u}}$ . Convergence of this iterative scheme is then proved using the Banach fixed point theorem.

The article is organized as follows. In §2, we first solve the linear problems corresponding to the iterative scheme described above in (1.12) and (1.13). Next, in §3, we prove the convergence of the iterative scheme using the contraction argument. Finally, in §4, we discuss the existence results for general singular and non-local offset functions; we formulate and prove our other main results—theorems 4.1 and 4.3.

#### 2. Linear theory

In this section, we solve linear problems corresponding to (1.12) and (1.13).

#### 2.1. Linear transport equation

Consider the linear transport equation

$$\eta_t + \mathbf{v} \cdot \nabla \eta = g \text{ on } \mathbb{T}^3 \times (0, T), \quad \eta|_{t=0} = \eta_0 \text{ on } \mathbb{T}^3$$
 (2.1)

with given vector field  $\mathbf{v}$  and unknown scalar valued  $\eta$ . Our goal is to prove the existence of a solution to (2.1) in the regularity framework corresponding to theorem 1.1. We will use the explicit form of the solution in Lagrangian coordinates given by

$$\frac{\partial X(t,y)}{\partial t} = \mathbf{v}(t, X(t,y)), \quad X(0,y) = y.$$
(2.2)

The first step is, therefore, to investigate the regularity properties of solutions to (2.2). For this purpose, we shall repeatedly use a basic fact that if a matrix is close to identity in the  $L_{\infty}$  in space-time norm, then the same holds for its inverse. In

particular, using the notation E(T), we have

$$\|A - \mathbb{I}\|_{L_{\infty}((0,T) \times \mathbb{T}^3)} \le E(T) \Longrightarrow \|A^{-1} - \mathbb{I}\|_{L_{\infty}((0,T) \times \mathbb{T}^3)} \le E(T)$$
(2.3)

for any function  $A: (0,T) \times \mathbb{T}^3 \to \mathbb{R}^{3 \times 3}$ . The following result improves [23, Lemma 3.2]:

LEMMA 2.1. Assume  $v \in L_2(0,T; H^3(\mathbb{T}^3))$ . Then there exists a continuous, positive function  $\phi(\cdot)$  denotes such that the solution to (2.2) satisfies

$$\|\nabla_y X - \mathbb{I}\|_{L_{\infty}((0,T) \times \mathbb{T}^3)} \le E(T), \qquad (2.4)$$

$$\|\nabla_y X\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \le \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}),$$
(2.5)

$$\|[\nabla_y X]^{-1} - \mathbb{I}\|_{L_{\infty}((0,T) \times \mathbb{T}^3)} \le E(T),$$

$$(2.6)$$

$$\|[\nabla_y X]^{-1}\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \le \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}),$$
(2.7)

$$\|\nabla_y X\|_{L_p(0,T;L_\infty(\mathbb{T}^3)} \le E(T) \quad \text{for } 1 \le p < \infty$$
(2.8)

$$\|\nabla_y^2 X\|_{L_{\infty}(0,T;L_6(\mathbb{T}^3))} \le \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))})$$
(2.9)

$$\|\nabla_y^2 X\|_{L_p(0,T;L_6(\mathbb{T}^3))} \le E(T) \quad \forall \ 1 \le p < \infty$$
(2.10)

$$\|\nabla_y^3 X\|_{L_{\infty}(0,T;L_2(\mathbb{T}^3))} \le \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))})$$
(2.11)

$$\|\nabla_y^3 X\|_{L_p(0,T;L_2(\mathbb{T}^3))} \le E(T) \quad \forall \ 1 \le p < \infty.$$
(2.12)

*Proof.* The first assertion was proved in [23, Lemma 3.2]. The second is derived and used in the proof of the aforementioned result. Then (2.6) results from (2.3), while (2.7) follows from the fact that the entries of the inverse an invertible matrix A are smooth functions of the entries of A

$$\|\nabla_y X\|_{L_p(0,T;L_\infty(\mathbb{T}^3))} \le T^{1/p} \|\nabla_y X\|_{L_\infty((0,T)\times\mathbb{T}^3)} \le E(T).$$

In order to prove the bounds for higher derivatives observe that differentiating the solution formula

$$X(t,y) = y + \int_0^t \mathbf{v}(s, X(s,y)) \,\mathrm{d}s$$

with respect to y we obtain

$$\nabla_y X(t,y) = y + \int_0^T \nabla_x \mathbf{v}(x, X(s,y)) \otimes \nabla_y X(s,y) \, \mathrm{d}s,$$

which is equivalent to

$$\partial_t \nabla_y X(t,y) = \nabla_x \mathbf{v}(x, X(t,y)) \otimes \nabla_y X(t,y).$$

Differentiating this identity in y, we obtain

$$\partial_t \nabla_y^2 X(t,y) \sim \nabla_x^2 \mathbf{v}(t, X(t,y)) (\nabla_y X(t,y))^2 + \nabla_x \mathbf{v}(t, X(t,y)) \nabla_y^2 X(t,y).$$
(2.13)

Multiplying the component corresponding to  $\partial_{y_i y_j}^2 X$  by  $|\partial_{y_i y_j}^2 X|^4 \partial_{y_i y_j}^2 X$ , summing over i, j and integrating over  $\mathbb{T}^3$  we get

$$\begin{aligned} \partial_t \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6 &\leq \int_{\mathbb{T}^3} |\nabla_x^2 \mathbf{v}(t, X(t, y))| (\nabla_y X(t, y))^2 |\nabla_y^2 X(t, y)|^5 \mathrm{d}y \\ &\int_{\mathbb{T}^3} |\nabla_x \mathbf{v}(t, X(t, y))| |\nabla_y^2 X(t, y)|^6 \mathrm{d}y \\ &\leq \| \nabla_x^2 \mathbf{v}(t, X(t, \cdot)) \|_{L_6(\mathbb{T}^3)} \| \nabla_y X(t, y) \|_{L_\infty(\mathbb{T}^3)}^2 \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^5 \\ &+ \| \nabla_x \mathbf{v}(t, X(t, \cdot)) \|_{L_\infty(\mathbb{T}^3)} \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6. \end{aligned}$$

$$(2.14)$$

By (2.4), for small T and any function f of the time variable with values in  $L_p(\mathbb{T}^3)$ for  $1 \leq p < \infty$ , we have

$$\begin{split} \|f(t,X(t,\cdot))\|_{L_{p}(\mathbb{T}^{3})} &= \left(\int_{\mathbb{T}^{3}} |f(t,X(t,y))|^{p} \, \mathrm{d}y\right)^{1/p} \\ &= \left(\int_{\mathbb{T}^{3}} |f(t,X(t,y))|^{p} |J_{y}X(t,y)| |J_{y}X(t,y)|^{-1} \, \mathrm{d}y\right)^{1/p} \\ &\leq \left(\sup_{y\in\mathbb{T}^{3}} |J_{y}X(t,y)|^{-1}\right)^{1/p} \left(\int_{\mathbb{T}^{3}} |f(t,x)|^{p} \, \mathrm{d}x\right)^{1/p} \leq C \|f(t,\cdot)\|_{L_{p}(\mathbb{T}^{3})}, \end{split}$$
(2.15)

and similarly

$$\|f(t, X(t, \cdot))\|_{L_{\infty}(\mathbb{T}^3)} \le C \|f(t, \cdot)\|_{L_{\infty}(\mathbb{T}^3)}.$$
(2.16)

By (2.15), (2.16), and Sobolev imbedding, applying (2.4) to the first term of the RHS of (2.14) we get

$$\begin{aligned} \partial_t \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6 &\leq \| \nabla_x^2 \mathbf{v}(t, \cdot) \|_{L_6(\mathbb{T}^3)} \| \nabla_y X(t, y) \|_{L_\infty(\mathbb{T}^3)}^2 \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^5 \\ &+ \| \nabla_x \mathbf{v}(t, \cdot) \|_{L_\infty(\mathbb{T}^3)} \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6 \\ &\leq C(\| \nabla_x^2 \mathbf{v}(t, X(t, \cdot)) \|_{L_6(\mathbb{T}^3)} \\ &+ \| \nabla_x \mathbf{v}(t, X(t, \cdot)) \|_{L_\infty(\mathbb{T}^3)}) \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6 \\ &\leq C \| \mathbf{v}(t, X(t, \cdot)) \|_{H^3(\mathbb{T}^3)} \| \nabla_y^2 X(t, \cdot) \|_{L_6(\mathbb{T}^3)}^6. \end{aligned}$$

The assumed integrability of  $\mathbf{v}$  allows to conclude (2.9) by Gronwall inequality:

$$\|\nabla_y^2 X(t,\cdot)\|_{L_6(\mathbb{T}^3)}^6 \le C \exp\left(\int_0^t \|\mathbf{v}(s,\cdot)\|_{H^3(\mathbb{T}^3)} \,\mathrm{d}s\right) \le \exp\left(\sqrt{t}\|\mathbf{v}\|_{L_2(0,t;H^3(\mathbb{T}^3))}\right).$$

This proves (2.9), which immediately implies (2.10). A remark is due here. In derivation of (2.9), we assumed for simplicity that

$$\|\nabla_y X\|_{L_{\infty}(\mathbb{T}^3)} \le C \|\nabla_y^2 X\|_{L_6(\mathbb{T}^3)}$$
(2.17)

which does not hold since we don't have Poincaré inequality. To make the proof fully precise, we would have to replace  $\|\nabla_y^2 X\|_{L_6(\mathbb{T}^3)}$  by  $\|\nabla_y X\|_{W_6^1(\mathbb{T}^3)}$  which is easy—it is enough to write estimate for  $\frac{\partial}{\partial t} \|\nabla_y X\|_{L_6(\mathbb{T}^3)}^6$ . Therefore to avoid additional obvious terms, we assume (2.17). Similar simplification is also used later in the proof.

In order to prove (2.11), we differentiate (2.13) in y obtaining

$$\partial_t \nabla_y^3 X(t,y) \sim \nabla_x^3 \mathbf{v}(t,X) (\nabla_y X)^3 + \nabla_x \mathbf{v}(t,X) \nabla_y X \nabla_y^2 X + \nabla_x \mathbf{v}(t,X) \nabla_y^2 X.$$

Multiplying the equation corresponding to  $\partial_{y_iy_jy_k}^3 X$  by  $\partial_{y_iy_jy_k}^3 X$  and summing over all i, j, k we get

$$\begin{split} \partial_t \| \nabla_y^3 X(t, \cdot) \|_{L_2(\mathbb{T}^3)}^2 &\leq \int_{\mathbb{T}^3} \nabla_x^3 \mathbf{v}(t, X(t, y)) |\nabla_y X(t, y)|^3 |\nabla_y^3 X(t, y)| \,\mathrm{d}y \\ &+ \int_{\mathbb{T}^3} |\nabla_x^2 \mathbf{v}(t, X(t, y))| \, |\nabla_y X(t, y)| |\nabla_y^2 X(t, y)| |\nabla_y^3 X(t, y)| \,\mathrm{d}y \\ &+ \int_{\mathbb{T}^3} \nabla_x \mathbf{v}(t, X(t, y)) |\nabla_y^3 X(t, y)|^2 \,\mathrm{d}y, \end{split}$$

from which, by Sobolev imbedding, (2.4), (2.15), and (2.16), we obtain

$$\begin{aligned} \partial_t \| \nabla_y^3 X(t, \cdot) \|_{L_2(\mathbb{T}^3)}^2 &\leq C \big( \| \nabla_x^3 \mathbf{v}(t, \cdot) \|_{L_2(\mathbb{T}^3)} + \| \nabla_x^2 \mathbf{v}(t, \cdot) \|_{L_6(\mathbb{T}^3)} \\ &+ \| \nabla_x \mathbf{v}(t, \cdot) \|_{L_\infty(\mathbb{T}^3)} \big) \| \nabla_y^3 X(t, \cdot) \|_{L_2(\mathbb{T}^3)}^2 \\ &\leq C \| \mathbf{v}(t, \cdot) \|_{H^3(\mathbb{T}^3)} \| \nabla_y^3 X(t, \cdot) \|_{L_2(\mathbb{T}^3)}^2, \end{aligned}$$

and by Gronwall inequality we conclude (2.11), which implies (2.12).

Now we are in a position to prove a series of estimates for the transport equation (2.1). As they may be of independent interest, we prove them in a possibly general form.

LEMMA 2.2. Assume  $v \in \mathcal{X}_3(T)$  defined in (1.10) and  $\eta_0 \in H^3(\mathbb{T}^3)$ . Then the solution to (2.1) satisfies

$$\begin{aligned} \|\eta\|_{L_{\infty}(0,T;W^{1}_{r}(\mathbb{T}^{3}))} &\leq \phi(\sqrt{T} \|\boldsymbol{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W^{1}_{r}(\mathbb{T}^{3})} + E(T) \|g\|_{L_{q}(0,T;W^{1}_{r}(\mathbb{T}^{3}))} \\ \forall 1 < q \leq \infty, 1 \leq r \leq \infty \end{aligned}$$
(2.18)

$$\|\nabla_x \eta\|_{L_p(0,T;L_r(\mathbb{T}^3))} \le E(T) \left( \|\nabla \eta_0\|_{L_\infty(\mathbb{T}^3)} + \|\nabla_x g\|_{L_1(0,T;L_r(\mathbb{T}^3))} \right)$$
  
$$\forall 1 \le p < \infty, \ 1 \le r \le \infty$$
(2.19)

Regular solutions to the dissipative Aw-Rascle system

$$\begin{aligned} \|\nabla_x^2 \eta\|_{L_{\infty}(0,T;L_6(\mathbb{T}^3))} &\leq \phi(\sqrt{T} \|\boldsymbol{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}) \|\eta_0\|_{W_6^2(\mathbb{T}^3)} \\ &+ E(T) \left( \|\nabla_x g\|_{L_q(0,T;W_6^1(\mathbb{T}^3))} \right) \quad \forall 1 < q \le \infty, \end{aligned}$$
(2.20)

$$\begin{aligned} \|\nabla_x^2 \eta\|_{L_p(0,T;L_6(\mathbb{T}^3))} &\leq E(T) \left( \|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;W_6^1(\mathbb{T}^3))} \right) \\ &\forall \ 1 \leq p \leq \infty, \ 1 < q \leq \infty, \end{aligned}$$
(2.21)

$$\begin{aligned} \|\nabla_x^2 \eta\|_{L_{\infty}(0,T;L_3(\mathbb{T}^3))} &\leq \phi(\sqrt{T} \|\boldsymbol{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}) \|\eta_0\|_{W_6^2(\mathbb{T}^3)} \\ &+ E(T) \left( \|\nabla_x g\|_{L_q(0,T;H^2(\mathbb{T}^3))} \right) \quad \forall 1 < q \leq \infty, \end{aligned}$$
(2.22)

$$\begin{aligned} \|\nabla_x^2 \eta\|_{L_p(0,T;L_3(\mathbb{T}^3))} &\leq E(T) \left( \|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;W_2^1(\mathbb{T}^3))} \right) \\ &\forall \, 1 \leq p < \infty, \, 1 < q \leq \infty, \end{aligned}$$
(2.23)

$$\begin{aligned} \|\eta\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))} &\leq \phi(\sqrt{T} \|\boldsymbol{v}\|_{\mathcal{X}_{3}(T)}) \|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} \\ &+ E(T) \left( \|\eta_{0}\|_{H^{3}(\mathbb{T}^{3})} + \|g\|_{L_{q}(0,T;H^{3}(\mathbb{T}^{3}))} \right) \quad \forall 1 < q \leq \infty, \end{aligned}$$
(2.24)

$$\begin{aligned} \|\eta\|_{L_p(0,T;H^3(\mathbb{T}^3))} &\leq E(T) \left( \|\eta_0\|_{H^3(\mathbb{T}^3)} + \|g\|_{L_q(0,T;H^3(\mathbb{T}^3))} \right), \\ \forall 1 \leq p < \infty, \ 1 < q \leq \infty, \end{aligned}$$
(2.25)

$$\begin{aligned} \|\eta_t\|_{L_{\infty}(0,T;H^2(\mathbb{T}^3))} &\leq \phi(\sqrt{T} \|\boldsymbol{v}\|_{\mathcal{X}_3(T)}) (\|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla g\|_{L_{\infty}(0,T;H^1(\mathbb{T}^3))}) \\ &+ E(T) \left(\|\eta_0\|_{H^3(\mathbb{T}^3)} + \|g\|_{L_q(0,T;H^3(\mathbb{T}^3))}\right) \quad \forall 1 < q \leq \infty, \end{aligned}$$

$$(2.26)$$

 $\begin{aligned} \|\eta_t\|_{L_p(0,T;H^2(\mathbb{T}^3))} &\leq E(T) \left( \|\nabla\eta_0\|_{H^3(\mathbb{T}^3)} + \|g\|_{L_q(0,T;H^3(\mathbb{T}^3))} + \|\nabla g\|_{L_\infty(0,T;H^1(\mathbb{T}^3))} \right) \\ \forall 1 \leq p < \infty, \ 1 < q \leq \infty. \end{aligned}$ (2.27)

*Proof.* We have

$$\eta(t, X(t, y)) = \eta_0(y) + \int_0^t g(s, X(s, y)) \, \mathrm{d}s.$$

Differentiating this identity in y, we obtain

$$\nabla_y X(t,y) \nabla_x \eta(t, X(t,y)) = \nabla_y \eta_0 + \int_0^t \nabla_x g(s, X(s,y)) \nabla_y X(s,y) \,\mathrm{d}s, \qquad (2.28)$$

which, by lemma 2.1, implies

$$\begin{aligned} \|\nabla_{x}\eta\|_{L_{\infty}(0,T;L_{r}(\mathbb{T}^{3}))} &\leq \|(\nabla_{y}X)^{-1}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla_{y}\eta_{0}\|_{L_{r}(\mathbb{T}^{3})} \\ &+ \|(\nabla_{y}X)^{-1}\int_{0}^{t}\nabla_{x}g(s,X(s,y))\nabla_{y}X(s,y)\,\mathrm{d}s\|_{L_{\infty}(0,T;L_{r}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))})\|\nabla\eta_{0}\|_{L_{r}(\mathbb{T}^{3})} + E(T)\|\nabla_{x}g\|_{L_{q}(0,T;L_{r}(\mathbb{T}^{3}))} \\ &\forall 1 < q \leq \infty, \quad 1 \leq r \leq \infty \end{aligned}$$

$$(2.29)$$

from which we obtain (2.18). Similarly, using (2.8), we obtain (2.19). In order to estimate  $\nabla_x^2 \eta$ , we differentiate (2.28) in y to obtain

$$\begin{split} &\frac{\partial^2 \eta_0(y)}{\partial y_j \partial y_k} + \int_0^t \sum_i \left[ g_{x_i}(s, X(s, y)) \frac{\partial^2 X_i(s, y)}{\partial y_j \partial y_k} \right. \\ &+ \left( \sum_l g_{x_i x_l}(s, X(s, y)) \frac{\partial X_l(s, y)}{\partial y_k} \right) \frac{\partial X_i(s, y)}{\partial y_j} \right] = \\ &= \partial_{y_k} \left( \sum_i \frac{\partial \eta}{\partial x_i}(t, X(t, y)) \frac{\partial X_i(t, y)}{\partial y_j}(t, y) \right) = \\ &= \sum_{i,l} \frac{\partial^2 \eta}{\partial x_i \partial x_l}(t, X(t, y)) \frac{\partial X_l}{\partial y_k}(t, y) \frac{\partial X_i}{\partial y_j}(t, y) + \nabla_x \eta(t, X(t, y)) \cdot \frac{\partial^2 X}{\partial y_j \partial y_k}(t, y) \end{split}$$

for  $j, k \in \{1, 2, 3\}$ . Rewriting the above system as

$$\sum_{i,l} \frac{\partial^2 \eta}{\partial x_i \partial x_l} (t, X(t, y)) \frac{\partial X_l}{\partial y_k} (t, y) \frac{\partial X_i}{\partial y_j} (t, y) =$$

$$= \frac{\partial^2 \eta_0(y)}{\partial y_j \partial y_k} + \int_0^t \sum_i \left[ g_{x_i}(s, X(s, y)) \frac{\partial^2 X_i(s, y)}{\partial y_j \partial y_k} + \left( \sum_l g_{x_i x_l}(s, X(s, y)) \frac{\partial X_l(s, y)}{\partial y_k} \right) \frac{\partial X_i}{\partial y_j} \right]$$

$$- \nabla_x \eta(t, X(t, y)) \cdot \frac{\partial^2 X(t, y)}{\partial y_i \partial y_k}$$
(2.30)

for  $k, j \in \{1, 2, 3\}$ , which is a linear system of nine equations for the unknown derivatives  $\frac{\partial^2 \eta}{\partial x_i \partial x_l}(t, X)$ . In order to solve it, we observe that the diagonal of this system corresponds to (i, l) = (j, k), which means that on the diagonal we have terms  $\frac{\partial X_k}{\partial y_k} \frac{\partial X_j}{\partial y_j}$ , while all entries outside the diagonal contains the terms which are not on the diagonal of  $\nabla_y X$ . Therefore, by (2.4), all terms on the diagonal of system (2.30) are close to 1 for short times, while all other terms are small.

## Regular solutions to the dissipative Aw-Rascle system

Therefore, system (2.30) is uniquely solvable and we obtain

$$\begin{aligned} |\nabla_x^2 \eta(t, X(t, y))| &\leq C \Big( |\nabla_y^2 \eta_0| + |\nabla_x \eta(t, X(t, y))| |\nabla_y^2 X(t, y)| \\ &+ \left| \int_0^t |\nabla_x g(s, X(s, y))| |\nabla_y^2 X(s, y)| + |\nabla_x^2 g(s, X(s, y))| |\nabla_y X(s, y)|^2 \, \mathrm{d}s \right| \Big). \end{aligned}$$
(2.31)

By (2.9) and (2.29), we have

$$\begin{split} \||\nabla_{x}\eta||\nabla_{y}^{2}X|\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} &\leq \|\nabla_{x}\eta\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla_{y}^{2}X\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \\ &\leq \left(\phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))})\|\nabla\eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + E(T)\|\nabla_{x}g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))}\right) \\ &\phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \\ &\leq \phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))})\|\nabla\eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + E(T)\|\nabla_{x}g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))} \quad \forall \ 1 < q \leq \infty. \end{aligned}$$

$$(2.32)$$

Next, by (2.9)

$$\begin{aligned} \left\| \int_{0}^{t} |\nabla_{x}g| |\nabla_{y}^{2}X| \right\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} &\leq \int_{0}^{T} \|\nabla_{x}g\|_{L_{\infty}(\mathbb{T}^{3})} \|\nabla_{y}^{2}X\|_{L_{6}(\mathbb{T}^{3})} \\ &\leq \phi(\sqrt{t}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \int_{0}^{T} \|\nabla_{x}g\|_{L_{\infty}(\mathbb{T}^{3})} \mathrm{d}t \leq E(T) \|\nabla_{x}g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))} \end{aligned}$$

$$\forall 1 < q \leq \infty, \end{aligned}$$

$$(2.33)$$

and finally

$$\begin{aligned} \left\| \int_{0}^{t} |\nabla_{x}^{2}g| |\nabla_{y}X|^{2} \mathrm{d}t \right\|_{L_{\infty}(0,T;L_{p}(\mathbb{T}^{3}))} &\leq \int_{0}^{T} \|\nabla_{x}^{2}g\|_{L_{p}(\mathbb{T}^{3})} \|\nabla_{y}X\|_{L_{\infty}}^{2} \\ &\leq \|\nabla_{y}X\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{2} \int_{0}^{T} \|\nabla_{x}^{2}g\|_{L_{p}(\mathbb{T}^{3})} \mathrm{d}t \leq E(T) \|\nabla_{x}^{2}g\|_{L_{q}(0,T;L_{p}(\mathbb{T}^{3}))} \\ &\forall \ 1 < q \leq \infty, \ 1 \leq p \leq 6. \end{aligned}$$

$$(2.34)$$

Combining (2.31), (2.32), (2.33), and (2.34), we obtain (2.20). Next, by (2.10) and (2.29), we have

$$\begin{aligned} \left\| |\nabla_x \eta| |\nabla_y^2 X| \right\|_{L_p(0,T;L_6(\mathbb{T}^3))} &\leq \|\nabla_x \eta\|_{L_\infty((0,T)\times\mathbb{T}^3)} \|\nabla_y^2 X\|_{L_p(0,T;L_6(\mathbb{T}^3))} \\ &\leq E(T)(\text{RHS of } (2.29)) \leq E(T) \left( \|\nabla\eta_0\|_{L_\infty(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;L_\infty(\mathbb{T}^3))} \right) \\ &\forall 1 < q \leq \infty. \end{aligned}$$

Combining this estimate with (2.31), (2.33), and (2.34), we arrive at (2.21). Next, similarly to (2.32), we obtain

$$\begin{split} &\|\nabla_{x}\eta\|\nabla_{y}^{2}X\|\|_{L_{\infty}(0,T;L_{3}(\mathbb{T}^{3}))} \leq \|\nabla_{y}^{2}X\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))}\|\nabla_{x}\eta\|_{L_{1}(0,T;L_{6}(\mathbb{T}^{3}))} \\ &\leq \left(\phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))})\|\nabla\eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + E(T)\|\nabla_{x}\eta\|_{L_{q}(0,T;L_{6}(\mathbb{T}^{3}))}\right)\phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \\ &\leq \phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))})\|\nabla\eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + E(T)\|\nabla_{x}g\|_{L_{q}(0,T;H^{1}(\mathbb{T}^{3}))} \quad \forall \ 1 < q \leq \infty \end{split}$$

$$(2.35)$$

and, in analogy to (2.33), we have

$$\begin{aligned} \left\| \int_{0}^{t} |\nabla_{x}g| |\nabla_{y}^{2}X| \right\|_{L_{\infty}(0,T;L_{3}(\mathbb{T}^{3}))} &\leq \int_{0}^{T} \|\nabla_{x}g\|_{L_{6}(\mathbb{T}^{3})} \|\nabla_{y}^{2}X\|_{L_{6}(\mathbb{T}^{3})} \\ &\leq \phi(\sqrt{t}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \int_{0}^{T} \|\nabla_{x}g\|_{L_{6}(\mathbb{T}^{3})} \mathrm{d}t \leq E(T) \|\nabla_{x}g\|_{L_{q}(0,T;H^{1}(\mathbb{T}^{3}))} \\ &\forall 1 < q \leq \infty. \end{aligned}$$

$$(2.36)$$

Combining (2.35), (2.36), and (2.34) with p = 2, we obtain (2.22). Next, by (2.10) and (2.29), we have

$$\begin{split} \left\| |\nabla_x \eta| |\nabla_y^2 X| \right\|_{L_p(0,T;L_3(\mathbb{T}^3))} &\leq \|\nabla_x \eta\|_{L_\infty(L_6(\mathbb{T}^3))} \|\nabla_y^2 X\|_{L_p(0,T;L_6(\mathbb{T}^3))} \\ &\leq E(T)(\text{RHS of } (2.29)) \leq E(T) \left( \|\nabla\eta_0\|_{L_\infty(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;L_6(\mathbb{T}^3))} \right) \quad \forall \ 1 < q \leq \infty, \end{split}$$

which combined with (2.36) and (2.34) for p = 2 gives (2.23).

In order to estimate the third order derivatives, we differentiate (2.30) w.r.t.  $y_m$ , which yields

$$\begin{split} &\sum_{i,l,n} \frac{\partial^3 \eta}{\partial x_i \partial x_l \partial x_n} (t, X) \frac{\partial X_n}{\partial y_m} \frac{\partial X_l}{\partial y_k} \frac{\partial X_i}{\partial y_j} = \\ &= \frac{\partial^3 \eta_0}{\partial y_j \partial y_k \partial y_m} - \partial_{y_m} \left( \nabla_x \eta(t, X) \frac{\partial^2 X}{\partial y_j \partial y_k} (t, y) \right) \\ &- \sum_{i,l} \frac{\partial^2 \eta}{\partial x_i \partial x_l} (t, X) \left( \frac{\partial^2 X_l}{\partial y_k \partial y_m} \frac{\partial X_i}{\partial y_l} + \frac{\partial^2 X_i}{\partial y_l \partial y_m} \frac{\partial X_l}{\partial y_k} \right) (t, y) \\ &+ \int_0^t \partial_{y_m} \left\{ \sum_i \left[ g_{x_i}(s, X) \frac{\partial^2 X_i}{\partial y_j \partial y_k} + \left( \sum_l g_{x_i x_l}(s, X) \frac{\partial X_l}{\partial y_k} \right) \frac{\partial X_i}{\partial y_j} \right] \right\}, \end{split}$$
(2.37)

where X = X(t, y) or X = X(s, y) according to (2.30). Similarly as in case of (2.30), it is a system of 27 linear equations for the third order derivatives of  $\eta$ . On the diagonal, we have terms corresponding to (i, l, n) = (j, k, m), which, again by (2.4), are close to one, while all other entries are small for small times. Therefore, (2.37) is uniquely solvable and we obtain

$$\begin{aligned} |\nabla_{x}^{3}\eta| &\leq C\left(|\nabla_{y}^{3}\eta_{0}| + |\nabla_{x}^{2}\eta||\nabla_{y}X||\nabla_{y}^{2}X| + |\nabla_{x}\eta||\nabla_{y}^{3}X| \\ &+ \int_{0}^{t} |\nabla_{x}g||\nabla_{y}^{3}X| + |\nabla_{x}^{2}g||\nabla_{y}X||\nabla_{y}^{2}X| + |\nabla_{x}^{3}g||\nabla_{y}X|^{3} \mathrm{d}t\right). \end{aligned}$$
(2.38)

Let us estimate the RHS of (2.38). For the second term, by (2.9) and (2.20), we have

$$\begin{split} \left\| |\nabla_{x}^{2} \eta| |\nabla_{y} X| |\nabla_{y}^{2} X| \right\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq C \|\nabla_{x}^{2} \eta\|_{L_{\infty}(0,T;L_{3}(\mathbb{T}^{3}))} \|\nabla_{y}^{2} X\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{t} \|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|RHS \text{ of } (2.20)] \\ &\leq \phi(\sqrt{t} \|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} + E(T) \left( \|\nabla_{x} g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))} \right) + \\ & \|\nabla_{x}^{2} g\|_{L_{q}(0,T;L_{6}(\mathbb{T}^{3}))}), \end{split}$$

$$(2.39)$$

and for the third, by (2.11) and (2.18)

$$\begin{split} & \left\| |\nabla_x \eta| |\nabla_y^3 X| \right\|_{L_{\infty}(0,T;L_2(\mathbb{T}^3))} \le \|\nabla_x \eta\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \|\nabla_y^3 X\|_{L_{\infty}(0,T;L_2(\mathbb{T}^3))} \\ & \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}) \text{ [RHS of (2.18)]} \\ & \le \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}) \|\eta_0\|_{W^1_{\infty}} + E(T) \|g\|_{L_q(0,T;W^1_{\infty}(\mathbb{T}^3))}. \end{split}$$

It remains to estimate the terms with g. By (2.11), we have

$$\begin{split} \left\| \int_{0}^{t} |\nabla_{x}g| |\nabla_{y}^{3}X| \right\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq \int_{0}^{T} \|\nabla_{x}g(t,\cdot)\|_{L_{\infty}} \|\nabla_{y}^{3}X(t,\cdot)\|_{L_{2}(\mathbb{T}^{3})} \mathrm{d}t \\ &\leq E(T) \|\nabla_{x}g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))} \quad \forall q < 1 \leq \infty. \end{split}$$

$$(2.40)$$

Next, by (2.9),

$$\begin{split} & \left\| \int_{0}^{t} |\nabla_{x}^{2}g| |\nabla_{y}X| |\nabla_{y}^{2}X| \right\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} \leq C \int_{0}^{T} \|\nabla_{x}^{2}g(t,\cdot)\|_{L_{3}(\mathbb{T}^{3})} \|\nabla_{y}^{2}X(t,\cdot)\|_{L_{6}(\mathbb{T}^{3})} \\ & \leq C \|\nabla_{y}^{2}X\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \int_{0}^{T} \|\nabla_{x}^{2}g(t,\cdot)\|_{L_{3}(\mathbb{T}^{3})} \mathrm{d}t \leq E(T) \|\nabla_{x}^{2}g\|_{L_{q}(0,T;L_{3}(\mathbb{T}^{3}))} \quad \forall 1 < q \leq \infty, \\ & (2.41) \end{split}$$

and finally

$$\begin{aligned} \left\| \left| \nabla_x^3 g \right| \left| \nabla_y X \right|^3 \right| \right\|_{L_{\infty}(0,T;L_2(\mathbb{T}^3))} &\leq C \int_0^T \| \nabla_x^3 g(t,\cdot) \|_{L_2(\mathbb{T}^3)} \mathrm{d}t \leq E(T) \| \nabla_x^3 g \|_{L_q(0,T;L_2(\mathbb{T}^3))} \\ \forall 1 < q \leq \infty. \end{aligned}$$

$$(2.42)$$

Combining (2.38)–(2.42) and applying Sobolev imbedding to estimate all terms containing g by a single norm, we obtain

$$\begin{aligned} \|\nabla_x^3 \eta\|_{L_{\infty}(0,T;L_2(\mathbb{T}^3))} &\leq \phi(\sqrt{T} \|\mathbf{v}\|_{L_2(0,T;H^3(\mathbb{T}^3))}) \|\eta_0\|_{W_6^2(\mathbb{T}^3)} \\ &+ E(T) \left( \|g\|_{L_q(0,T;H^3(\mathbb{T}^3))} \right) \quad \forall \, 1 < q \le \infty, \end{aligned}$$

which together with estimates on lower order derivatives of  $\eta$  completes the proof of (2.24).

In order to prove of (2.25) observe that, for any finite p, by (2.9) and (2.21) we have

$$\begin{aligned} \|\nabla_x^2 \eta \nabla_y X \nabla_y^2 X\|_{L_p(0,T;L_2(\mathbb{T}^3))} &\leq C \|\nabla_y^2 X\|_{L_\infty(0,T;L_6(\mathbb{T}^3))} \|\nabla_x^2 \eta(t)\|_{L_p(0,T;L_3(\mathbb{T}^3))} \\ &\leq E(T) \left( \|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;L_\infty)} + \|\nabla_x^2 g\|_{L_q(0,T;L_6(\mathbb{T}^3))} \right) \quad \forall \ 1 < q \leq \infty. \end{aligned}$$

Similarly by (2.11) and (2.19), we obtain

$$\begin{aligned} \|\nabla_{x}\eta\nabla_{y}^{3}X\|_{L_{p}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq C\|\nabla_{y}^{3}X\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}\|\nabla_{x}\eta\|_{L_{p}(0,T;L_{\infty}(\mathbb{T}^{3}))} \\ &\leq E(T)\left(\|\nabla\eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + \|\nabla_{x}g\|_{L_{1}(0,T;L_{\infty}(\mathbb{T}^{3}))}\right). \end{aligned}$$
(2.44)

For the terms with g on the RHS of (2.38), we use the estimates (2.40)–(2.42). Combining them with (2.43)–(2.44), we obtain

$$\|\nabla_x^3\eta\|_{L_p(0,T;L_2(\mathbb{T}^3))} \le E(T) \left( \|\eta_0\|_{H^3(\mathbb{T}^3)} + \|g\|_{L_q(0,T;H^3(\mathbb{T}^3))} \right) \quad \forall 1 \le p < \infty, \ 1 < q \le \infty,$$

which completes the proof of (2.25). Now we can use (2.1) to prove the estimates for  $\eta_t$ . First we immediately get

$$\|\eta_t\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \le \phi(\sqrt{t}\|\mathbf{v}\|_{\mathcal{X}_3(T)}), \quad \|\eta_t\|_{L_p(0,T;L_{\infty}(\mathbb{T}^3))} \le E(T) \quad \forall \ 1 \le p < \infty.$$

Next we differentiate (2.1) in the space variable to obtain

$$\nabla \eta_t \sim \nabla \mathbf{v} \nabla \eta + \mathbf{v} \nabla^2 \eta + \nabla g. \tag{2.45}$$

By (2.18), we have

$$\begin{aligned} \|\nabla \mathbf{v} \nabla \eta\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} &\leq \|\nabla \mathbf{v}\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \|\nabla \eta\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \\ &\leq \phi(\sqrt{T} \|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W^{1}_{\infty}(\mathbb{T}^{3})} + E(T) \|g\|_{L_{q}(0,T;W^{1}_{\infty}(\mathbb{T}^{3}))} \quad \forall 1 < q \leq \infty, \end{aligned}$$

$$(2.46)$$

by (2.20)

$$\begin{aligned} \|\mathbf{v}\nabla^{2}\eta\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} &\leq \|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla^{2}\eta\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{T}\|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} + E(T) \left(\|\nabla_{x}g\|_{L_{q}(0,T;L_{\infty}(\mathbb{T}^{3}))} + \|\nabla_{x}^{2}g\|_{L_{q}(0,T;L_{6}(\mathbb{T}^{3}))}\right) \quad \forall 1 < q \leq \infty, \end{aligned}$$

$$(2.47)$$

by (2.19)

$$\begin{aligned} \|\nabla \mathbf{v} \nabla \eta(t)\|_{L_{p}(0,T;L_{6}(\mathbb{T}^{3}))} &\leq \|\nabla \mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla \eta(t)\|_{L_{p}(0,T;L_{6}(\mathbb{T}^{3}))} \\ &\leq E(T) \left(\|\nabla \eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + \|\nabla_{x}g\|_{L_{1}(0,T;L_{\infty}(\mathbb{T}^{3}))}\right) \quad \forall 1 \leq p < \infty, \end{aligned}$$
(2.48)

and finally by (2.21)

$$\int_{0}^{T} \|\mathbf{v}(t,\cdot)\nabla^{2}\eta(t,\cdot)\|_{L_{6}(\mathbb{T}^{3})}^{p} dt \leq \|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{p} \int_{0}^{T} \|\nabla^{2}\eta(t,\cdot)\|_{L_{6}(\mathbb{T}^{3})}^{p} dt \\
\leq E(T) \left(\|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} + \|\nabla_{x}g\|_{L_{q}(0,T;W_{6}^{1}(\mathbb{T}^{3}))}\right) \quad \forall 1 \leq p < \infty, \ 1 < q \leq \infty,$$
(2.49)

14

so altogether we obtain

$$\begin{split} \|\nabla\eta_t\|_{L_{\infty}(0,T;L_6(\mathbb{T}^3))} &\leq \phi(\|\mathbf{v}\|_{\mathcal{X}_3(T)})(\|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla_x g\|_{L_{\infty}(0,T;L_6(\mathbb{T}^3))}) \\ &+ E(T)\|\nabla_x g\|_{L_q(0,T;W_6^1(\mathbb{T}^3))} \quad \forall 1 < q \leq \infty, \\ \|\nabla\eta_t\|_{L_p(0,T;L_6(\mathbb{T}^3))} &\leq E(T)\left(\|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla g\|_{L_q(0,T;W_6^1(\mathbb{T}^3))} \\ &+ \|\nabla g\|_{L_{\infty}(0,T;L_6(\mathbb{T}^3))}\right) \quad \forall 1 \leq p < \infty, 1 < q \leq \infty. \end{split}$$

Finally we differentiate (2.45) once more in space:

$$\nabla^2 \eta_t \sim \nabla^2 \mathbf{v} \nabla \eta + \nabla \mathbf{v} \nabla^2 \eta + \mathbf{v} \nabla^3 \eta + \nabla^2 g_t$$

For the first term we have, by (2.18),

$$\begin{aligned} \|\nabla^{2} \mathbf{v} \nabla \eta\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq \|\nabla \eta\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla^{2} \mathbf{v}\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{T} \|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W^{1}_{\infty}(\mathbb{T}^{3})} + E(T) \|g\|_{L_{q}(0,T;W^{1}_{\infty}(\mathbb{T}^{3}))} \quad \forall 1 < q \leq \infty, \end{aligned}$$

$$(2.50)$$

and, by (2.19),

$$\begin{aligned} \|\nabla^{2} \mathbf{v} \nabla \eta(t)\|_{L_{p}(0,T;L_{2}(\mathbb{T}^{3}))}^{p} &\leq \|\nabla \eta\|_{L_{p}(0,T;L_{\infty}(\mathbb{T}^{3}))} \int_{0}^{T} \|\nabla^{2} \mathbf{v}(t)\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq E(T) \left(\|\nabla \eta_{0}\|_{L_{\infty}(\mathbb{T}^{3})} + \|\nabla_{x}g\|_{L_{1}(0,T;L_{\infty}(\mathbb{T}^{3}))}\right) \quad \forall \ 1 \leq p < \infty. \end{aligned}$$

$$(2.51)$$

For the second term, by (2.20),

$$\begin{aligned} \|\nabla \mathbf{v} \nabla^{2} \eta\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq \|\nabla \mathbf{v}\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))} \|\nabla^{2} \eta\|_{L_{\infty}(0,T;L_{3}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{T} \|\mathbf{v}\|_{L_{2}(0,T;H^{3}(\mathbb{T}^{3}))}) \|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} + E(T) \left( \|\nabla_{x}g\|_{L_{q}(0,T;W_{6}^{1}(\mathbb{T}^{3}))} \right) \quad \forall 1 < q \leq \infty \end{aligned}$$

$$(2.52)$$

and by (2.21)

$$\begin{aligned} \|\nabla \mathbf{v} \nabla^2 \eta(t)\|_{L_p(0,T;L_2(\mathbb{T}^3))} &\leq \|\nabla \mathbf{v}\|_{L_\infty(0,T;L_6(\mathbb{T}^3))} \|\nabla^2 \eta(t)\|_{L_p(0,T;L_6(\mathbb{T}^3))} \\ &\leq E(T) \left( \|\eta_0\|_{W_6^2(\mathbb{T}^3)} + \|\nabla_x g\|_{L_q(0,T;W_6^1(\mathbb{T}^3))} \right) \quad \forall \ 1 \leq p < \infty, \ 1 < q \leq \infty. \end{aligned}$$
(2.53)

Finally, to estimate the last term, we apply (2.24) to get

$$\begin{aligned} \|\mathbf{v}\nabla^{3}\eta\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq \|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla^{3}\eta\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq \phi(\sqrt{t}\|\mathbf{v}\|_{\mathcal{X}_{3}(T)})\|\eta_{0}\|_{W_{6}^{2}(\mathbb{T}^{3})} + E(T)\left(\|\eta_{0}\|_{H^{3}(\mathbb{T}^{3})} + \|g\|_{L_{q}(0,T;H^{3}(\mathbb{T}^{3}))}\right) \quad \forall 1 < q \leq \infty \end{aligned}$$

$$(2.54)$$

and (2.25) to obtain

$$\|\mathbf{v}\nabla^{3}\varrho(t)\|_{L_{p}(0,T;L_{2}(\mathbb{T}^{3}))} \leq \|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{p} \int_{0}^{T} \|\nabla^{3}\eta(t)\|_{L_{p}(0,T;L_{2}(\mathbb{T}^{3}))}$$
  
$$\leq E(T)\left(\|\eta_{0}\|_{H^{3}(\mathbb{T}^{3})} + \|g\|_{L_{q}(0,T;H^{3}(\mathbb{T}^{3}))}\right) \quad \forall 1 \leq p < \infty, \ 1 < q \leq \infty.$$

$$(2.55)$$

Combining (2.46), (2.47), (2.50), (2.52), and (2.54), we obtain (2.26). Finally, (2.48), (2.49), (2.51), (2.53), and (2.55) allow to conclude (2.27), which completes the proof.

#### 2.2. Linear continuity equation with dissipation

In this section, we investigate the linear problem

$$\varrho_t + \operatorname{div}(\varrho \mathbf{v}) - \operatorname{div}(\mathfrak{a}\nabla \varrho) = \mathfrak{b}, \quad \varrho|_{t=0} = \varrho_0.$$
(2.56)

Concerning the regularity of the data, we keep in mind that the above system corresponds to the first equation of (1.7). Therefore, taking into account lemma 2.2, it is sufficient to assume  $\mathbf{v} \in L_{\infty}(0, T; H^3(\mathbb{T}^3))$ . The parabolic maximal regularity then leads to the following result.

LEMMA 2.3. Assume  $\varrho_0 \in H^2(\mathbb{T}^3)$ ,  $v \in L_\infty(0,T;H^2(\mathbb{T}^3))$ ,  $\mathfrak{a} \in L_\infty(0,T;W^1_\infty(\mathbb{T}^3)) \cap L_2(0,T;H^2(\mathbb{T}^3))$  with  $\mathfrak{a} \ge c > 0$ , and  $\mathfrak{b} \in L_2(0,T;H^1(\mathbb{T}^3))$ . Then (2.56) admits a unique solution satisfying

$$\begin{split} \|\varrho\|_{\mathcal{V}_{2}(T)} &\leq C \Big[ \|\varrho_{0}\|_{H^{2}(\mathbb{T}^{3})} + \|\mathfrak{b}\|_{L_{2}(0,T;H^{1}(\mathbb{T}^{3}))} \\ &+ \|\varrho\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \Big( T \big( \|\boldsymbol{v}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\nabla\mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \big) + \|\nabla^{2}\mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} ] \Big) \Big]. \end{split}$$

$$(2.57)$$

$$If \ \varrho_{0} \ \in \ H^{3}(\mathbb{T}^{3}), \ \mathfrak{b} \ \in \ L_{2}(0,T;H^{2}(\mathbb{T}^{3})) \ \boldsymbol{v} \ \in \ L_{\infty}(0,T;H^{3}(\mathbb{T}^{3})) \ and \ \mathfrak{a} \ \in \ L_{\infty}(W^{1}_{\infty}(\mathbb{T}^{3})) \cap L_{2}(0,T;H^{3}(\mathbb{T}^{3})) \ with \ \mathfrak{a} \ge c > 0, \ then$$

$$\begin{aligned} \|\varrho\|_{\mathcal{V}_{3}(T)} &\leq C \Big[ \|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})} + \|\mathfrak{b}\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))} \\ &+ \|\varrho\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))} \Big( T \Big( \|\boldsymbol{v}\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))} + \|\nabla\mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \Big) + \|\nabla^{2}\mathfrak{a}\|_{L_{2}(0,T;H^{1}(\mathbb{T}^{3}))} \Big] \Big) \Big]. \end{aligned}$$

$$(2.58)$$

Moreover, if we assume

$$\mathfrak{b} = \operatorname{div} (\varrho \mathbf{b}) \quad \text{with} \quad \|\operatorname{div} \mathbf{b}\|_{L_1(0,T;L_\infty(\mathbb{T}^3))} \le C \tag{2.59}$$

and the initial data is strictly positive, i.e.,  $\inf_{x \in \mathbb{T}^3} \varrho_0 > 0$  then we have

$$\min_{(t,x)\in[0,T]\times\mathbb{T}^3} \varrho(t,x) > 0.$$
(2.60)

REMARK 1. Clearly, the constant C in (2.57) and (2.58) depends on c and it goes to infinity when  $c \rightarrow 0$ .

*Proof.* Rewriting (2.56) as

$$\varrho_t - \operatorname{div}\left(\mathfrak{a}\nabla\varrho\right) = -\operatorname{div}\left(\varrho\mathbf{v}\right) + \mathfrak{b}$$
(2.61)

we immediately obtain the bound

$$\begin{aligned} \|\varrho\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))} &+ \|\varrho\|_{L_{\infty}(0,T;H^{1}(\mathbb{T}^{3}))} + \|\varrho_{t}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq C\left(\|\varrho_{0}\|_{H^{1}(\mathbb{T}^{3})} + \|\operatorname{div}\left(\varrho\mathbf{v}\right)\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} + \|\mathfrak{b}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}\right). \end{aligned}$$

$$(2.62)$$

The above estimate (2.62) is a direct consequence of parabolic regularity theorem in  $L^2$ - setup (see Ladyženskaya, Solonikov, and Ural'ceva [24, Chapter 3], even for more general boundary conditions). Differentiating (2.61) in  $x_i$ , we obtain

$$(\varrho_{x_i})_t - \operatorname{div}\left(\mathfrak{a}\nabla \varrho_{x_i}\right) = -(\operatorname{div}\left(\varrho\mathbf{v}\right))_{x_i} + \operatorname{div}\left(\mathfrak{a}_{x_i}\nabla \varrho\right) + \mathfrak{b}_{x_i} =: F_1^i,$$

therefore

$$\begin{aligned} \|\varrho_{x_i}\|_{L_2(0,T;H^2(\mathbb{T}^3))} + \|\varrho_{x_i}\|_{L_\infty(0,T;H^1(\mathbb{T}^3))} + \|\partial_t \varrho_{x_i}\|_{L_2(0,T;L_2(\mathbb{T}^3))} \\ &\leq C\left(\|\varrho_0\|_{H^2(\mathbb{T}^3)} + \|F_1^i\|_{L_2(0,T;L_2(\mathbb{T}^3))}\right). \end{aligned}$$
(2.63)

Under the assumed regularity of **v** and  $\mathfrak{a}$ , using (2.62) we can find appropriate bound on  $||F_1||_{L_2(0,T;L_2(\mathbb{T}^3))}$ . Namely,

$$F_1 \sim \mathbf{v} \nabla^2 \varrho + \nabla \mathbf{v} \nabla \varrho + \varrho \nabla^2 \mathbf{v} + \nabla^2 \mathfrak{a} \nabla \varrho + \nabla \mathfrak{a} \nabla^2 \varrho + \nabla \mathfrak{b}.$$

We have

$$\begin{split} \|\mathbf{v}\nabla^{2}\varrho\|_{L_{2}(0,T,L_{2}(\mathbb{T}^{3}))} &\leq T\|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\nabla\mathbf{v}\nabla\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\nabla\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\varrho\nabla^{2}\mathbf{v}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\mathbf{v}\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\nabla\varrho\nabla^{2}\mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq \|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\nabla^{2}\varrho\nabla\mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\nabla\mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}, \end{split}$$

which gives

$$\begin{split} \|F_1\|_{L_2(0,T;L_2(\mathbb{T}^3))} &\leq \|\varrho\|_{L_\infty(0,T;H^2(\mathbb{T}^3))} \Big(T[\|\mathbf{v}\|_{L_\infty(0,T;H^2(\mathbb{T}^3))} + \|\nabla\mathfrak{a}\|_{L_\infty((0,T)\times\mathbb{T}^3)}] \\ &+ \|\nabla^2\mathfrak{a}\|_{L_2(0,T;L_2(\mathbb{T}^3))} + \|\mathfrak{b}\|_{L_2(0,T;H^1(\mathbb{T}^3))}\Big), \end{split}$$

which together with (2.63) implies (2.57).

Next, let  $\alpha$  be any multi-index with  $|\alpha| = 2$ . Applying  $D^{\alpha}$  to (2.61), we obtain

$$(D^{\alpha}\varrho)_{t} - \operatorname{div}\left(\mathfrak{a}\nabla D^{\alpha}\varrho\right) = F_{2}^{\alpha}, \quad D^{\alpha}\varrho|_{t=0} = D^{\alpha}\varrho_{0}, \tag{2.64}$$

where

$$F_2^\alpha \sim \mathbf{v} \nabla^3 \varrho + \nabla \mathbf{v} \nabla^2 \varrho + \nabla^2 \mathbf{v} \nabla \varrho + \varrho \nabla^3 \mathbf{v} + \sum_{k=1}^3 \nabla^k \mathfrak{a} \nabla^{4-k} \varrho + \nabla^2 \mathfrak{b}.$$

In order to prove (2.58), we have to estimate the  $L_2(0,T;L_2(\mathbb{T}^3))$  norm of  $F_2^{\alpha}$ . We have

$$\begin{split} \|\mathbf{v}\nabla^{3}\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{3}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\nabla\mathbf{v}\nabla^{2}\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\nabla\mathbf{v}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))},\\ \|\nabla^{2}\mathbf{v}\nabla\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\nabla^{2}\mathbf{v}\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))}\|\nabla\varrho\|_{L_{\infty}(0,T;H_{2}(\mathbb{T}^{3}))},\\ \|\varrho\nabla^{3}\mathbf{v}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq T\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{3}\mathbf{v}\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}. \end{split}$$

Combining these estimates with Sobolev imbedding, we obtain

$$\|\mathbf{v}\nabla^{3}\varrho + \nabla\mathbf{v}\nabla^{2}\varrho + \nabla^{2}\mathbf{v}\nabla\varrho + \varrho\nabla^{3}\mathbf{v}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \leq T\|\varrho\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))}\|\mathbf{v}\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))}.$$
(2.65)

The terms with  ${\mathfrak a}$  can be treated as follows

18

$$\begin{split} \|\nabla \varrho \nabla^{3} \mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} &\leq \|\nabla \varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{2} \|\nabla^{3} \mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2}, \\ \|\nabla^{2} \varrho \nabla^{2} \mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} &\leq \|\nabla^{2} \varrho\|_{L_{\infty}(0,T;L_{6}(\mathbb{T}^{3}))}^{2} \|\nabla^{2} \mathfrak{a}\|_{L_{2}(0,T;L_{3}(\mathbb{T}^{3}))}^{2}, \\ \|\nabla^{3} \varrho \nabla \mathfrak{a}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} &\leq T \|\nabla \mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{2} \|\nabla^{3} \varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}, \end{split}$$

which together with Sobolev imbedding yields

$$\|\sum_{k=1}^{3} \nabla^{m} \mathfrak{a} \nabla^{4-m} \varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}$$

$$\leq \|\varrho\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))} [T\|\nabla \mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} + \|\nabla^{2} \mathfrak{a}\|_{L_{2}(0,T;H^{1}(\mathbb{T}^{3}))}].$$
(2.66)

Combining (2.65) and (2.66) with the maximal regularity estimate for (2.64), we obtain

$$\begin{aligned} \|\partial_t \nabla^2 \varrho\| + \|\nabla^2 \varrho\|_{L_{\infty}(0,T;H^1(\mathbb{T}^3))} + \|\nabla^2 \varrho\|_{L_2(0,T;H^2(\mathbb{T}^3))} &\leq C \Big[ \|\varrho_0\|_{H^3(\mathbb{T}^3)} + \|\mathfrak{b}\|_{L_2(0,T;H^2(\mathbb{T}^3))} \\ &+ \|\varrho\|_{L_{\infty}(0,T;H^3(\mathbb{T}^3))} \Big( T \big( \|\mathbf{v}\|_{L_{\infty}(0,T;H^3(\mathbb{T}^3))} + \|\nabla\mathfrak{a}\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \big) + \|\nabla^2\mathfrak{a}\|_{L_2(0,T;H^1(\mathbb{T}^3))} \Big) \Big]. \end{aligned}$$

Combining this estimate with (2.62) and (2.63), we obtain (2.58).

It remains to prove (2.60) under additional assumption (2.59). Consider a function  $\psi : \mathbb{R} \to \mathbb{R}$  as

$$\psi(\lambda) = \begin{cases} \frac{1}{2}\lambda^2, & \lambda \le 0\\ 0, & \lambda > 0. \end{cases}$$

Clearly,  $\psi \in C^1(\mathbb{R})$  with

$$\psi^{'}(\lambda) = egin{cases} \lambda, & \lambda \leq 0 \\ 0, & \lambda > 0. \end{cases}$$

Let  $t \in [0,T]$ , define  $K = \underline{\rho_0} := \min_{x \in \mathbb{T}^3} \rho_0 > 0$ ,  $M := \|\operatorname{div} (\mathbf{v} - \mathbf{b})(t)\|_{L_{\infty}(\mathbb{T}^3)}$ . Now, we consider a function  $\overline{h_z} : [0,T) \to \mathbb{R}$  as

$$h_z(t) = \int_{\mathbb{T}^3} \psi\left(\varrho(t,x) - \widetilde{K}\right) \, \mathrm{d}x,$$

where  $\widetilde{K} = K \exp(-Mt)$ . Since,  $\varrho \in \mathcal{V}_3(T)$ , a straightforward computation yields

$$h'_{z}(t) = \int_{\mathbb{T}^{3}} \psi'(\varrho(t,x) - \widetilde{K}) \partial_{t} \varrho(t,x) \, \mathrm{d}x + \int_{\mathbb{T}^{3}} \psi'(\varrho(t,x) - \widetilde{K}) M \widetilde{K}.$$

Using Eq. (2.56), assumption (2.59), and integration by parts, we obtain

$$\begin{split} h_{z}'(t) &= \int_{\mathbb{T}^{3}} \psi''(\varrho(t,x) - \widetilde{K})\varrho\left(\mathbf{v} - \mathbf{b}\right) \nabla(\varrho(t,x) - \widetilde{K}) \, \mathrm{d}x \\ &- \int_{\mathbb{T}^{3}} \mathfrak{a} \psi''(\varrho(t,x) - \widetilde{K}) |\nabla_{x} \varrho(t,x)|^{2} \, \mathrm{d}x \\ &+ \int_{\mathbb{T}^{3}} \psi'(\varrho(t,x) - \widetilde{K}) M \widetilde{K} \mathrm{d}x \\ &= \int_{\mathbb{T}^{3}} \psi''(\varrho(t,x) - \widetilde{K}) \left(\varrho - \widetilde{K}\right) (\mathbf{v} - \mathbf{b}) \nabla(\varrho(t,x) - \widetilde{K}) \, \mathrm{d}x \\ &- \int_{\mathbb{T}^{3}} \mathfrak{a} \psi''(\varrho(t,x) - \widetilde{K}) |\nabla_{x} \varrho(t,x)|^{2} \, \mathrm{d}x \\ &+ \int_{\mathbb{T}^{3}} \psi'(\varrho(t,x) - \widetilde{K}) \widetilde{K} \left(M + \operatorname{div} \left(\mathbf{v} - \mathbf{b}\right)\right) \mathrm{d}x. \end{split}$$

From the assumption on  $\mathfrak{a}$ , we get

$$\int_{\mathbb{T}^3} \mathfrak{a} \psi''(\varrho(t,x) - \widetilde{K}) |\nabla(\varrho(t,x))|^2 \, \mathrm{d} x \ge 0.$$

Moreover, the choice of  $\psi$  and M gives

$$\int_{\mathbb{T}^3} \psi'(\varrho(t,x) - \widetilde{K}) \widetilde{K} \left( M + \operatorname{div} \left( \mathbf{v} - \mathbf{b} \right) \right) \mathrm{d}x \le 0.$$

Now the identity  $\lambda \psi''(\lambda) = \psi'(\lambda)$  for  $\lambda \in \mathbb{R}$  implies

$$\begin{split} h'_{z}(t) &\leq \int_{\mathbb{T}^{3}} \psi'(\varrho(t,x) - \widetilde{K}) \left(\mathbf{v} - \mathbf{b}\right) \nabla(\varrho(t,x) - \widetilde{K}) \, \mathrm{d}x \\ &= \int_{\mathbb{T}^{3}} \nabla(\psi(\varrho(t,x) - \widetilde{K})) \left(\mathbf{v} - \mathbf{b}\right) \, \mathrm{d}x \\ &= -\int_{\mathbb{T}^{3}} \psi(\varrho(t,x) - \widetilde{K}) \, \mathrm{div} \left(\mathbf{v} - \mathbf{b}\right) \, \mathrm{d}x \\ &\leq \| \operatorname{div} \left(\mathbf{v} - \mathbf{b}\right) \|_{L_{\infty}(\mathbb{T}^{3})} h_{z}(t). \end{split}$$

Here, we apply Grönwall's inequality along with  $\inf_{x\in\mathbb{T}^3} \varrho_0 = K$  to deduce

$$h_z(t) = 0$$
 for a.e.  $t \in (0, T)$ .

Therefore, we have  $\varrho(t) \geq \widetilde{K} = \exp\left(-Mt\right) \inf_{x \in \mathbb{T}^3} \varrho_0 > 0$  in  $(0,T) \times \mathbb{T}^3$ .

Next, we state another lemma, related to the maximum and minimal principles of (2.56) with b = 0.

LEMMA 2.4. Assume  $\varrho_0 \in H^2(\mathbb{T}^3)$ ,  $\boldsymbol{v} \in L_{\infty}(0,T;H^2(\mathbb{T}^3))$ , and  $\mathfrak{a} \in L_{\infty}(0,T;W^1_{\infty}) \cap L_2(0,T;H^2(\mathbb{T}^3))$  with  $\mathfrak{a} \geq c > 0$ . Then the unique solution of (2.56) with  $\mathfrak{b} = 0$ 

admits satisfying

$$\inf_{\mathbb{T}^3} \varrho_0 \exp\left(-\int_0^t \|\operatorname{div} \mathbf{v}(s)\|_{L_{\infty}(\mathbb{T}^3)} \, ds\right) \le \varrho(t, x) \le \sup_{\mathbb{T}^3} \varrho_0 \exp\left(\int_0^t \|\operatorname{div} \mathbf{v}(s)\|_{L_{\infty}(\mathbb{T}^3)} \, ds\right),\tag{2.68}$$

for  $0 \leq t \leq T$ .

*Proof. Sketch of proof.* The proof is similar to the proof non-negativity property in lemma 2.3 and extensively used in literatures like Novotný and Straskraba [30, Proposition 7.39, Feireisl and Novotny [19, Lemma 3.1], where they have consider a special case with **a** is constant. For the sake of completeness, we just highlight the key steps:

- Define : R = (sup<sub>T<sup>3</sup></sub> ρ<sub>0</sub>) exp (∫<sub>0</sub><sup>t</sup> || div **v**(s)||<sub>L∞(T<sup>3</sup>)</sub> ds).
  Then R satisfies R'(t) || div **v**(t)||<sub>L∞(T<sup>3</sup>)</sub>R(t) = 0 with R(0) = sup<sub>T<sup>3</sup></sub> ρ<sub>0</sub>.
- Now consider  $W(t, x) = \varrho(t, x) R(t)$  and it satisfies

$$\partial_t W + \operatorname{div}(W\mathbf{v}) - \operatorname{div}(\mathfrak{a}\nabla\varrho) \le 0$$
a.e. in  $(0,T) \times \mathbb{R}^d$ , (2.69)

with  $W(0, x) = \varrho_0 - \sup_{\mathbb{T}^3} \varrho_0 \leq 0.$ 

• Now we test Eq. (2.69) with  $\Psi'(W)$  and integrating over space to obtain

$$\frac{d}{dt} \int_{\mathbb{T}^3} \Psi(W) + 2 \int_{\mathbb{T}^3} \mathfrak{a} \Psi''(W) |\nabla W|^2 \le \|\operatorname{div} \mathbf{v}\|_{L_{\infty}(\mathbb{T}^3)} \Psi(W)$$

where

$$\Psi(\lambda) = \begin{cases} \frac{1}{2}\lambda^2, & \lambda \ge 0\\ 0, & \lambda < 0. \end{cases}$$

• Since,  $\Psi(0) \leq 0$ , from Grönwall's inequality, we conclude

$$\varrho(t,x) \leq \sup_{\mathbb{T}^3} \varrho_0 \exp\left(\int_0^t \|\operatorname{div} \mathbf{v}(s)\|_{L_{\infty}(\mathbb{T}^3)} \, ds\right).$$

For the other side of the inequality (2.68), we need to consider

$$r(t) = \inf_{\mathbb{T}^3} \varrho_0 \exp\left(-\int_0^t \|\operatorname{div} \mathbf{v}(s)\|_{L_{\infty}(\mathbb{T}^3)} \, ds\right)$$

and proceed analogously by considering

$$w(t, x) = \varrho(t, x) - r(t)$$

We have collected all the necessary tools to use in order to prove the convergence of iterative scheme (1.11).

### 3. Convergence of the iterative scheme

#### 3.1. Boundedness of the sequence of approximations

The estimates from the previous section allow us to show

LEMMA 3.1. Let  $(\mathbf{w}^n, \varrho^n)$  be the sequence defined in (1.11) with  $\mathbf{u}^1(t, x) = \mathbf{u}_0(x)$ for  $(t, x) \in (0, T) \times \mathbb{T}^3$  and initial data satisfying (1.4) or (1.5). There exists  $M = M(\|\varrho_0\|_{H^4(\mathbb{T}^3)}, \|\mathbf{u}_0\|_{H^3(\mathbb{T}^3)})$  in case of (1.4) or  $M = M(\|\varrho_0\|_{H^3(\mathbb{T}^3)}, \|\mathbf{u}_0 + \nabla p(\varrho_0)\|_{H^3(\mathbb{T}^3)})$  in case of (1.5) and T = T(M) > 0 such that

$$\|\boldsymbol{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{Y}_3(T)} \le M \quad \forall n \in \mathbb{N},$$
(3.1)

where the spaces  $\mathcal{Y}_k(T)$  and  $\mathcal{V}_k(T)$  are defined in (1.10).

*Proof.* Recall that  $\phi(\cdot)$  denotes an increasing, positive function, precise form of which may vary from line to line. For the purpose of this proof, we also introduce more precise notation  $\phi_1, \phi_2, \phi_3$  to denote given, increasing, positive functions. The first equation of system (1.11) is exactly (2.56) with  $b = 0, \rho = \rho^{n+1}, \mathbf{v} = \mathbf{w}^{n+1}$ , and  $\mathbf{a} = \rho^n p'(\rho^n)$ , while the subequation of the second line of (1.11) corresponding to the *i*th component of  $\mathbf{w}^{n+1}$  is nothing but (2.1) with  $g = 0, \eta = \mathbf{w}_i^{n+1}$ , and  $\mathbf{v} = \mathbf{u}^n$ .

By lemma 2.2, recalling that  $\mathbf{u}^n = \mathbf{w}^n + \nabla p(\varrho^n)$  we have

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_{3}(T)} \le \phi(\sqrt{T}\|\mathbf{u}^{n}\|_{\mathcal{X}_{3}(T)}) = \phi_{1}\left(\sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)})\right).$$
(3.2)

In order to use (2.58) to estimate  $\|\varrho^{n+1}\|_{\mathcal{V}_3(T)}$ , we need to have a closer look at  $\mathfrak{a}(\varrho) = \varrho p'(\varrho)$ 

$$\begin{split} \nabla \mathfrak{a}(\varrho) &= \left[ p'(\varrho) + \varrho p^{(2)}(\varrho) \right] \nabla \varrho, \\ \nabla^2 \mathfrak{a}(\varrho) &= \left[ p'(\varrho) + \varrho p^{(2)}(\varrho) \right] \nabla^2 \varrho + \left[ 2p^{(2)}(\varrho) + \varrho p^{(3)}(\varrho) \right] \nabla \varrho \otimes \nabla \varrho, \\ \nabla^3 \mathfrak{a}(\varrho) &\sim Q(\varrho, p'(\varrho), p^{(2)}(\varrho), p^{(3)}(\varrho), p^{(4)}(\varrho)) [\nabla^3 \varrho + \nabla^2 \varrho \nabla \varrho + |\nabla \varrho|^3], \end{split}$$

where Q is some polynomial, a precise form of which is not relevant. We have  $\rho \in L_{\infty}((0,T) \times \mathbb{T}^3)$ . Therefore, as  $p \in C^5$ ,

$$\begin{split} \|\nabla^{2}\mathfrak{a}(\varrho)\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} &\leq \phi(\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}) \left[\|\nabla^{2}\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} + \||\nabla\varrho|^{2}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2}\right] \\ &\leq T\phi(\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}) \left(\|\nabla^{2}\varrho\|_{L_{\infty}(0,T;H^{1}(\mathbb{T}^{3}))}^{2} + C\|\nabla\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{2}\right) \end{split}$$

and

$$\begin{split} \|\nabla^{3}\mathfrak{a}(\varrho)\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} \\ &\leq \phi(\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})})\left(\|\nabla^{3}\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} + \|\nabla^{2}\varrho\nabla\varrho\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} + \||\nabla\varrho|^{3}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2}\right) \\ &\leq T\phi(\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}) \\ &\times \left(\|\nabla^{3}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} + \|\nabla\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}^{2}\|\nabla^{2}\varrho\|_{L_{\infty}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} + C\||\nabla\varrho|^{3}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\right). \end{split}$$

These bounds with an obvious estimate on  $\|\nabla \mathfrak{a}(\varrho)\|_{L_2(0,T;L_2(\mathbb{T}^3))}$  imply

 $\|\nabla \mathfrak{a}(\varrho)\|_{L_2(0,T;H^2(\mathbb{T}^3))} \le T\phi(\|\varrho\|_{L_\infty((0,T)\times\mathbb{T}^3)})\|\varrho\|_{L_\infty(0,T;H^3(\mathbb{T}^3))}.$  (3.3)

Obviously, we also have

$$\|\nabla \mathfrak{a}(\varrho)\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \leq \phi(\|\varrho\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})})\|\varrho\|_{L_{\infty}(0,T;H^{3}(\mathbb{T}^{3}))}.$$
(3.4)

Applying (2.58) to the first equation of (1.11), we obtain

$$\begin{split} \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} &\leq \phi(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}) \Big[ C(T) + \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} \\ &\times [T(\|\mathbf{w}^{n+1}\|_{\mathcal{V}_{3}(T)} + \|\nabla \mathfrak{a}(\varrho^{n})\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}) + \|\nabla \mathfrak{a}(\varrho^{n})\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))}] \Big], \end{split}$$

which together with (3.3), (3.4), and (3.2) gives

$$\begin{aligned} \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} &\leq \phi_{2}(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}) \left[ C(T) + T\|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} \left( \|\mathbf{w}^{n+1}\|_{\mathcal{Y}_{3}(T)} + \phi_{3}(\|\varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})})\|\varrho^{n}\|_{\mathcal{V}_{3}(T)} \right) \right] \\ &\leq \phi_{2}(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}) \left[ C(T) + T\|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} \left[ \phi_{1} \left( \sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)}) \right) \\ &+ \phi_{3}(\|\varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})})\|\varrho^{n}\|_{\mathcal{V}_{3}(T)} \right] \right]. \end{aligned}$$

$$(3.5)$$

Let us take

$$M = 2 \max\{\sup_{s \in [0,1]} \phi_1(s), \phi_2(\|\varrho_0\|_{H^3(\mathbb{T}^3)})C(T)\},\$$

where C(T) is the constant from (3.5). Then, assuming that

$$\|\mathbf{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{V}_3(T)} \le M,$$

for sufficiently small T we can assure that

$$\begin{split} \phi_1(\sqrt{T}(\|\mathbf{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{V}_3(T)}) &\leq \frac{M}{2} \\ T\phi_2(\|\varrho_0\|_{H^3}) \left[\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)} + \phi_3(\|\varrho^n\|_{L_{\infty}((0,T)\times\mathbb{T}^3)})\|\varrho^n\|_{\mathcal{V}_3(T)}\right] &\leq \frac{1}{2}, \end{split}$$

which together with (3.2) and (3.5) implies

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)} + \|\varrho^{n+1}\|_{\mathcal{V}_3(T)} \le M,$$

thus we have (3.1).

## 3.2. Contraction argument. Proof of theorem 1.1.

LEMMA 3.2. Let us denote

$$\delta \mathbf{w}^n = \mathbf{w}^n - \mathbf{w}^{n-1}, \quad \delta \varrho^n = \varrho^n - \varrho^{n-1}, \quad \delta \mathbf{u}^n = \mathbf{u}^n - \mathbf{u}^{n-1}.$$

Then, under the assumptions of lemma 3.1, we have

$$\|\delta \boldsymbol{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n+1}\|_{\mathcal{V}_{2}(T)} \leq E(T) \left(\|\delta \boldsymbol{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n}\|_{\mathcal{V}_{2}(T)}\right),$$
(3.6)

where notation E(t) is described in §1.1.

22

*Proof.* Subtracting (1.11) for  $(\mathbf{w}^{n+1}, \varrho^{n+1})$  and  $(\mathbf{w}^n, \varrho^n)$ , we obtain

$$\begin{cases} \delta \varrho_t^{n+1} + \operatorname{div} \left( \delta \varrho^{n+1} \mathbf{w}^{n+1} \right) - \operatorname{div} \left( \varrho^n p'(\varrho^n) \nabla \delta \varrho^{n+1} \right) = R_n, \\ \delta \mathbf{w}_t^{n+1} + \mathbf{u}^n \cdot \nabla \delta \mathbf{w}^{n+1} = -\delta \mathbf{u}^n \cdot \nabla \mathbf{w}^n, \\ \left( \delta \varrho^{n+1}, \delta \mathbf{w}^{n+1} \right) |_{t=0} = (0, \mathbf{0}), \end{cases}$$
(3.7)

where

$$R_n = \operatorname{div}\left[\left(p'(\varrho^{n-1})\delta\varrho^n + \varrho^n(p'(\varrho^n) - p'(\varrho^{n-1}))\right)\nabla\varrho^n - \varrho^n\delta\mathbf{w}^{n+1}\right].$$
 (3.8)

Each equation of the second line of (3.7) corresponds to (2.1) with  $g \sim \delta \mathbf{u}^n \nabla \mathbf{w}^n$ . Therefore, taking into account (3.2), we can differentiate the right hand side in space only twice. For this purpose, we show contraction in lower regularity then the estimate (3.2). This approach is well known in the regularity theory of the compressible and inhomogeneous Navier–Stokes systems to overcome the limitations coming from the presence of the gradient of the density in the continuity equation (see among others [21], [27], [32], [14], [15], [16]). Combining (2.18) and (2.22) for  $\eta_0 = 0$ , we obtain

$$\|\eta\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \leq E(T)\|g\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))},$$

which applied to (3.7) implies

$$\|\delta \mathbf{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \leq E(T) \|\delta \mathbf{u}^{n} \cdot \nabla \mathbf{w}^{n}\|_{L_{q}(0,T;H^{2}(\mathbb{T}^{3}))} \quad \forall, \ 1 < q \leq \infty.$$
(3.9)

The first equation of (3.7) is (2.56) with  $\mathbf{v} = \mathbf{w}^{n+1}$ ,  $\mathfrak{a} = \varrho^n p'(\varrho^n)$ ,  $b = R_n$ , and  $\varrho_0 = 0$ .

Therefore, (2.57) implies

$$\begin{split} \|\delta\varrho^{n+1}\|_{\mathcal{V}_{2}(T)} &\leq C \Big[ \|R_{n}\|_{L_{2}(0,T;H^{1}(\mathbb{T}^{3}))} + \|\delta\varrho^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \\ &\times \Big( T \Big( \|\mathbf{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\nabla\big(\varrho^{n}p'(\varrho^{n})\big)\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \Big) \\ &+ \|\nabla^{2}\big(\varrho^{n}p'(\varrho^{n})\big)\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \Big] \Big]. \end{split}$$
(3.10)

As  $p(\cdot)$  is sufficiently smooth, we have

$$\nabla^2(\varrho^n p'(\varrho^n)) \sim |\nabla \varrho^n|^2 + |\varrho^n| \, |\nabla^2 \varrho^n|.$$

By (3.1), we have

$$\||\nabla \varrho^{n}|^{2}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}^{2} = \int_{0}^{T} \|\nabla \varrho^{n}\|_{L_{4}(\mathbb{T}^{3})}^{4} \leq T \|\nabla \varrho^{n}\|_{L_{\infty}(0,T;L_{4}(\mathbb{T}^{3}))} \leq M E(T),$$

and

$$\|\varrho^n \nabla^2 \varrho^n\|_{L_2(0,T;L_2(\mathbb{T}^3))} \le \|\varrho^n\|_{L_\infty((0,T)\times\mathbb{T}^3)} \sqrt{T} \|\nabla^2 \varrho^n\|_{L_\infty(0,T;L_2(\mathbb{T}^3))} \le M E(T).$$

Therefore choosing T sufficiently small, we can ensure that

2nd line of 
$$(3.10) \le \frac{1}{2}$$
.

Choosing such T and adding (3.10) to (3.9), we obtain

$$\begin{aligned} \|\delta \mathbf{w}^{n+1}\|_{L_{\infty}(0,T;W_{2}^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n+1}\|_{\mathcal{V}_{2}(T)} &\leq C \|R_{n}\|_{L_{2}(0,T;H^{1}(\mathbb{T}^{3}))} \\ + E(T)\|\delta \mathbf{u}^{n} \cdot \nabla \mathbf{w}^{n}\|_{L_{q}(0,T;W_{2}^{2}(\mathbb{T}^{3}))}. \end{aligned}$$
(3.11)

Recalling (3.8), we have

$$\begin{aligned} \|R_n\|_{L_2(0,T;H^1(\mathbb{T}^3))} &\leq \|p'(\varrho^{n-1})\delta\varrho^n \nabla \varrho^n\|_{L_2(0,T;H^2(\mathbb{T}^3))} \\ &+ \|\varrho^n(p'(\varrho^n) - p'(\varrho^{n-1}))\nabla \varrho^n\|_{L_2(0,T;H^2(\mathbb{T}^3))} \\ &+ \|\varrho^n \delta \mathbf{w}^{n+1}\|_{L_2(0,T;H^2(\mathbb{T}^3))} =: A_1 + A_2 + A_3. \end{aligned}$$
(3.12)

The last term can be estimated directly:

$$A_{3} \leq \|\varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\delta\mathbf{w}^{n+1}\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))} \leq E(T) \|\delta\mathbf{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))},$$
(3.13)

Let us proceed with  $A_1$ . We estimate the second order derivatives, which are the most restrictive. We have

$$\nabla^2 \left( p'(\varrho^{n-1}) \delta \varrho^n \nabla \varrho^n \right) \sim \nabla^2 \varrho^{n-1} \delta \varrho^n \nabla \varrho^n + \nabla \varrho^{n-1} \delta \varrho^n \nabla^2 \varrho^n + \nabla \varrho^{n-1} \nabla \delta \varrho^n \nabla \varrho^n \\ + \varrho^{n-1} \nabla^2 \delta \varrho^n \nabla \varrho^n + \varrho^{n-1} \nabla \delta \varrho^n \nabla^2 \varrho^n + \varrho^{n-1} \delta \varrho^n \nabla^3 \varrho^n .$$

The first two terms have the same structure and can be bounded as follows applying (3.1):

$$\begin{split} \|\nabla^{2} \varrho^{n-1} \delta \varrho^{n} \nabla \varrho^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq \|\nabla \varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla^{2} \varrho^{n-1}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \|\delta \varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \\ &\leq M^{2} E(T) \|\delta \varrho^{n}\|_{\mathcal{V}_{2}(T)}. \end{split}$$

For the third term, we have

$$\begin{split} \|\nabla \varrho^{n-1} \nabla \delta \varrho^n \nabla \varrho^n \|_{L_2(0,T;L_2(\mathbb{T}^3))} \\ &\leq \|\nabla \varrho^{n-1} \|_{L_\infty((0,T)\times\mathbb{T}^3)} \|\nabla \varrho^n \|_{L_\infty((0,T)\times\mathbb{T}^3)} \|\nabla \delta \varrho^n \|_{L_2(0,T;L_2(\mathbb{T}^3))} \\ &\leq M^2 E(T) \|\delta \varrho^n \|_{\mathcal{V}_2(T)}, \end{split}$$

and the same estimate holds for the fourth one. Next,

$$\begin{split} &\|\varrho^{n-1}\nabla\delta\varrho^{n}\nabla^{2}\varrho^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq \sqrt{T}\|\varrho^{n-1}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{2}\varrho^{n}\|_{L_{\infty}(0,T;L_{4}(\mathbb{T}^{3}))}\|\nabla\delta\varrho^{n}\|_{L_{\infty}(0,T;L_{4}(\mathbb{T}^{3}))} \\ &\leq M^{2}E(T)\|\delta\varrho^{n}\|_{\mathcal{V}_{2}(T)}, \end{split}$$

and finally, for the last term, we have

$$\begin{split} \|\varrho^{n-1}\delta\varrho^n\nabla^3\varrho^n\|_{L_2(0,T;L_2(\mathbb{T}^3))} &\leq C \|\varrho^{n-1}\|_{L_\infty((0,T)\times\mathbb{T}^3)}T\|\nabla^3\varrho^n\|_{L_\infty(0,T;L_2(\mathbb{T}^3))}\|\delta\varrho^n\|_{L_\infty((0,T)\times\mathbb{T}^3)} &\leq M^2E(T)\|\delta\varrho^n\|_{\mathcal{V}_2(T)}. \end{split}$$

Plugging the above estimates into (3.12), we obtain and observing that  $A_1$  and  $A_2$  have the same structure due to assumed regularity of the pressure we obtain

$$A_1 + A_2 \le M^2 E(T) \| \delta \varrho^n \|_{\mathcal{V}_2(T)},$$

which together with (3.13) gives

$$\|R_n\|_{L_2(0,T;H^1(\mathbb{T}^3))} \le E(T) \left( \|\delta \varrho^n\|_{\mathcal{V}_2(T)} + \|\delta \mathbf{w}^{n+1}\|_{L_\infty(0,T;H^2(\mathbb{T}^3))} \right).$$
(3.14)

Now in order to close the contraction argument, it is enough to estimate the second term on the RHS of (3.11). Recalling (1.2), we have

$$\delta \mathbf{u}^n = \delta \mathbf{w}^n - p'(\varrho^n) \nabla \delta \varrho^n - [p'(\varrho^n) - p'(\varrho^{n-1})] \nabla \varrho^{n-1}.$$

We have

$$\begin{split} \|\nabla^{2}(\delta \mathbf{w}^{n} \cdot \nabla \mathbf{w}^{n})\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} &\leq C\left(\|\nabla^{2} \delta \mathbf{w}^{n} \cdot \nabla \mathbf{w}^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &+ \|\nabla \delta \mathbf{w}^{n} \cdot \nabla^{2} \mathbf{w}^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &+ \|\delta \mathbf{w}^{n} \cdot \nabla^{3} \mathbf{w}^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))}\right) =: B_{1} + B_{2} + B_{3}. \end{split}$$

We again apply (3.1) to obtain

$$B_{1} \leq C\sqrt{T} \|\nabla \mathbf{w}_{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \leq ME(T) \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))},$$
  

$$B_{2} \leq C\sqrt{T} \|\mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} \leq ME(T) \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))},$$
  

$$B_{3} \leq C \|\delta \mathbf{w}^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})} \|\nabla^{3} \mathbf{w}^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \leq ME(T) \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))},$$
  
which gives

$$\|\delta \mathbf{w}^{n} \cdot \nabla \mathbf{w}^{n}\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))} \leq E(T) \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))}.$$
(3.16)

Next,

$$\nabla^{2} \left( p'(\varrho^{n}) \nabla \delta \varrho^{n} \right) \sim \nabla^{2} \varrho^{n} \nabla \delta \varrho^{n} \nabla \mathbf{w}^{n} + \nabla \varrho^{n} \nabla^{2} \delta \varrho^{n} + \nabla \varrho^{n} \nabla \delta \varrho^{n} \nabla^{2} \mathbf{w}^{n} + \varrho^{n} \nabla^{3} \delta \varrho^{n} \nabla \mathbf{w}^{n} + \varrho^{n} \nabla^{2} \delta \varrho^{n} \nabla^{2} \mathbf{w}^{n} + \varrho^{n} \nabla^{3} \delta \varrho^{n} \nabla \mathbf{w}^{n}.$$
(3.17)

Let us focus on the last term:

$$\begin{split} &\|\varrho^{n}\nabla^{3}\delta\varrho^{n}\cdot\nabla\mathbf{w}^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq C(T)\|\varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla\mathbf{w}^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}\|\nabla^{3}\delta\varrho^{n}\|_{L_{2}(0,T;L_{2}(\mathbb{T}^{3}))} \\ &\leq CM^{2}\|\delta\varrho^{n}\|_{L_{q}(0,T;H^{3}(\mathbb{T}^{3}))} \leq CM^{2}\|\delta\varrho^{n}\|_{\mathcal{V}_{2}(T)}. \end{split}$$

The lack of small constant is not a problem, since the term which we are estimating is already multiplied by a small constant in (3.11). The other terms in (3.17)

can be estimated similarly, here we even get additional smallness in time. Summing up, we obtain

$$\|p'(\varrho^n)\nabla\delta\varrho^n\cdot\nabla\mathbf{w}^n\|_{L_2(0,T;H^2(\mathbb{T}^3))} \le C(M,T)\|\delta\varrho^n\|_{\mathcal{V}_2(T)}.$$
(3.18)

Similarly, we can show

$$\left\| [p'(\varrho^n) - p'(\varrho^{n-1})] \nabla \varrho^{n-1} \cdot \mathbf{w}^n \right\|_{L_2(0,T;H^2(\mathbb{T}^3))} \le E(T) \| \delta \varrho^n \|_{\mathcal{V}_2(T)}.$$
(3.19)

Combining (3.16), (3.18), and (3.19), we get

$$\|\delta \mathbf{u}^{n} \cdot \nabla \mathbf{w}^{n}\|_{L_{2}(0,T;H^{2}(\mathbb{T}^{3}))} \leq E(T) \|\delta \mathbf{w}^{n}\|_{L_{\infty}(0,T;H^{2})} + C \|\delta \varrho^{n}\|_{\mathcal{V}_{2}(T)}.$$
 (3.20)

Plugging (3.14) and (3.20) into (3.11), we finally conclude (3.6).

Now we complete the proof of theorem 1.1 in a standard way. Inequality (3.6) implies

$$(\mathbf{w}^n, \varrho^n) \to (\mathbf{w}, \varrho)$$
 strongly in  $L_{\infty}(0, T; H^2(\mathbb{T}^3)) \times \mathcal{V}_2(T)$ .

On the other hand, the estimate (3.1) implies existence of a subsequence, which we can still denote  $(\mathbf{w}^n, \varrho^n)$ 

$$(\mathbf{w}^n, \varrho^n) \to (\mathbf{w}, \varrho)$$
 weakly in  $\mathcal{Y}_3(T) \times \mathcal{V}_3(T)$ .

Setting  $\mathbf{u} = \mathbf{w} - \nabla p(\varrho)$  we easily verify that the limit satisfies (1.7).

## 4. Existence theory for general velocity offsets

The goal of this section is to extend the previous result to the case of different velocity offset functions  $p(\varrho)$ .

• Singular offset: As mentioned in the introduction, Aceves et al. [1] derive a variant of dissipative Aw-Rascle system from a microscopic model of pedestrian dynamics. This corresponds to (1.7) with the offset function in the form (1.9). Drawing motivation from this, we consider a more general offset function and the system (1.6), namely

$$p(\varrho) = a \frac{\varrho^{\alpha}}{(1-\varrho)^{\beta}} \text{ with } a > 0, \ \alpha > 0 \text{ and } \beta > 1.$$
(4.1)

Similar form was considered in [12] where the existence of regular solutions for certain approximation of this function was proven in one space dimension. The authors also performed the singular limit passage  $a \rightarrow 0$  obtaining in the limit the hard congestion system. Similar limit has been postulated in the multi-dimensional case, see [1], but to our knowledge it has not yet been proved rigorously. Nevertheless, it is expected that in the hard congestion limit, in the saturated region where  $\rho = 1$ , one cannot expect a regular solution due to appearance of extra forcing term. Therefore here we keep a a positive constant, which results in restriction to unsaturated flow provided that the initial density is strictly below 1. This case is discussed in §4.1. • Non-local velocity offset: It turns out that considering the non-local offset function

$$p(\varrho) = K * \varrho, \tag{4.2}$$

where K is some non-local kernel, leads to reformulation of the system that generalizes the pressureless Euler-alignment model, see [13] for further details. Unfortunately, we are not able to apply our construction along with the linear theory (described in §2 and 3) directly to the system (1.7) with (4.2). Instead, we consider a system where the offset function  $p(\varrho)$  is a combination of local and a particular non-local component that corresponds to the Newtonian potential, i.e.,

$$-\Delta K * \varrho \approx \varrho.$$

More precisely, the closure relation is now of the form

$$\mathbf{w} = \mathbf{u} + \nabla p(\varrho) + \nabla \Phi_{\varrho} \text{ with } p \in C^{5}(\mathbb{R}_{+}) \text{ and } -\Delta \Phi_{\varrho} = \varrho - \langle \varrho \rangle,$$

where  $\langle \varrho \rangle = \frac{1}{|\mathbb{T}^3|} \int_{\mathbb{T}^3} \varrho \, \mathrm{d}x$ , and so (1.7) can be rewritten as follows

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{w}) - \operatorname{div}(\rho (\nabla p(\rho) + \nabla \Phi_\rho)) = 0, \qquad (4.4a)$$

$$\left\{ \partial_t(\rho \mathbf{w}) + \operatorname{div}(\rho \mathbf{w} \otimes \mathbf{w}) = \operatorname{div}(\rho \mathbf{w} \otimes \nabla p(\rho)) + \operatorname{div}(\rho \mathbf{w} \otimes \nabla \Phi_\rho), \quad (4.4b) \right\}$$

$$-\Delta\Phi_{\rho} = \rho - \langle \rho \rangle . \tag{4.4c}$$

This case is discussed in  $\S4.2$ .

#### 4.1. Construction for the system with singular velocity offset

Here, we consider (1.6a)–(1.6b) with (4.1). Moreover, along with the hypothesis on initial data (1.4) or (1.5), we need

$$0 < \varrho_0(x) < 1 \text{ for } x \in \mathbb{T}^3.$$

$$(4.5)$$

THEOREM 4.1. Assume the initial data satisfies (1.4) or (1.5) with (4.5). Then there exists T > 0 such that system (1.7) admits a unique solution  $(\varrho, \boldsymbol{w}) \in \mathcal{V}_3(T) \times \mathcal{Y}_3(T)$  with the estimate

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\boldsymbol{w}\|_{\mathcal{Y}_{3}(T)} \leq C(\|\varrho_{0}\|_{H^{4}(\mathbb{T}^{3})}, \|\boldsymbol{u}_{0}\|_{H^{3}(\mathbb{T}^{3})}) and \ 0 < \varrho < 1 in \ [0, T) \times \mathbb{T}^{3}$$

in case of (1.4) or

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\boldsymbol{w}\|_{\mathcal{Y}_{3}(T)} \leq C(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}, \|\boldsymbol{u}_{0} + \nabla p(\varrho_{0})\|_{H^{3}(\mathbb{T}^{3})}) and \ 0 < \varrho < 1 in \ [0, T) \times \mathbb{T}^{3}$$
  
in case of (1.5).

The proof of the theorem 4.1 consists of two parts: the proof of boundedness of approximate solutions and the compactness argument. The iteration scheme for this case is the same as in (1.11), but the proof of analogue of lemma 3.1 requires some alterations that we explain below. The contraction argument is similar to the one described in §3.2, more precisely lemma 3.2, and we skip the details here.

The main lemma for the uniform bounds reads:

LEMMA 4.2. Let  $(\mathbf{w}^n, \varrho^n)$  be the sequence defined in (1.11) with  $\mathbf{u}^1(t, x) = \mathbf{u}_0(x)$ for  $(t, x) \in (0, T) \times \mathbb{T}^3$ . There exists  $M = M(\|\varrho_0\|_{H^4(\mathbb{T}^3)}, \|\mathbf{u}_0\|_{H^3(\mathbb{T}^3)})$  in case of (1.4) or  $M = M(\|\varrho_0\|_{H^3(\mathbb{T}^3)}, \|\mathbf{u}_0 + \nabla p(\varrho_0)\|_{H^3(\mathbb{T}^3)})$  in case of (1.5) and T(=T(M)) > 0such that

$$\|\boldsymbol{w}^n\|_{\mathcal{Y}_3(T)} + \|\boldsymbol{\varrho}^n\|_{\mathcal{V}_3(T)} \le M \text{ and } 0 < \boldsymbol{\varrho}_n(t,x) < 1 \ \forall n \in \mathbb{N}, \ \forall (t,x) \in (0,T) \times \mathbb{T}^3,$$

$$(4.6)$$

where the spaces  $\mathcal{Y}_k(T)$  and  $\mathcal{V}_k(T)$  are defined in (1.10). More precisely, it holds that there exists  $0 < \vartheta < 1$  such that

$$\frac{\vartheta}{2} < \varrho_n(t,x) < 1 - \frac{\vartheta}{2} \quad \forall n \in \mathbb{N}, \ \forall (t,x) \in (0,T) \times \mathbb{T}^3.$$
(4.7)

*Proof.* Similarly to lemma (3.1), we prove this lemma with the help of an induction hypothesis: we interpret the first equation in (1.11) as Eq. (2.56) with  $\mathfrak{b} = 0$ ,  $\varrho = \varrho^{n+1}$ ,  $\mathbf{v} = \mathbf{w}^{n+1}$ , and  $\mathfrak{a} = \varrho^n p'(\varrho^n)$ , while the *i*th row of the second equation in (1.11) as (2.1) with g = 0,  $\eta = \mathbf{w}_i^{n+1}$  and  $\mathbf{v} = \mathbf{u}^n$ .

Clearly, there exists  $0 < \vartheta < 1$  such that

$$\vartheta \leq \inf_{\mathbb{T}^3} \varrho_0 \quad \text{and} \quad \sup_{\mathbb{T}^3} \varrho_0 \leq 1 - \vartheta.$$

Now, we assume that  $(\varrho_n, \mathbf{w}_n)$  satisfy (4.6) along with (4.7). This yields,

$$\mathfrak{a} = \varrho^n p'(\varrho^n) > c(\vartheta) > 0$$

and  $\mathfrak{a} \in L_{\infty}(0,T; W^{1}_{\infty}(\mathbb{T}^{3})) \cap L_{2}(0,T; H^{2}(\mathbb{T}^{3}))$ . Here, we perform the similar estimates for  $\mathbf{w}^{n+1}$  and obtain

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_{3}(T)} \le \phi(\sqrt{T}\|\mathbf{u}^{n}\|_{\mathcal{X}_{3}(T)}) = \phi_{1}\left(\sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)})\right).$$
(4.8)

Similarly, for  $\rho^{n+1}$ , we have

$$\begin{aligned} \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} &\leq \phi_{2}(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}) \left[ C(T) + T \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} \left[ \phi_{1} \left( \sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)}) \right) \right. \\ &+ \phi_{3}(\|\varrho^{n}\|_{L_{\infty}((0,T)\times\mathbb{T}^{3})}) \|\varrho^{n}\|_{\mathcal{V}_{3}(T)} \right] \right]. \end{aligned}$$

$$(4.9)$$

Applying lemma 2.4, we obtain

$$\inf_{\mathbb{T}^3} \varrho_0 \exp\left(-\int_0^t \|\operatorname{div} \mathbf{w}^{n+1}(s)\|_{L_{\infty}} ds\right) \\
\leq \varrho^{n+1}(t,x) \leq \sup_{\mathbb{T}^3} \varrho_0 \exp\left(\int_0^t \|\operatorname{div} \mathbf{w}^{n+1}(s)\|_{L_{\infty}} ds\right),$$

for  $t \in (0, T)$ . Because

$$\|\operatorname{div} \mathbf{w}^{n+1}\|_{L_{\infty}((0,T)\times\mathbb{T}^3)} \le C \|\mathbf{w}^{n+1}\|_{\mathcal{V}_3(T)},$$

with C independent of time, we further deduce

$$\vartheta \exp\left(-tC\|\mathbf{w}^{n+1}\|_{\mathcal{V}_3(T)}\right) \le \varrho^{n+1}(t,x) \le (1-\vartheta) \exp\left(tC\|\mathbf{w}^{n+1}\|_{\mathcal{V}_3(T)}\right).$$
(4.10)

We set

$$M = 2 \max\{\sup_{s \in [0,1]} \phi_1(s), \phi_2(\|\varrho_0\|_{H^3(\mathbb{T}^3)})C(T)\},\$$

where C(T) is the constant from (4.9). Assuming that

$$\|\mathbf{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{V}_3(T)} \le M,$$

for sufficiently small T, we can show that

$$\begin{split} \phi_1(\sqrt{T}(\|\mathbf{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{V}_3(T)}) &\leq \frac{M}{2}, \\ T\phi_2(\|\varrho_0\|_{H^3(\mathbb{T}^3)}) \left[\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)} + \phi_3(\|\varrho^n\|_{L_{\infty}((0,T)\times\mathbb{T}^3)})\|\varrho^n\|_{\mathcal{V}_3(T)}\right] &\leq \frac{1}{2}. \end{split}$$

This in turn, along with (4.8) and (4.9), implies

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)} + \|\varrho^{n+1}\|_{\mathcal{V}_3(T)} \le M,$$

thus we have the first part of (4.6). Finally, we choose a sufficiently small T depending only on M, such that from (4.10) we obtain that

$$\frac{\vartheta}{2} < \varrho^{n+1}(t,x) < 1 - \frac{\vartheta}{2} \quad \text{on } (0,T) \times \mathbb{T}^3.$$

This finishes the proof.

#### 4.2. Construction for the system with non-local velocity offset

In this part, we prove the following existence result for system (4.4):

THEOREM 4.3. Assume the initial data satisfies (1.4) or (1.5). Then there exists T > 0 such that system (4.4) admits a unique solution  $(\varrho, \boldsymbol{w}, \Phi_{\varrho}) \in \mathcal{V}_3(T) \times \mathcal{Y}_3(T) \times \mathcal{V}_5(T)$  with the estimate

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\boldsymbol{w}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{\varrho}\|_{\mathcal{V}_{5}(T)} \le C(\|\varrho_{0}\|_{H^{4}(\mathbb{T}^{3})}, \|\boldsymbol{u}_{0}\|_{H^{3}(\mathbb{T}^{3})})$$

in case of (1.4) or

$$\|\varrho\|_{\mathcal{V}_{3}(T)} + \|\boldsymbol{w}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{\varrho}\|_{\mathcal{V}_{5}(T)} \le C(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}, \|\boldsymbol{u}_{0} + \nabla p(\varrho_{0})\|_{H^{3}(\mathbb{T}^{3})})$$

in case of (1.5).

Again the proof of the above theorem is similar to the proof of theorem 1.1.

In this case, however, we use the following iteration scheme for construction of the approximate solution:

$$\begin{cases} \varrho_t^{n+1} + \operatorname{div}\left(\varrho^{n+1}\mathbf{w}^{n+1}\right) = \operatorname{div}\left(\varrho^n p'(\varrho^n)\nabla\varrho^{n+1}\right) + \operatorname{div}\left(\varrho^{n+1}\nabla\Phi_n\right),\\ \mathbf{w}_t^{n+1} + \mathbf{u}^n \cdot \nabla\mathbf{w}^{n+1} = 0,\\ -\Delta\Phi_{n+1} = \varrho^{n+1} - \langle \varrho^{n+1} \rangle,\\ (\varrho^{n+1}, \mathbf{w}^{n+1})|_{t=0} = (\varrho_0, \mathbf{u}_0 + \nabla p(\varrho_0) + \nabla\Phi_{\varrho_0}), \end{cases}$$
(4.11)

where  $-\Delta \Phi_{\varrho_0} = \varrho_0 - \langle \varrho_0 \rangle$ ,  $(\varrho^0, \mathbf{w}^0, \Phi_0) = (\varrho_0, \mathbf{u}_0 + \nabla p(\varrho_0) + \nabla \Phi_{\varrho_0}, \Phi_{\varrho_0})$  and

$$\mathbf{w}^{n} = \mathbf{u}^{n} + \nabla p(\varrho^{n}) + \nabla \Phi_{n}.$$
(4.12)

The key observation is that the equation for  $\Phi_{n+1}$  in this scheme is an elliptic equation. Therefore, the regularity class for  $\{\Phi_n\}$  is higher. Additionally, we have uniform bounds for the approximate sequence  $(\varrho^n, \mathbf{w}^n, \Phi_n)$  in spaces  $\mathcal{V}_3(T) \times \mathcal{V}_3(T) \times \mathcal{V}_5(T)$ , as stated in the following Lemma.

LEMMA 4.4. Let  $(\boldsymbol{w}^n, \varrho^n, \Phi_n)$  be the sequence defined in (4.11) with  $\boldsymbol{u}^1(t, x) = \boldsymbol{u}_0(x)$  for  $(t, x) \in (0, T) \times \mathbb{T}^3$ . There exists  $M = M(\|\varrho_0\|_{H^4(\mathbb{T}^3)}, \|\boldsymbol{u}_0\|_{H^3(\mathbb{T}^3)})$  in case of (1.4) or  $M = M(\|\varrho_0\|_{H^3}, \|\boldsymbol{u}_0 + \nabla p(\varrho_0)\|_{H^3(\mathbb{T}^3)})$  in case of (1.5) and T = T(M) > 0 such that

$$\|\boldsymbol{w}^{n}\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{Y}_{3}(T)} + \|\Phi_{n}\|_{\mathcal{Y}_{5}(T)} \le M \quad \forall n \in \mathbb{N},$$

where the spaces  $\mathcal{Y}_k(T)$  and  $\mathcal{V}_k(T)$  are defined in (1.10).

Proof. Sketch of the proof. The key steps of the proof are the following.

• The bounds on  $\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)}$  and  $\|\varrho^{n+1}\|_{\mathcal{V}_3(T)}$  are obtained by calculations similar to proof of lemma 3.1, but keeping in mind that the relation between  $\mathbf{w}_n$  and  $\mathbf{u}_n$  has been modified (4.12). Ultimately, we obtain

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_{3}(T)} \leq \phi(\sqrt{T}\|\mathbf{u}^{n}\|_{\mathcal{X}_{3}(T)}) = \phi_{1}\left(\sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{n}\|_{\mathcal{V}_{5}(T)})\right),$$
(4.13)

and

$$\begin{split} \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} &\leq \phi_{2}(\|\varrho_{0}\|_{H^{3}(\mathbb{T}^{3})}) \left[ C(T) + T \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} \\ &\left[ \phi_{1} \left( \sqrt{T}(\|\mathbf{w}^{n}(t)\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{n}\|_{\mathcal{V}_{5}(T)}) \right) \\ &+ \phi_{3}(\|\varrho^{n}\|_{L_{\infty}((0,T) \times \mathbb{T}^{3})}) \left( \|\varrho^{n}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{n}\|_{\mathcal{V}_{5}(T)} \right) \right] \right]. \end{split}$$

$$(4.14)$$

https://doi.org/10.1017/prm.2025.10033 Published online by Cambridge University Press

30

• The estimate of  $\|\Phi_{n+1}\|_{\mathcal{V}_5(T)}$  follows from the elliptic regularity estimates for the solutions of the problem

$$-\Delta \Phi_{n+1} = \varrho^{n+1} - \langle \varrho^{n+1} \rangle .$$

Following Evans [18, Chapter 6], we obtain

$$\|\Phi_{n+1}\|_{\mathcal{V}_5(T)} \le C_{\text{ell}} \|\varrho^{n+1}\|_{\mathcal{V}_3(T)}.$$
(4.15)

• The strict positivity of  $\rho^{n+1}$  is a direct consequence of lemma 2.3, where we use the particular form of  $\mathfrak{b} = \operatorname{div}(\rho^n \nabla \Phi_n)$  in the first line of (4.11), i.e.,

$$\varrho_t^{n+1} + \operatorname{div}\left(\varrho^{n+1}\mathbf{w}^{n+1}\right) - \operatorname{div}\left(\varrho^n p'(\varrho^n)\nabla \varrho^{n+1}\right) = \mathfrak{b} = \operatorname{div}\left(\varrho^{n+1}\nabla \Phi_n\right).$$

• Now we combine estimates (4.13)-(4.15) and choose

$$M = 2(1 + C_{\text{ell}}) \max\{\sup_{s \in [0,1]} \phi_1(s), \phi_2(\|\varrho_0\|_{H^3(\mathbb{T}^3)})C(T)\},\$$

where C(T) is the constant from (4.14) and  $C_{\text{ell}}$  is from (4.15). Now, assuming that

$$\|\mathbf{w}^{n}\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{n}\|_{\mathcal{V}_{5}(T)} \le M,$$

for sufficiently small T, first we can show that

$$\begin{split} \phi_1(\sqrt{T}(\|\mathbf{w}^n\|_{\mathcal{Y}_3(T)} + \|\varrho^n\|_{\mathcal{V}_3(T)} + \|\Phi_{n+1}\|_{\mathcal{V}_5(T)}) &\leq \frac{M}{2} \\ T\phi_2(\|\varrho_0\|_{H^3(\mathbb{T}^3)}) \left[\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_3(T)} + \phi_3(\|\varrho^n\|_{L_{\infty}((0,T)\times\mathbb{T}^3)}) \\ \left(\|\varrho^n\|_{\mathcal{V}_3(T)} + \|\Phi_n\|_{\mathcal{V}_5(T)}\right) \right] &\leq \frac{1}{2}. \end{split}$$

This, along with (4.8) and (4.9), implies

$$\|\mathbf{w}^{n+1}\|_{\mathcal{Y}_{3}(T)} + \|\varrho^{n+1}\|_{\mathcal{V}_{3}(T)} + \|\Phi_{n+1}\|_{\mathcal{V}_{5}(T)} \le M,$$

and the proof is complete.

The second part of the proof of theorem 4.3 is to provide an analogue of the contraction argument described in §3.2. We denote

$$\delta \mathbf{w}^n = \mathbf{w}^n - \mathbf{w}^{n-1}, \quad \delta \varrho^n = \varrho^n - \varrho^{n-1}, \quad \delta \mathbf{u}^n = \mathbf{u}^n - \mathbf{u}^{n-1}, \quad \delta \Phi_n = \Phi_n - \Phi_{n-1}.$$

The system of difference equation corresponding to (4.11) is thus

$$\begin{cases} \delta \varrho_t^{n+1} + \operatorname{div} \left( \delta \varrho^{n+1} \mathbf{w}^{n+1} \right) - \operatorname{div} \left( \varrho^n p'(\varrho^n) \nabla \delta \varrho^{n+1} \right) = \operatorname{div} \left( \delta \varrho^{n+1} \nabla \Phi_n \right) + \widetilde{R}_n, \\ \delta \mathbf{w}_t^{n+1} + \mathbf{u}^n \cdot \nabla \delta \mathbf{w}^{n+1} = -\delta \mathbf{u}^n \cdot \nabla \mathbf{w}^n, \\ -\Delta \delta \Phi_{n+1} = \delta \varrho^{n+1}, \\ \left( \delta \varrho^{n+1}, \delta \mathbf{w}^{n+1} \right) |_{t=0} = (0, \mathbf{0}), \end{cases}$$

where

$$\widetilde{R}_n = \operatorname{div}\left[\left(p'(\varrho^{n-1})\delta\varrho^n + \varrho^n(p'(\varrho^n) - p'(\varrho^{n-1}))\right)\nabla\varrho^n - \varrho^n\delta\mathbf{w}^{n+1}\right] + \operatorname{div}\left(\varrho^n\nabla\delta\Phi_n\right).$$

Performing the estimates similar to those from the proof of lemma 3.2, for the term

$$\|\delta \mathbf{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n+1}\|_{\mathcal{V}_{2}(T)} + \|\delta \Phi_{n+1}\|_{\mathcal{V}_{4}(T)}.$$

We can show the following lemma:

LEMMA 4.5. Under the assumptions of lemma 4.4, we have

$$\begin{aligned} \|\delta \boldsymbol{w}^{n+1}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n+1}\|_{\mathcal{V}_{2}(T)} + \|\delta \Phi_{n+1}\|_{\mathcal{V}_{4}(T)} \\ &\leq E(T) \left(\|\delta \boldsymbol{w}^{n}\|_{L_{\infty}(0,T;H^{2}(\mathbb{T}^{3}))} + \|\delta \varrho^{n}\|_{\mathcal{V}_{2}(T)} + \|\delta \Phi_{n}\|_{\mathcal{V}_{4}(T)}\right), \end{aligned}$$

where E(t) is described in §1.1.

This concludes the contraction argument and also the proof of theorem 4.3.

#### Acknowledgements

The work of N.C. was funded by the 'Excellence Initiative Research University (IDUB)' program at the University of Warsaw. T.P.'s work was supported by the National Science Centre (NCN) project 2022/45/B/ST1/03432. The work of N.C. and E.Z. was supported by the EPSRC Early Career Fellowship no. EP/V000586/1.

#### References

- 1 P. Aceves-Sanchez, R. Bailo, P. Degond and Z. Mercier. Pedestrian models with congestion effects. *Math. Models Methods Appl. Sci.* **34** (2024), 1001–1041.
- 2 S. Agrawal, V. Kanagaraj and M. Treiber. Two-dimensional LWR model for lane-free traffic. *Physica A: Statistical Mechanics and Its Applications* 625 (2023), 128990.
- 3 A. Aw, A. Klar, M. Rascle and T. Materne. Derivation of continuum traffic flow models from microscopic follow-the-leader models. *SIAM J. Math. Anal.* **63** (2002), 259–278.
- 4 A. Aw and M. Rascle. Resurrection of second order models of traffic flow. *SIAM J. Appl. Math.* **60** (2000), 916–938.
- 5 N. Bellomo and F. Brezzi. Active particle methods towards a mathematics of living systems. Math. Models Methods Appl. Sci. 34 (2024), 993–999.
- 6 N. Bellomo, L. Gibelli, A. Quaini and A. Reali. Towards a mathematical theory of behavioral human crowds. *Math. Models Methods Appl. Sci.* **32** (2022), 321–358.
- 7 N. Bellomo, J. Liao, A. Quaini, L. Russo and C. Siettos. Human behavioral crowds review, critical analysis and research perspectives. *Math. Models Methods Appl. Sci.* **33** (2023), 1611–1659.
- 8 F. Berthelin, P. Degond, M. Delitata and M. Rascle. A model for the formation and evolution of traffic jams. *Arch. Ration. Mech. Anal.* **187** (2008), 185–220.
- 9 N. Chaudhuri, E. Feireisl and E. Zatorska. Nonuniqueness of weak solutions to the dissipative Aw-Rascle model. *Appl. Math. Optim.* **90** (2024), 19.
- 10 N. Chaudhuri, P. Gwiazda and E. Zatorska. Analysis of the generalised Aw-Rascle model. Comm. Partial Differential Equations 48 (2023), 440–477.
- 11 N. Chaudhuri, M. A. Mehmood, C. Perrin and E. Zatorska. Duality solutions to the hard-congestion model for the dissipative Aw-Rascle system. *Comm. Partial Differential Equations* 49 (2024), 671—697.
- 12 N. Chaudhuri, L. Navoret, C. Perrin and E. Zatorska. Hard congestion limit of the dissipative Aw-Rascle system. *Nonlinearity* 37 (2024), 04.

32

- 13 N. Chaudhuri, J. Peszek, M. Szlenk and E. Zatorska. Non-local dissipative Aw-Rascle model and its relation with matrix-valued communication in Euler alignment. 2024 arXiv:2409.07593.
- 14 R. Danchin and P. B. Mucha. Incompressible flows with piecewise constant density. Arch. Ration. Mech. Anal. 207 (2013), 991–1023.
- 15 R. Danchin and P. B. Mucha. The incompressible Navier-Stokes equations in vacuum. Comm. Pure Appl. Math. 72 (2019), 1351–1385.
- 16 R. Danchin, P. B. Mucha and T. Piasecki. Stability of the density patches problem with vacuum for incompressible inhomogeneous viscous flows. Ann. Inst. H. Poincaré C Anal. Non Linéaire 41 (2024), 897–931.
- 17 Y. Enomoto and Y. Shibata. On the  $\mathcal{R}$ -sectoriality and the initial boundary value problem for the viscous compressible fluid flow. *Funkcial. Ekvac.* 56 (2013), 441–505.
- 18 L. C. Evans. Partial differential equations, Graduate Studies in Mathematics, Vol.19 (Providence, RI: American Mathematical Soc., 2022).
- 19 E. Feireisl and A. Novotný. Singular limits in thermodynamics of viscous fluids, Advances in Mathematical Fluid Mechanics. 2nd edition (2017) (Birkhäuser Verlag, Basel, 2009).
- 20 M. Herty, A. Fazekas and G. Visconti. A two-dimensional data-driven model for traffic flow on highways. *Networks and Heterogeneous Media* 13 (2018), 217–240.
- 21 D. Hoff. Uniqueness of weak solutions of the Navier-Stokes equations of multidimensional, compressible flow. SIAM J. Math. Anal. 37 (2006), 1742–1760.
- 22 J. Kim. A Cucker-Smale flocking model with the Hessian communication weight and its first-order reduction. J. Nonlinear Sci. **32**, 2022.
- 23 O. Kreml, S. Nečasová and T. Piasecki. Local existence of strong solutions and weak-strong uniqueness for the compressible Navier-Stokes system on moving domains. *Proc. Royal Soc. Edinburgh A* **150** (2020), 2255–2300.
- 24 O. A. Ladyženskaja, V. A. Solonnikov, and N. N. Ural'ceva. Linear and quasi-linear equations of parabolic type, Translations of Mathematical Monographs, Vol. 23 (Providence, RI: American Mathematical Soc., 1968).
- 25 A. Lefebvre-Lepot and B. Maury. Micro-macro modelling of an array of spheres interacting through lubrication forces. Advances in Mathematical Sciences and Applications 21 (2011), 535–557.
- 26 M. A. Mehmood. Hard congestion limit of the dissipative Aw-Rascle system with a polynomial offset function. J. Math. Anal. Appl. 533 (2024), 128028.
- 27 P. B. Mucha and T. Piasecki. Stationary compressible Navier-Stokes equations with inflow condition in domains with piecewise analytical boundaries. *Pure Appl. Anal.* 2 (2020), 123–155.
- 28 P. B. Mucha and W. M. Zajaczkowski. On a  $L_p$ -estimate for the linearized compressible Navier-Stokes equations with the Dirichlet boundary conditions. J. Differential Equations 186 (2002), 377–393.
- 29 P. B. Mucha and W. M. Zajaczkowski. Global existence of solutions of the Dirichlet problem for the compressible Navier-Stokes equations. ZAMM Z. Angew. Math. Mech. 84 (2004), 417–424.
- 30 A. Novotný and I. Straškraba. Introduction to the mathematical theory of compressible flow, Of Oxford Lecture Series in Mathematics and its Applications, Vol.27 (Oxford University Press, Oxford, 2004).
- 31 J. Peszek and D. Poyato. Measure solutions to a kinetic Cucker-Smale model with singular and matrix-valued communication. arXiv:2207.14764.
- 32 T. Piasecki and M. Pokorný. Strong solutions to the Navier-Stokes-Fourier system with slip-inflow boundary conditions. ZAMM Z. Angew. Math. Mech. 94 (2014), 1035–1057.
- 33 L. Tumash, C. Canudas-de-Wit and M. L. D. Monache. Multi-directional continuous traffic model for large-scale urban networks. *Transportation Research Part B: Methodological* 158 (2022), 374–402.
- 34 A. Valli and W. M. Zajaczkowski. Navier-Stokes equations for compressible fluids: Global existence and qualitative properties of the solutions in the general case. *Comm. Math. Phys.* 103 (1986), 259–296.
- 35 W. M. Zajaczkowski. On nonstationary motion of a compressible barotropic viscous fluid with boundary slip condition. J. Appl. Anal. 4 (1998), 167–204.