

# **Hopf bifurcations for a delayed discrete single population patch model in advective environments**

Weiwei Liu<sup>1</sup>, Zuolin Shen<sup>[2](https://orcid.org/0000-0002-8094-9069)</sup> and Shanshan Chen<sup>2</sup>

<sup>1</sup> School of Mathematics, Harbin Institute of Technology, Harbin, Heilongjiang, P.R. China <sup>2</sup>Department of Mathematics, Harbin Institute of Technology, Weihai, Shandong, P.R. China **Corresponding author:** Shanshan Chen; Email: [chenss@hit.edu.cn](mailto:chenss@hit.edu.cn)

**Received:** 18 November 2023; **Revised:** 15 May 2024; **Accepted:** 04 August 2024

**Keywords:** Hopf bifurcations; delay; directed drift; random movement

**2020 Mathematics Subject Classification:** 34K18 (Primary); 92D25 (Secondary)

## **Abstract**

In this paper, we consider a delayed discrete single population patch model in advective environments. The individuals are subject to both random and directed movements, and there is a net loss of individuals at the downstream end due to the flow into a lake. Choosing time delay as a bifurcation parameter, we show the existence of Hopf bifurcations for the model. In homogeneous non-advective environments, it is well known that the first Hopf bifurcation value is independent of the dispersal rate. In contrast, for homogeneous advective environments, the first Hopf bifurcation value depends on the dispersal rate. Moreover, we show that the first Hopf bifurcation value in advective environments is larger than that in non-advective environments if the dispersal rate is large or small, which suggests that directed movements of the individuals inhibit the occurrence of Hopf bifurcations.

## **1. Introduction**

Time delays can induce periodic solutions through Hopf bifurcations, and this phenomenon can explain population oscillations in the natural world. Take the following *n*-patch single population model

$$
\frac{du_j}{dt} = \sum_{k=1}^n L_{jk} u_k + u_j [a_j - b_j u_j (t - \tau)], \ j = 1, \cdots, n, \ t > 0
$$
\n(1.1)

for instance. Here,  $u_i$  denotes the population density in patch *j*,  $(L_{ik})$  is the dispersion matrix,  $\tau$  is the time delay and represents the maturation time, and  $a_i$  and  $b_j$  represent the intrinsic growth rate and intraspecific competition of the species in patch *j*, respectively. It was shown that time delay can induce Hopf bifurcations for model [\(1.1\)](#page-0-0) when the dispersion matrix  $(L_{jk}) = \epsilon(\hat{L}_{jk})$  with  $0 < \epsilon \ll 1$  or  $\epsilon \gg 1$  [\[8,](#page-29-0) [24\]](#page-30-0). Especially, if  $n = 2$ , delay-induced Hopf bifurcations can occur for a wider range of parameters [\[31\]](#page-30-1).

The species in streams are subject to both random and directed movements. For example, the following Figure [1](#page-1-0) represents stream to a lake, and the diffusive flux into and from the lake balances. Therefore, the flux into the lake at the downstream end is only the advective flux, and one can refer to [\[34,](#page-30-2) [35,](#page-30-3) [49\]](#page-31-0) for more biological explanation. For the river network illustrated in Figure [1,](#page-1-0) the dispersion matrix  $(L_{ik})$ in model  $(1.1)$  takes the following form:

<span id="page-0-1"></span>
$$
(L_{jk}) = dD + qQ. \tag{1.2}
$$

Here, *d* and *q* are the random diffusion rate and drift rate, respectively; and  $D = (D_{ii})$  and  $Q = (Q_{ii})$ represent the diffusion pattern and directed movement pattern of individuals, respectively, where

<sup>C</sup> The Author(s), 2024. Published by Cambridge University Press. This is an Open Access article, distributed under the terms of the Creative Commons Attribution licence [\(http://creativecommons.org/licenses/by/4.0/\)](http://creativecommons.org/licenses/by/4.0/), which permits unrestricted re-use, distribution and reproduction, provided the original article is properly cited.

<span id="page-0-2"></span><span id="page-0-0"></span>

<span id="page-1-0"></span>

*Figure 1. A sample river network.*

$$
D_{jk} = \begin{cases} 1, & j = k - 1 \text{ or } j = k + 1, \\ -2, & j = k = 2, \dots, n - 1, \\ -1, & j = k = 1, n, \\ 0, & \text{otherwise,} \end{cases}
$$
(1.3)

<span id="page-1-1"></span>and

$$
Q_{jk} = \begin{cases} 1, & j = k + 1, \\ -1, & j = k = 1, \cdots, n, \\ 0, & \text{otherwise.} \end{cases}
$$
 (1.4)

The population dynamics in streams have been studied extensively, and it can be modelled by discrete patch models or partial differential equations (PDE) models. It is well known that the stream flow takes individuals to the downstream locations, which is unfavourable for their persistence, see, for example, [\[35,](#page-30-3) [36,](#page-30-4) [45,](#page-31-1) [49\]](#page-31-0) for PDE models. The directed drift is also a disadvantage for two competing species, and to win the competition, the species need a faster random movement rate to compensate the net loss induced by directed movements, see, for example, [\[3,](#page-29-1) [13,](#page-30-5) [35,](#page-30-3) [36,](#page-30-4) [53\]](#page-31-2) for PDE models and, for example,  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  $[6, 22, 25, 26, 42]$  for discrete patch models. A natural question is:

 $(Q_1)$  Whether model [\(1.1\)](#page-0-0) undergoes Hopf bifurcations with  $(L_{ik})$  defined in [\(1.2\)](#page-0-1) and how directed movements of individuals affect Hopf bifurcations?

We remark that *d* and *q* are normally not proportional, and consequently results in [\[8,](#page-29-0) [24\]](#page-30-0) cannot apply to this type of dispersion matrix.

In this paper, we aim to answer question  $(Q_1)$  and consider the river network illustrated in Figure [1.](#page-1-0) To emphasise the effect of directed drift, we exclude the effect of spatial heterogeneity, and let  $a_1 = \cdots = a_n = a$  and  $b_1 = \cdots = b_n = b$  in model [\(1.1\)](#page-0-0). Then model (1.1) is reduced to the following form:

<span id="page-1-2"></span>
$$
\begin{cases}\n\frac{du_j}{dt} = \sum_{k=1}^n (dD_{jk} + qQ_{jk}) u_k + u_j (a - bu_j (t - \tau)), & t > 0, j = 1, \dots, n, \\
u(t) = \psi(t) \ge 0, & t \in [-\tau, 0],\n\end{cases}
$$
\n(1.5)

where  $n \geq 2$  is the number of patches, *d* and *q* are the random diffusion rate and drift rate, respectively,  $D = (D_{ik})$  and  $Q = (Q_{ik})$  are defined in [\(1.3\)](#page-0-2) and [\(1.4\)](#page-1-1), respectively, and parameters *a*, *b*,  $\tau > 0$  have the same meanings as the above model  $(1.1)$ .

<span id="page-1-3"></span>Our study is also motivated by some researches on reaction–diffusion models with time delay. One can refer to [\[1,](#page-29-3) [2,](#page-29-4) [9,](#page-29-5) [12,](#page-30-10) [19,](#page-30-11) [20,](#page-30-12) [23,](#page-30-13) [46,](#page-31-3) [50,](#page-31-4) [51\]](#page-31-5) and [\[7,](#page-29-6) [11,](#page-30-14) [27,](#page-30-15) [30,](#page-30-16) [33,](#page-30-17) [38,](#page-30-18) [41,](#page-30-19) [47\]](#page-31-6) for reaction–diffusion models without and with advection term, respectively. The following reaction–diffusion model with time delay

$$
\begin{cases}\n u_t = \hat{d}u_{xx} - \hat{q}u_x + u (a - bu(t - \tau)), & x \in (0, l), t > 0 \\
\hat{d}u_x(0, t) - \hat{q}u(0, t) = 0, & t > 0 \\
\hat{d}u_x(l, t) - \hat{q}u(l, t) = -\beta \hat{q}u(l, t), & t > 0 \\
u(x, t) = \psi(x, t) \ge 0, & x \in (0, l), t \in [-\tau, 0]\n\end{cases}
$$
\n(1.6)

models population dynamics in streams, where  $\beta \ge 0$  represents the loss of individuals at the down-stream end. Actually, model [\(1.5\)](#page-1-2) can be viewed as a discrete version of model [\(1.6\)](#page-1-3) with  $\beta = 1$ , which describes streams into a lake at the downstream end. Divide the interval  $[0, l]$  into  $n + 1$  sub-intervals with equal length  $\Delta x = l/(n + 1)$ , and denote the endpoints by 0, 1,  $\cdots$ ,  $n + 1$ . Discretising the spatial variable of the first equation of [\(1.6\)](#page-1-3) at endpoints  $j = 1, \dots, n$ , we obtain the following equation:

$$
\frac{du_j}{dt} = \hat{d}\frac{u_{j+1} - 2u_j + u_{j-1}}{(\Delta x)^2} - \hat{q}\frac{u_j - u_{j-1}}{\Delta x} + u_j(a - bu_j(t - \tau)), \ \ j = 1, \cdots, n,
$$
\n(1.7)

where  $u_j(t)$  is the population density at endpoint *j*. Let  $d = \hat{d}/(\Delta x)^2$  and  $q = \hat{q}/\Delta x$ . Then we obtain [\(1.5\)](#page-1-2) from [\(1.7\)](#page-2-0) for  $j = 2, \dots, n - 1$ . At the upstream end, we discretise the no-flux boundary condition and obtain that

<span id="page-2-2"></span><span id="page-2-0"></span>
$$
d(u_1 - u_0) - qu_0 = 0. \t\t(1.8)
$$

<span id="page-2-1"></span>Plugging [\(1.8\)](#page-2-1) into [\(1.7\)](#page-2-0) with  $j = 1$ , we obtain [\(1.5\)](#page-1-2) for  $j = 1$ . Similarly, discretising the boundary condition at the downstream end with  $\beta = 1$ , we have  $u_n = u_{n+1}$ . Plugging it into [\(1.7\)](#page-2-0), we obtain [\(1.5\)](#page-1-2) for  $j = n$ . The above discretisation for model [\(1.6\)](#page-1-3) with  $\tau = 0$  can be found in [\[10,](#page-29-7) [34\]](#page-30-2), and here we include it for the sake of completeness.

For the non-advective case ( $\hat{q} = 0$ ), model [\(1.6\)](#page-1-3) admits a unique positive steady state  $u = a/b$ , and it was shown in  $[40, 52]$  $[40, 52]$  $[40, 52]$  that large delay can make such a constant steady state unstable and induce Hopf bifurcations. If *d* and  $\hat{q}$  are proportional with  $\hat{q} \neq 0$ , delay-induced Hopf bifurcations can also be investigated. Letting  $\tilde{u} = e^{-\frac{\tilde{q}}{\tilde{q}}/\tilde{d}x}u$  and  $\tilde{t} = dt$ , model [\(1.6\)](#page-1-3) can be transformed as follows:

$$
\begin{cases}\n\tilde{u}_{\tilde{t}} = e^{-\lambda x} (e^{\lambda x} \tilde{u}_x)_x + r \tilde{u} \left( a - b e^{\lambda x} \tilde{u} (\tilde{t} - \tilde{\tau}) \right), & x \in (0, l), \ \tilde{t} > 0, \\
\tilde{u}_x(0, \tilde{t}) = 0, & \tilde{t} > 0, \\
\tilde{u}_x(l, \tilde{t}) = -\beta \lambda \tilde{u}, & \tilde{t} > 0,\n\end{cases}
$$
\n(1.9)

where  $\lambda = \frac{\partial}{\partial q}$ ,  $r = 1/\hat{d}$  and  $\tilde{\tau} = \hat{d}\tau$ . For the case of  $\beta = 0$ , it was shown in [\[11\]](#page-30-14) that delay can induce Hopf bifurcations for model  $(1.9)$  if  $r \ll 1$ , which implies that delay-induced Hopf bifurcations can occur if  $\hat{d}$  and  $\hat{q}$  are proportional and both large for the original model [\(1.6\)](#page-1-3). To our knowledge, the case that  $\hat{d}$ and  $\hat{q}$  are not proportional is also unknown for model [\(1.6\)](#page-1-3). Our study on question (Q<sub>1</sub>) also solves this problem in a discrete setting.

The main results of the paper are summarised as follows. It is well known that large delay can induce Hopf bifurcations for model [\(1.5\)](#page-1-2) if the directed drift rate  $q = 0$  (the non-advective case), and the first Hopf bifurcation value is  $\tau_{non} = \pi/2a$ , which is independent of the random diffusion rate *d* (Proposition [5.1\)](#page-24-0). In contrast, if  $q \neq 0$  (the advection case), the first Hopf bifurcation value  $\tau_{adv}$  depends on *d* and is strictly monotone decreasing in  $d \in (d_3, \infty)$  with  $d_3 \gg 1$  (Proposition [5.2\)](#page-25-0). Moreover, we show that  $\tau_{adv} > \tau_{non}$  for  $d \gg 1$  or  $d \ll 1$  (Propositions [5.2](#page-25-0) and [5.3\)](#page-26-0), which suggests that directed movements of the individuals inhibit the occurrence of Hopf bifurcations. The comparison of Hopf bifurcation values between non-advective and advective cases is illustrated in Figure [2.](#page-3-0) Moreover, we obtain that the total population size is strictly increasing in  $d \in [\delta, \infty)$  with  $\delta \gg 1$  (Proposition [2.6](#page-8-0) and remark [2.7\)](#page-10-0).

For patch models in homogeneous non-advective environments, one can refer to  $[16-18]$  $[16-18]$  for the framework of Turing and Hopf bifurcations, see also [\[4,](#page-29-8) [15\]](#page-30-23) for cross diffusion-induced Turing bifurcations and [\[5,](#page-29-9) [32,](#page-30-24) [39,](#page-30-25) [43,](#page-30-26) [44,](#page-30-27) [48\]](#page-31-8) for delay-induced Hopf bifurcations. For homogeneous advective

<span id="page-3-0"></span>

*Figure 2. The comparison of Hopf bifurcation values.*

environments, one cannot obtain the explicit expressions for the positive equilibria. This brings some difficulties in bifurcation analysis, and we overcome them by constructing equivalent eigenvalue problems in this paper.

The rest of the paper is organised as follows. In Section [2,](#page-3-1) we give some preliminaries and obtain some properties for the unique positive equilibrium  $u_d$  of model [\(1.5\)](#page-1-2). In Section [3,](#page-10-1) we study the eigenvalue problem associated with the positive equilibrium  $u_d$  for three cases. In Section [4,](#page-21-0) we obtain the local dynamics and the existence of Hopf bifurcations for model [\(1.5\)](#page-1-2). Finally, we show the effect of drift rate on Hopf bifurcation values and give some numerical simulations in Section [5.](#page-24-1)

## <span id="page-3-1"></span>**2. Preliminary**

We first list some notations for later use. Denote  $\mathbf{1} = (1, \dots, 1)^T$  and define the real and imaginary parts of  $\mu \in \mathbb{C}$  by  $\mathcal{R}e\mu$  and  $\mathcal{I}m\mu$ , respectively. For a space *Z*, we denote complexification of *Z* to be  $Z_{\mathbb{C}} := \{x_1 + \mathrm{i}x_2 | x_1, x_2 \in \mathbb{Z}\}$ . For a linear operator *T*, we define the domain, the range and the kernel of *T* by  $\mathscr{D}(T)$ ,  $\mathscr{D}(T)$  and  $\mathscr{M}(T)$ , respectively. For  $\mathbb{C}^n$ , we choose the inner product  $\langle u, v \rangle = \sum_{j=1}^n \overline{u}_j v_j$  for  $u, v \in \mathbb{C}^n$  and denote

$$
\|\boldsymbol{u}\|_{\infty} = \max_{j=1,\cdots,n} |u_j|, \ \|\boldsymbol{u}\|_2 = \left(\sum_{j=1}^n |u_j|^2\right)^{1/2}
$$

.

For  $u = (u_1, \dots, u_n)^T \in \mathbb{R}^n$ , we write  $u \gg 0$  if  $u_j > 0$  for all  $j = 1, \dots, n$ . For an  $n \times n$  real-valued matrix *M* we denote the spectral bound of *M* by: *M*, we denote the spectral bound of *M* by:

 $s(M) := \max\{ \mathcal{R}e\mu : \mu \text{ is an eigenvalue of } M \},$ 

and the spectral radius of *M* by:

$$
\rho(M) := \max\{|\mu| : \mu \text{ is an eigenvalue of } M\}.
$$

An eigenvalue of *M* with a positive eigenvector is called the principal eigenvalue of *M*. A real-valued square matrix *M* with non-negative off-diagonal entries is referred to as the essentially non-negative matrix. If *M* is an irreducible essentially non-negative matrix, then there exists  $c > 0$  such that  $M + cI$ is an irreducible non-negative matrix. It follows from [\[28,](#page-30-28) Theorem 2.1] that

- (i)  $\rho(M + cI)$  is positive and is an algebraically simple eigenvalue of  $M + cI$  with a positive eigenvector.
- (ii)  $\rho(M + cI)$  is the unique eigenvalue with a non-negative eigenvector.

By (i), we have  $s(M) + c = s(M + cI) = \rho(M + cI)$ , and consequently,  $s(M)$  is an algebraically simple eigenvalue of *M* with a positive eigenvector. By (ii), we see that  $s(M)$  is the unique principal eigenvalue of *M*.

Consider the following eigenvalue problem:

<span id="page-4-0"></span>
$$
\sum_{k=1}^{n} dD_{jk}\phi_k + \sum_{k=1}^{n} qQ_{jk}\phi_k + a\phi_j = \lambda \phi_j, \ \ j = 1, \cdots, n,
$$
\n(2.1)

where  $(D_{ik})$  and  $(Q_{ik})$  are defined in [\(1.3\)](#page-0-2) and [\(1.4\)](#page-1-1), respectively. Since  $dD + qQ + aI$  is an irreducible and essentially non-negative matrix, it follows that  $(2.1)$  admits a unique principal eigenvalue  $\lambda_1(d, q)$ with

$$
\lambda_1(d,q) = s (dD + qQ + aI).
$$

The global dynamics of model [\(1.5\)](#page-1-2) for  $\tau = 0$  is determined by the sign of  $\lambda_1(d, q)$  (see [\[14,](#page-30-29) [29,](#page-30-30) [37\]](#page-30-31) for the proof).

<span id="page-4-2"></span>**Lemma 2.1.** *Suppose that*  $\tau = 0$ *. If*  $\lambda_1(d, q) \leq 0$ *, then the trivial equilibrium* 0 *of model* [\(1.5\)](#page-1-2) *is globally asymptotically stable; if*  $\lambda_1(d, q) > 0$ , *then model* [\(1.5\)](#page-1-2) *admits a unique positive equilibrium, which is globally asymptotically stable.*

For later use, we cite the following result from [\[10\]](#page-29-7).

<span id="page-4-1"></span>**Lemma 2.2.** *Let*  $\lambda_1(d, q)$  *be the principal eigenvalue of* [\(2.1\).](#page-4-0) *Then the following statements hold:* 

- *(i)* For fixed  $d > 0$ ,  $\lambda_1(d, q)$  is strictly decreasing with respect to q in  $[0, \infty)$ , and there exists  $q_d^* > 0$ *such that*  $\lambda_1(d, q_d^*) = 0$ ,  $\lambda_1(d, q) < 0$  *for*  $q > q_d^*$ , and  $\lambda_1(d, q) > 0$  *for*  $q < q_d^*$ ,
- *(ii)*  $q_d^*$  *is strictly increasing in*  $d \in (0, \infty)$  *with*  $\lim_{d\to 0} q_d^* = a$  *and*  $\lim_{d\to \infty} q_d^* = na$ .

Here, we remark that Lemma [2.2](#page-4-1) (i) follows from [\[10,](#page-29-7) Lemma [3.1](#page-11-0) and Proposition 3.2 (i)], and Lemma  $2.2$  (ii) follows from [\[10,](#page-29-7) Lemma  $3.7$ ]. The following result is deduced directly from Lemma [2.2.](#page-4-1)

<span id="page-4-3"></span>**Lemma 2.3.** Let  $\lambda_1(d, q)$  be the principal eigenvalue of [\(2.1\).](#page-4-0) Then the following statements hold:

- *(i) If*  $q \in (0, a]$ *, then*  $\lambda_1(d, q) > 0$  *for all*  $d > 0$ *;*
- (ii) If  $q \in (a, na)$ , then there exists  $d_q^* > 0$  such that  $\lambda_1(d_q^*, q) = 0$ ,  $\lambda_1(d, q) < 0$  for  $0 < d < d_q^*$  and  $\lambda_1(d, q) > 0$  for  $d > d_q^*$ ;
- *(iii) If*  $q \in [na, \infty)$ *, then*  $\lambda_1(d, q) < 0$  *for all*  $d > 0$ *.*

By Lemmas [2.1](#page-4-2) and [2.3,](#page-4-3) we obtain the global dynamics of model [\(1.5\)](#page-1-2) for  $\tau = 0$ .

<span id="page-4-4"></span>**Proposition 2.4.** *Suppose that*  $d$ ,  $q$ ,  $a$ ,  $b > 0$  *and*  $\tau = 0$ *. Then the following statements hold:* 

- *(i)* If  $q \in (0, a]$ , then model [\(1.5\)](#page-1-2) admits a unique positive equilibrium  $u_d \gg 0$  for all  $d > 0$ *, which is globally asymptotically stable;*
- *(ii)* If  $q \in (a, na)$ , then the trivial equilibrium **0** *of model* [\(1.5\)](#page-1-2) *is globally asymptotically stable for d* ∈ (0, *d*<sup>∗</sup><sub>1</sub>); and for *d* ∈ (*d*<sup>∗</sup><sub>1</sub></sub>, ∞), model [\(1.5\)](#page-1-2) admits a unique positive equilibrium  $u_d \gg 0$ , which is alobally asymptotically stable: *is globally asymptotically stable;*
- *(iii)* If  $q \in [na, \infty)$ , then the trivial equilibrium **0** of model [\(1.5\)](#page-1-2) is globally asymptotically stable.

Clearly,  $\boldsymbol{u}_d$  satisfies

<span id="page-5-7"></span>
$$
\sum_{k=1}^{n} (dD_{jk} + qQ_{jk})u_k + u_j (a - bu_j) = 0, \ \ j = 1, \cdots, n.
$$
 (2.2)

For simplicity, we first list some notations. Define

<span id="page-5-9"></span><span id="page-5-1"></span>
$$
\mathcal{L} = (\mathcal{L}_{jk}) := d_q^* D + qQ + aI,\tag{2.3}
$$

where  $d_q^*$  is defined in Lemma [2.3.](#page-4-3) It follow from Lemma [2.3](#page-4-3) (ii) that 0 is the principal eigenvalue of  $\mathcal L$ and a corresponding eigenvector is

<span id="page-5-0"></span>
$$
\eta = (\eta_1, \dots, \eta_n)^T \gg 0
$$
 with  $\sum_{i=1}^n \eta_i = 1.$  (2.4)

Clearly, 0 is also the principal eigenvalue of  $\mathcal{L}^T$  and a corresponding eigenvector is

$$
\hat{\boldsymbol{\eta}} = (\hat{\eta}_1, \cdots, \hat{\eta}_n)^T \gg 0 \text{ with } \sum_{i=1}^n \hat{\eta}_i = 1.
$$
 (2.5)

Here, we remark that 0 is an algebraically simple eigenvalue of  $\mathcal L$  and  $\mathcal L^T$ , and the corresponding eigenvector is unique up to multiplying by a scalar. Then, we have the following decompositions:

<span id="page-5-4"></span>
$$
\mathbb{R}^n = \text{span}\{\boldsymbol{\eta}\} \oplus X_1 = \text{span}\{\mathbf{1}\} \oplus X_1,\tag{2.6}
$$

<span id="page-5-2"></span>where

$$
X_1 := \left\{ (x_1, \cdots, x_n)^T \in \mathbb{R}^n : \sum_{i=1}^n \hat{\eta}_i x_i = 0 \right\}.
$$
 (2.7)

In fact, for any  $y = (y_1, \dots, y_n)^T \in \mathbb{R}^n$ ,  $y = c_1 1 + \xi_1 = c_2 \eta + \xi_2$ , where

$$
c_1 = \sum_{i=1}^n \hat{\eta}_i y_i, \ \xi_1 = \mathbf{y} - c_1 \mathbf{1} \in X_1, \ \ c_2 = \frac{\sum_{i=1}^n \hat{\eta}_i y_i}{\sum_{i=1}^n \eta_i \hat{\eta}_i}, \ \xi_2 = \mathbf{y} - c_2 \mathbf{\eta} \in X_1.
$$

Now, we explore further properties on the positive equilibrium *<sup>u</sup><sup>d</sup>*.

<span id="page-5-8"></span>**Proposition 2.5.** *Let*  $\eta$ ,  $\hat{\eta}$  *and*  $X_1$  *be defined in* [\(2.4\),](#page-5-0) [\(2.5\)](#page-5-1) *and* [\(2.7\),](#page-5-2) *respectively. Then the following statements hold:*

*(i)* For fixed  $q \in (0, a)$ ,  $u_d$  is continuously differentiable with respect to  $d \in [0, \infty)$ , where  $u_d = u_0$ *for*  $d = 0$ *, and*  $u_0$  *is the unique solution of* 

$$
\begin{cases}\n u_{0,1} = \frac{a - q}{b}, & qu_{0,j-1} = -u_{0,j} \left( a - q - bu_{0,j} \right) \text{ for } j = 2, \cdots, n, \\
 u_{0,j} > 0 \text{ for } j = 1, \cdots, n.\n\end{cases} \tag{2.8}
$$

<span id="page-5-10"></span>*Moreover,*

<span id="page-5-5"></span><span id="page-5-3"></span>
$$
u_{0,1} < \ldots < u_{0,n};\tag{2.9}
$$

*(ii)* For fixed  $q \in (a, na)$ ,  $\mathbf{u}_d$  can be represented as follows:

$$
\boldsymbol{u}_d = (d - d_q^*)(\alpha_d \boldsymbol{\eta} + \boldsymbol{\xi}_d) \text{ for } d > d_q^*.
$$
 (2.10)

*Here,*  $(\alpha_d, \xi_d) \in \mathbb{R} \times X_1$  *is continuously differentiable with respect to*  $d \in [d^*_q, \infty)$ *, and for*  $d = d^*_q$ *,*  $(\alpha_d, \xi_d) = (\alpha^*, \mathbf{0})$  *with* 

<span id="page-5-6"></span>
$$
\alpha^* = \frac{\sum_{j,k=1}^n D_{jk} \eta_k \hat{\eta}_j}{b \sum_{j=1}^n \eta_j^2 \hat{\eta}_j} > 0.
$$
\n(2.11)

**Proof.** (i) It follows from Proposition [2.4](#page-4-4) that  $u_d$  is the unique positive equilibrium of [\(1.5\)](#page-1-2), which is stable (non-degenerate). Therefore, by the implicit function theorem, we obtain that  $u_d$  is continuously differentiable for  $d \in (0, \infty)$ . Then we need to show that  $u_d$  is continuously differentiable for  $d \in [0, d_1]$ with  $0 < d_1 \ll 1$ .

Define

$$
H(d, u) = \begin{pmatrix} d \sum_{k=1}^{n} D_{1k} u_k - qu_1 + u_1 (a - bu_1) \\ d \sum_{k=1}^{n} D_{2k} u_k + q (u_1 - u_2) + u_2 (a - bu_2) \\ \vdots \\ d \sum_{k=1}^{n} D_{nk} u_k + q (u_{n-1} - u_n) + u_n (a - bu_n) \end{pmatrix}.
$$

Clearly,  $H(0, u_0) = 0$ , where  $u_0$  is defined by [\(2.8\)](#page-5-3). Let  $D_u H(0, u_0)$  be the Jacobian matrix of  $H(d, u)$ with respect to *u* at  $(0, u_0)$ . A direct computation implies that  $D_u H(0, u_0) = (h_{ik})$  with

$$
h_{j,k} = \begin{cases} a - q - 2bu_{0,j} & j = k = 1, \cdots, n, \\ q, & j = k + 1, \\ 0, & \text{otherwise.} \end{cases}
$$

By [\(2.8\)](#page-5-3), we have  $\frac{a-q}{b} = u_{0,1} < \ldots < u_{0,n}$ , which implies that  $D_u H(0, u_0)$  is invertible. Then we see from the implicit function theorem that there exist  $d_1 > 0$ , and a continuously differentiable mapping

$$
d \in [0, d_1] \mapsto \mathbf{u}(d) = (u_1(d), \cdots, u_n(d))^T \in \mathbb{R}^n
$$

such that  $H(d, u(d)) = 0$ ,  $u(d) \gg 0$ , and  $u(0) = u_0$ . Therefore,  $u(d)$  is the positive equilibrium of model [\(1.5\)](#page-1-2) for small *d*. This combined with Proposition [2.4](#page-4-4) implies that  $u(d) = u_d$ . Consequently,  $u_d$  is continuously differentiable for  $d \in [0, \infty)$ .

(ii) Using similar arguments as in (i), we see that  $u_d$  is continuously differentiable for  $d \in (d^*_\tau, \infty)$ . By first decomposition in (2.6), we see that *u*, can be represented as in (2.10), where  $\alpha$ , and **k**, are als the first decomposition in [\(2.6\)](#page-5-4), we see that  $u_d$  can be represented as in [\(2.10\)](#page-5-5), where  $\alpha_d$  and  $\xi_d$  are also continuously differentiable for  $d \in (d_q^*, \infty)$ . Then we need to show that  $\alpha_d$  and  $\xi_d$  are also continuously differentiable for  $d \in [d_q^*, \tilde{d}_1)$  with  $0 < \tilde{d}_1 - d_q^* \ll 1$ .

We first show that  $\alpha^* > 0$ . A direct computation yields

<span id="page-6-0"></span>
$$
\sum_{j,k=1}^{n} D_{jk} \eta_k \hat{\eta}_j = \sum_{j=1}^{n-1} (\eta_{j+1} - \eta_j) (\hat{\eta}_j - \hat{\eta}_{j+1}). \qquad (2.12)
$$

Noticing that *η* (respectively,  $\hat{\eta}$ ) is an eigenvector of *L* (respectively, *L<sup>T</sup>*) corresponding to eigenvalue 0, we have

$$
\begin{aligned} (d_q^* + q)(\eta_{n-1} - \eta_n) &= -a\eta_n, \\ (d_q^* + q)(\eta_{j-1} - \eta_j) &= -a\eta_j + d_q^*(\eta_j - \eta_{j+1}), \ j = 2, \cdots, n-1, \end{aligned}
$$

and

$$
\begin{aligned} (d_q^* + q)(\hat{\eta}_2 - \hat{\eta}_1) &= -a\hat{\eta}_1, \\ (d_q^* + q)(\hat{\eta}_{j+1} - \hat{\eta}_j) &= -a\eta_j + d_q^*(\hat{\eta}_j - \hat{\eta}_{j-1}), \ j = 2, \cdots, n-1, \end{aligned}
$$

which implies that  $\eta_1 < \eta_2 < \ldots < \eta_n$  and  $\hat{\eta}_1 > \hat{\eta}_2 > \ldots > \hat{\eta}_n$ . This combined with [\(2.11\)](#page-5-6) and [\(2.12\)](#page-6-0) yields  $\alpha^* > 0$ .

#### 8 *W. Liu et al.*

By the definition of  $\mathcal{L}$ , we rewrite [\(2.2\)](#page-5-7) as follows:

<span id="page-7-0"></span>
$$
\sum_{k=1}^{n} \mathcal{L}_{jk} u_k + (d - d_q^*) \sum_{j=1}^{n} D_{jk} u_k - b u_j^2 = 0.
$$
 (2.13)

From the first decomposition in  $(2.6)$ , we see that  $\boldsymbol{u}$  in  $(2.13)$  can be represented as follows:

<span id="page-7-2"></span>
$$
\mathbf{u} = (d - d_q^*)(\alpha \eta + \xi) \text{ for } d > d_q^*.
$$
 (2.14)

<span id="page-7-1"></span>٦

Plugging  $(2.14)$  into  $(2.13)$ , we have

$$
\sum_{k=1}^{n} \mathcal{L}_{jk}\xi_k + (d - d_q^*) \left[ \sum_{j=1}^{n} D_{jk}(\alpha \eta_k + \xi_k) - b(\alpha \eta_j + \xi_j)^2 \right] = 0, \quad j = 1, \cdots, n. \tag{2.15}
$$

Denoting the left side of [\(2.15\)](#page-7-2) by *yj*, we see from the second decomposition in [\(2.6\)](#page-5-4) that

$$
y = (y_1, \dots, y_n)^T = c\mathbf{1} + z
$$
 with  $c = \sum_{j=1}^n \hat{\eta}_j y_j$  and  $z \in X_1$ .

Therefore,  $y = 0$  if and only if  $c = 0$  and  $z = 0$ . Since  $\mathcal{L} \xi \in X_1$ , it follows that

$$
c = (d - d_q^*)G_2
$$
 and  $z = (G_{1,1}, \cdots, G_{1,n})^T$ ,

where

$$
G_{1,j}(d, \alpha, \xi) = \sum_{k=1}^{n} \mathcal{L}_{jk}\xi_k + (d - d_q^*) \left[ \sum_{j=1}^{n} D_{jk}(\alpha \eta_k + \xi_k) - b(\alpha \eta_j + \xi_j)^2 - (d - d_q^*)G_2(d, \alpha, \xi), \quad j = 1, \cdots, n,
$$
  

$$
G_2(d, \alpha, \xi) = \sum_{j=1}^{n} \hat{\eta}_j \left[ \sum_{k=1}^{n} D_{jk}(\alpha \eta_k + \xi_k) - b(\alpha \eta_j + \xi_j)^2 \right].
$$

Define *G*(*d*, α, *ξ*):  $\mathbb{R}^2 \times X_1 \to X_1 \times \mathbb{R}$  by *G* := (*G*<sub>1,1</sub>, ···, *G*<sub>1,*n*</sub>, *G*<sub>2</sub>)<sup>*T*</sup>. It follows that, for *d* > *d*<sup>∗</sup><sub>*g*</sub>, *u*<sub>*g*</sub> (*represented in* (2.14)) is a solution of (2.13) if and only if (represented in [\(2.14\)](#page-7-1)) is a solution of [\(2.13\)](#page-7-0) if and only if  $(\alpha, \xi) \in \mathbb{R} \times X_1$  is a solution of  $G(d, \alpha, \xi) = 0$ .

Now we consider the equivalent problem  $G(d, \alpha, \xi) = 0$ . Clearly,  $G(d^*_q, \alpha^*, 0) = 0$ . Let  $T\left(\check{\alpha}, \check{\xi}\right) =$  $(T_{1,1}, \dots, T_{1,n}, T_2)^T : \mathbb{R} \times X_1 \mapsto X_1 \times \mathbb{R}$  be the Fréchet derivative of *G* with respect to  $(\alpha, \xi)$  at  $(d_q^*, \alpha^*, \mathbf{0})$ .<br>Then we compute that Then we compute that

$$
T_{1,j}\left(\check{\alpha},\check{\xi}\right) = \sum_{k=1}^n \mathcal{L}_{jk}\check{\xi}_k, \ j = 1,\cdots,n,
$$
  

$$
T_2\left(\check{\alpha},\check{\xi}\right) = \sum_{j=1}^n \hat{\eta}_j\left(\sum_{k=1}^n D_{jk}\eta_k - 2b\alpha^*\eta_j^2\right)\check{\alpha} + \sum_{j=1}^n \hat{\eta}_j\left(\sum_{k=1}^n D_{jk}\check{\xi}_k - 2b\alpha^*\eta_j\check{\xi}_j\right).
$$

By [\(2.11\)](#page-5-6), we see that *T* is a bijection from  $\mathbb{R} \times X_1$  to  $X_1 \times \mathbb{R}$ . Then it follows from the implicit function theorem that there exists  $\tilde{d}_1 > d_q^*$  and a continuously differentiable mapping  $d \in \left[d_q^*, \tilde{d}_1\right] \mapsto \left(\tilde{\alpha}_d, \tilde{\xi}_d\right) \in$  $\mathbb{R} \times X_1$  such that  $G(d, \tilde{\alpha}_d, \tilde{\xi}_d) = 0$ , and  $\tilde{\alpha}_d = \alpha^*$  and  $\tilde{\xi}_d = 0$  for  $d = d^*_d$ . It follows from Proposition [2.4](#page-4-4) and Eq. [\(2.6\)](#page-5-4) that the unique positive equilibrium  $u_d$  can be represented as [\(2.10\)](#page-5-5) for  $d > d_q^*$ . Then we obtain that  $\alpha_d = \tilde{\alpha}_d$ ,  $\xi_d = \tilde{\xi}_d$  for  $d \in (d^*_q, \tilde{d}_1]$ . Therefore,  $\alpha_d$  and  $\xi_d$  are continuously differentiable for  $d \in [d_q^*, \infty)$  if we define  $(\alpha_d, \xi_d) = (\alpha^*, \mathbf{0})$  for  $d = d_q^*$ .  $\Box$  Now, we consider the case  $d \gg 1$ . Clearly,  $u_d$  satisfies

<span id="page-8-1"></span>
$$
\sum_{k=1}^{n} D_{jk} u_k + \lambda \left[ q \sum_{k=1}^{n} Q_{jk} u_k + u_j (a - bu_j) \right] = 0, \ j = 1, \cdots, n
$$
 (2.16)

with  $d = 1/\lambda$ . To avoid confusion, we denote  $u_d$  by  $u^{\lambda}$  for the case  $d \gg 1$ . Then the properties of  $u_d$ for  $d \gg 1$  is equivalent to  $u^{\lambda}$  for  $0 < \lambda \ll 1$ . Clearly,  $s(D) = 0$  is the principal eigenvalue of *D*, and a corresponding eigenvector is corresponding eigenvector is

<span id="page-8-5"></span>
$$
\mathbf{S} = (S_1, \cdots, S_n)^T \text{ with } S_j = \frac{1}{n} \text{ for all } j = 1, \cdots, n. \tag{2.17}
$$

<span id="page-8-6"></span>Define

$$
\widetilde{X}_1 = \left\{ (x_1, \cdots, x_n)^T \in \mathbb{R}^n : \sum_{j=1}^n x_j = 0 \right\},
$$
\n(2.18)

and  $\mathbb{R}^n$  also has the following decomposition:

<span id="page-8-2"></span>
$$
\mathbb{R}^n = \text{span}\{\zeta\} \oplus \widetilde{X}_1. \tag{2.19}
$$

<span id="page-8-0"></span>**Proposition 2.6.** *Suppose that*  $q \in (0, na)$ *, let*  $u^{\lambda}$  *be the unique positive solution of* (2.16)*, and define*  $u^0 = (u_1^0, \cdots, u_1^0)^T$ , where

$$
u_j^0 = \frac{na - q}{nb} \text{ for all } j = 1, \cdots, n.
$$

*Then the following statements hold:*

<span id="page-8-7"></span>(*i*)  $u^{\lambda} = (u_1^{\lambda}, u_2^{\lambda}, \cdots, u_n^{\lambda})$  is continuously differentiable for  $\lambda \in [0, \lambda_q^*);$ <br>*ii*) *(ii)*

$$
\sum_{j=1}^{n} (u_j^{\lambda})' \big|_{\lambda=0} = -\frac{q^2(n+1)(n-1)}{6b} < 0,\tag{2.20}
$$

*and the total population size*  $\sum_{j=1}^{n} u_j^{\lambda}$  *is strictly decreasing in*  $\lambda \in (0, \epsilon)$  *with*  $\epsilon \ll 1$ *.* 

*Here, is the derivative with respect to* λ *and*

<span id="page-8-3"></span>
$$
\lambda_q^* = \begin{cases} 1/d_q^*, & \text{if } q \in (a, na) \\ \infty, & \text{if } q \in (0, a] \end{cases}
$$

*with d*<sup>∗</sup> *<sup>q</sup> defined in Proposition [2.5.](#page-5-8)*

**Proof.** (i) By Proposition [2.5,](#page-5-8) we see that  $u^{\lambda}$  is continuously differentiable for  $\lambda \in (0, \lambda^*_{\xi})$ . Then we need to show that  $u^{\lambda}$  is continuously differentiable for  $\lambda \in [0, \tilde{\lambda}_1]$  with  $0 < \tilde{\lambda}_1 \ll 1$ . From the decomposition in (2.10) we see that  $u = (u_1, \ldots, u_n)^T$  in (2.16) can be represented as follows: in [\(2.19\)](#page-8-2), we see that  $\mathbf{u} = (u_1, \dots, u_n)^T$  in [\(2.16\)](#page-8-1) can be represented as follows:

<span id="page-8-4"></span>
$$
\mathbf{u} = r\mathbf{g} + \mathbf{v} \text{ with } r = \sum_{j=1}^{n} u_j \in \mathbb{R} \text{ and } \mathbf{v} \in \widetilde{X}_1. \tag{2.21}
$$

Plugging  $(2.21)$  into  $(2.16)$ , we have

$$
\sum_{k=1}^{n} D_{jk} v_k + \lambda \left[ q \sum_{k=1}^{n} Q_{jk} (r \zeta_k + v_k) + (r \zeta_j + v_j) (a - b (r \zeta_j + v_j)) \right] = 0 \qquad (2.22)
$$

for  $j = 1, \dots, n$ . Denoting the left side of [\(2.22\)](#page-8-4) by  $\tilde{y}_i$ , we see from [\(2.19\)](#page-8-2) that

$$
\widetilde{\mathbf{y}} = (\widetilde{y}_1, \cdots, \widetilde{y}_n)^T = \widetilde{c}_{\mathbf{S}} + \widetilde{\mathbf{z}}
$$
 with  $\widetilde{c} = \sum_{j=1}^n y_j$  and  $\widetilde{\mathbf{z}} \in \widetilde{X}_1$ .

Therefore,  $\tilde{y} = 0$  if and only if  $\tilde{c} = 0$  and  $\tilde{z} = 0$ . This combined with [\(2.22\)](#page-8-4) implies that

$$
\widetilde{c} = \lambda \widetilde{G}_2
$$
 and  $\widetilde{z} = (\widetilde{G}_{1,1}, \ldots, \widetilde{G}_{1,n})^T$ ,

where

$$
\widetilde{G}_{1,j}(\lambda, r, \nu) = \sum_{k=1}^{n} D_{jk}v_k + \lambda \left[ q \sum_{k=1}^{n} Q_{jk} (r_{\mathcal{S}_k} + v_k) + a (r_{\mathcal{S}_j} + v_j) - b (r_{\mathcal{S}_j} + v_j)^2 \right] \n- \frac{\lambda}{n} \widetilde{G}_2(\lambda, r, \nu), \quad j = 1, \cdots, n, \n\widetilde{G}_2(\lambda, r, \nu) = \sum_{j=1}^{n} \left[ q \sum_{k=1}^{n} Q_{jk} (r_{\mathcal{S}_k} + v_k) + a (r_{\mathcal{S}_j} + v_j) - b (r_{\mathcal{S}_j} + v_j)^2 \right].
$$

Define  $\widetilde{G}(\lambda, r, v) : \mathbb{R}^2 \times \widetilde{X}_1 \to \widetilde{X}_1 \times \mathbb{R}$  by  $\widetilde{G} := (\widetilde{G}_{1,1}, \cdots, \widetilde{G}_{1,n}, G_2)^T$ . It follows that, for  $\lambda \in [0, \lambda_q^*)$ , *u*<br>(geographical in (2,21)) is a solution of (2,16) if and only if (*u*, *u*) (represented in [\(2.21\)](#page-8-3)) is a solution of [\(2.16\)](#page-8-1) if and only if  $(r, v) \in \mathbb{R} \times \widetilde{X}_1$  is a solution of  $\widetilde{G}(\lambda, r, v) = 0$ . Then using similar arguments as in the proof of Proposition [2.5,](#page-5-8) we can show that there exists  $\tilde{\lambda}_1 > 0$ and a continuously differentiable mapping  $\lambda \in [0, \tilde{\lambda}_1] \mapsto (r^{\lambda}, v^{\lambda}) \in \mathbb{R} \times \tilde{X}_1$  such that

$$
(r^{0}, v^{0}) = \left(\frac{na-q}{b}, 0\right) \text{ and } \widetilde{G}(\lambda, r^{\lambda}, v^{\lambda}) = 0 \text{ for } \lambda \in [0, \widetilde{\lambda}_{1}].
$$
 (2.23)

This combined with [\(2.21\)](#page-8-3) implies that, for  $\lambda \in (0, \lambda_1)$ ,

<span id="page-9-1"></span><span id="page-9-0"></span>
$$
\boldsymbol{u}^{\lambda} = \boldsymbol{r}^{\lambda} \boldsymbol{\varsigma} + \boldsymbol{v}^{\lambda} \text{ with } \boldsymbol{r}^{\lambda} = \sum_{j=1}^{n} u_j^{\lambda} \text{ and } \boldsymbol{v}^{\lambda} \in \widetilde{X}_1. \tag{2.24}
$$

Therefore, *u*<sup> $\lambda$ </sup> is continuously differentiable for  $\lambda \in [0, \tilde{\lambda}_1]$  if we defined  $u^0 = r^0 g$ .<br>(ii) Now we compute  $\sum_{n=1}^n (u^{\lambda})^{\lambda}$ . Differentiating (2.23) with respect to  $\lambda$ 

(ii) Now we compute  $\sum_{j=1}^{n} (u_j^{\lambda})' \big|_{\lambda=0}$ . Differentiating [\(2.23\)](#page-9-0) with respect to  $\lambda$  at  $\lambda = 0$  and noticing that  $v^0 = 0$ , we have

$$
\sum_{k=1}^{n} D_{jk} (v_k^{\lambda})' \Big|_{\lambda=0} + q \sum_{k=1}^{n} Q_{jk} r^0 \zeta_k + ar^0 \zeta_j - b (r^0 \zeta_j)^2 - \frac{1}{n} \widetilde{G}_2(0, r^0, v^0) = 0,
$$
  

$$
q \sum_{j=1}^{n} \sum_{k=1}^{n} Q_{jk} [(r^{\lambda})' \zeta_k + (v_k^{\lambda})'] \Big|_{\lambda=0} + \sum_{j=1}^{n} [a - 2b (r^0 \zeta_j + v_j^0)] [(r^{\lambda})' \zeta_j + (v_j^{\lambda})'] \Big|_{\lambda=0} = 0.
$$

Noting that  $r^0 = \frac{n a - q}{b}$ ,  $(v^{\lambda})' \in \widetilde{X}_1$  and  $\widetilde{G}_2(0, r^0, v^0) = 0$ , we have

$$
\begin{cases}\n\sum_{k=1}^{n} D_{1k}(\nu_k^{\lambda})' \big|_{\lambda=0} - \frac{na - q}{nb} \cdot q + \frac{na - q}{nb} \cdot \frac{q}{n} = 0, \\
\sum_{k=1}^{n} D_{jk}(\nu_k^{\lambda})' \big|_{\lambda=0} + \frac{na - q}{nb} \cdot \frac{q}{n} = 0, \ j = 2, \cdots, n, \\
\left( -a + \frac{q}{n} \right) (\nu^{\lambda})' \big|_{\lambda=0} - q(\nu_n^{\lambda})' \big|_{\lambda=0} = 0.\n\end{cases}
$$

By a tedious computation (see Proposition  $5.4$  in the appendix), we obtain that

$$
(v_n^{\lambda})'_{\lambda=0} = \frac{q(na-q)(n+1)(n-1)}{6nb}, \ \ (r^{\lambda})'_{\lambda=0} = -\frac{q^2(n+1)(n-1)}{6b}.
$$

This, combined with [\(2.24\)](#page-9-1), implies that

$$
\sum_{j=1}^n (u_j^{\lambda})' \big|_{\lambda=0} = (r^{\lambda})' \big|_{\lambda=0} = -\frac{q^2(n+1)(n-1)}{6b}.
$$

This completes the proof.

<span id="page-10-0"></span>**Remark 2.7.** *Note that*  $\lambda = 1/d$ *. Then the total population size*  $\sum_{j=1}^{n} u_{dj}$  *for* [\(2.2\)](#page-5-7) *is strictly increasing in*  $d \in [\delta, \infty)$  *with*  $\delta \gg 1$ .

## **3. Eigenvalue problem**

Linearising model  $(1.5)$  at  $u_d$ , we have

<span id="page-10-2"></span>
$$
\frac{dv}{dt} = dDv + qQv + \text{diag}(a - bu_{d,j})v - \text{diag}(bu_{d,j})v(t - \tau).
$$
\n(3.1)

It follows from [\[21,](#page-30-32) Chapter 7] that the infinitesimal generator  $A<sub>\tau</sub>(d)$  of the solution semigroup of [\(3.1\)](#page-10-2) is defined by:

$$
A_{\tau}(d)\Psi = \dot{\Psi} \tag{3.2}
$$

with the domain

<span id="page-10-7"></span><span id="page-10-6"></span>
$$
\mathscr{D}(A_{\tau}(d)) = \left\{ \Psi \in C^{1}([-\tau, 0], \mathbb{C}^{n}) : \dot{\Psi}(0) = dD\Psi(0) + qQ\Psi(0) + d\Psi(0) + \text{diag} (a - bu_{d,j}) \Psi(0) - \text{diag} (bu_{d,j}) \Psi(-\tau) \right\}.
$$
\n(3.3)

Then  $\mu \in \sigma_p(A_\tau(d))$  (resp.,  $\mu$  is an eigenvalue of  $A_\tau(d)$ ) if and only if there exists  $\psi = (\psi_1, \dots, \psi_n)^T (\neq$ **0**)  $\in \mathbb{C}^n$  such that  $\Delta(d, \mu, \tau)$  $\psi = 0$ , where matrix

$$
\Delta(d, \mu, \tau) := dD + qQ + \text{diag}(a - bu_{d,j}) - e^{-\mu \tau} \text{diag}(bu_{d,j}) - \mu I. \tag{3.4}
$$

To determine the distribution of the eigenvalues of  $A<sub>\tau</sub>(d)$ , one need to consider whether

<span id="page-10-3"></span>
$$
\sigma_p(A_\tau(d)) \cap \{x + iy : x = 0\} \neq \emptyset.
$$

By Proposition [2.4,](#page-4-4) we have

$$
0 \notin \sigma_p(A_\tau(d))
$$
 for all  $\tau \ge 0$ ,  
 $\sigma_p(A_\tau(d)) \subset \{x + iy : x < 0\}$  for  $\tau = 0$ .

In fact, if  $0 \in \sigma_p(A_{\tau_0}(d))$  for some  $\tau_0 \ge 0$ , then 0 is an eigenvalue of matrix

$$
dD + qQ + \text{diag}(a - 2bu_{d,j}),
$$

<span id="page-10-4"></span>which contradicts Proposition [2.4.](#page-4-4) By [\(3.4\)](#page-10-3), we see that  $i\nu(\nu > 0) \in \sigma_p(A_\tau(d))$  for some  $\tau > 0$  if and only if

<span id="page-10-5"></span>
$$
\begin{cases}\n\mathcal{M}(d, \nu, \theta)\mathbf{\psi} = \mathbf{0} \\
\nu > 0, \ \theta \in [0, 2\pi), \ \mathbf{\psi}(\neq \mathbf{0}) \in \mathbb{C}^n\n\end{cases}
$$
\n(3.5)

admits a solution  $(v, \theta, \psi)$ , where matrix

$$
\mathcal{M}(d, v, \theta) = dD + qQ + \text{diag}(a - bu_{d,j}) - e^{-i\theta} \text{diag}(bu_{d,j}) - i\nu I. \tag{3.6}
$$

It follows from Proposition [2.5](#page-5-8) that the properties of  $u_d$  are different for the following three cases (see Figure [3\)](#page-11-1):

> **Case I:**  $q \in (a, na)$  and  $0 < d - d_q^* \ll 1$ ; **Case II:**  $q \in (0, a)$  and  $0 < d \ll 1$ ;

**Case III:** 
$$
q \in (0, na)
$$
 and  $d \gg 1$ .

<span id="page-10-1"></span>Therefore, the following analysis on eigenvalue problem [\(3.5\)](#page-10-4) is divided into three cases.

<span id="page-11-1"></span>

*Figure 3. Diagram for parameter ranges of Cases I–III.*

## *3.1. A priori estimates*

In this subsection, we give *a priori* estimates for solutions of [\(3.5\)](#page-10-4).

<span id="page-11-0"></span>**Lemma 3.1.** Let  $(v^d, \theta^d, \psi^d)$  be a solution of [\(3.5\).](#page-10-4) Then the following statements hold: *(i)* For fixed  $q \in (a, na)$ , ν*d d* − *d*<sup>∗</sup> *q is bounded for*  $d \in (d_q^*, \tilde{d}_1]$  *with*  $0 < \tilde{d}_1 - d_q^* \ll 1$ ;

- *(ii)* For fixed  $q \in (0, a)$ ,  $|v^d|$  *is bounded for*  $d \in (0, \tilde{d}_2)$  *with*  $0 < \tilde{d}_2 \ll 1$ ;
- *(iii)* For fixed  $q \in (0, na)$ ,  $|v^d|$  *is bounded for*  $d \in (\tilde{d}_3, \infty)$  *with*  $\tilde{d}_3 \gg 1$ *.*

**Proof.** We only prove (i), and (ii)–(iii) can be proved similarly. Define matrix  $\rho := (\rho_{jk})$  with

$$
\varrho_{jk} = \begin{cases} \left(\frac{d}{d+q}\right)^{j-1}, & j = k = 1, \cdots, n, \\ 0, & \text{otherwise.} \end{cases}
$$

Substituting  $(v, \theta, \psi) = (v^d, \theta^d, \psi^d)$  into [\(3.5\)](#page-10-4), and multiplying both sides of (3.5) by  $(\overline{\psi_1^d}, \overline{\psi_2^d}, \ldots, \overline{\psi_n^d})$   $\varrho$  to the left, we have

$$
S + \sum_{k=1}^{n} \left[ (a - bu_{d,k}) - e^{-i\theta^d} bu_{d,k} - i\nu^d \right] \left( \frac{d}{d+q} \right)^{k-1} |\psi_k^d|^2 = 0, \tag{3.7}
$$

where

<span id="page-11-2"></span>
$$
S:=\left(\overline{\psi_1^d},\overline{\psi_2^d},\cdots,\overline{\psi_n^d}\right)\varrho(dD+qQ)\boldsymbol{\psi}^d.
$$

Since  $\rho(dD + qQ)$  is symmetric, we see that  $S \in \mathbb{R}$ . This combined with [\(3.7\)](#page-11-2) yields

$$
\nu^d \sum_{k=1}^n \left(\frac{d}{d+q}\right)^{k-1} |\psi_k^d|^2 = (\sin \theta^d) \sum_{k=1}^n bu_{d,k} \left(\frac{d}{d+q}\right)^{k-1} |\psi_k^d|^2.
$$
 (3.8)

By Proposition [2.5](#page-5-8) (ii), we see that there exists  $M > 0$  such that  $\frac{\|u_d\|_{\infty}}{d - d_q^*}$  $< M$  for  $d \in (d_q^*, \tilde{d}_1]$  with 0 <  $d_1 - d_q^* \ll 1$ . This combined with [\(3.8\)](#page-11-3) implies that

$$
\left|\frac{v^d}{d-d_q^*}\right| \le bM \text{ for } d \in (d_q^*, \tilde{d}_1].
$$

#### This completes the proof.

<https://doi.org/10.1017/S0956792524000342>Published online by Cambridge University Press

<span id="page-11-3"></span> $\Box$ 

#### 3.2. Case I *3.2. Case I*

For this case, the positive equilibrium  $u_d$  can be represented as [\(2.10\)](#page-5-5). Plugging (2.10) into [\(3.5\)](#page-10-4), we rewrite the eigenvalue problem [\(3.5\)](#page-10-4) as follows:

<span id="page-12-0"></span>
$$
\begin{cases} \sum_{k=1}^{n} \mathcal{L}_{jk} \psi_k + (d - d_q^*) f_j(\boldsymbol{\psi}, \theta, d) - i \nu \psi_j = 0, \ \ j = 1, \cdots, n, \\ \nu > 0, \ \theta \in [0, 2\pi), \ \boldsymbol{\psi}(\neq \mathbf{0}) \in \mathbb{C}^n, \end{cases}
$$
\n
$$
(3.9)
$$

where  $\mathcal L$  is defined in [\(2.3\)](#page-5-9), and

<span id="page-12-2"></span>
$$
f_j(\boldsymbol{\psi}, \theta, d) = \sum_{k=1}^n D_{jk} \psi_k - b(\alpha_d \eta_j + \xi_{dj}) \psi_j - e^{-i\theta} b(\alpha_d \eta_j + \xi_{dj}) \psi_j
$$
(3.10)

with  $\alpha_d$  and  $\xi_d$  defined in [\(2.10\)](#page-5-5). By [\(2.6\)](#page-5-4), we see that, ignoring a scalar factor,  $\psi$ ( $\neq$ **0**)  $\in$   $\mathbb{C}^n$  in [\(3.9\)](#page-12-0) can be represented as follows:

<span id="page-12-1"></span>
$$
\begin{cases} \n\boldsymbol{\psi} = \beta \boldsymbol{\eta} + z \text{ with } z \in (X_1)_{\mathbb{C}}, \ \beta \ge 0, \\
\|\boldsymbol{\psi}\|_2^2 = \beta^2 \|\boldsymbol{\eta}\|_2^2 + \beta \sum_{j=1}^n \eta_i (z_i + \overline{z_i}) + \|z\|_2^2 = \|\boldsymbol{\eta}\|_2^2,\n\end{cases} \tag{3.11}
$$

where  $\eta$  is defined in [\(2.4\)](#page-5-0). Then we obtain an equivalent problem of [\(3.9\)](#page-12-0) in the following.

**Lemma 3.2.** Assume that  $d > d_q^*$  and  $q \in (a, na)$ . Then  $(\psi, v, \theta)$  solves [\(3.9\),](#page-12-0) where  $\psi$  is defined in [\(3.11\)](#page-12-1) *and*  $v = (d - d_q^*)\omega$ , *if and only if*  $(\beta, \omega, \theta, z)$  *solves* 

<span id="page-12-4"></span>
$$
\begin{cases}\nF(\beta, \varpi, \theta, z, d) = 0, \\
\beta \ge 0, \varpi > 0, \ \theta \in [0, 2\pi), \ z \in (X_1)_{\mathbb{C}}.\n\end{cases}
$$
\n(3.12)

*Here*

<span id="page-12-5"></span>
$$
\bm{F}(\beta, \bm{\varpi}, \theta, \bm{z}, d) = (F_{1,1}, \cdots, F_{1,n}, F_2, F_3)^T
$$

*is a continuously differentiable mapping from*  $\mathbb{R}^3 \times (X_1)_\mathbb{C} \times [d_q^*, \infty]$  *to*  $(X_1)_\mathbb{C} \times \mathbb{C} \times \mathbb{R}$ *, and* 

$$
\begin{cases}\nF_{1,j} := \sum_{k=1}^{n} \mathcal{L}_{jk} z_k \\
+ (d - d_q^*) \left[ f_j(\beta \eta + z, \theta, d) - i \varpi \left( \beta \eta_j + z_j \right) - F_2(\beta, \varpi, \theta, z, d) \right], \ j = 1, \cdots, n, \\
F_2 := \sum_{j=1}^{n} \hat{\eta}_j \left[ f_j(\beta \eta + z, \theta, d) - i \varpi \left( \beta \eta_j + z_j \right) \right], \\
F_3 := (\beta^2 - 1) \|\eta\|_2^2 + \beta \sum_{i=j}^{n} \eta_j (z_j + \overline{z}_j) + \|z\|_2^2,\n\end{cases} \tag{3.13}
$$

*where*  $f_j$  ( $j = 1, \dots, n$ ) *are defined in* [\(3.10\).](#page-12-2)

**Proof.** Clearly,  $F_3 = 0$  is equivalent to second equation of [\(3.11\)](#page-12-1). Substituting (3.11) and  $v = (d - d_q^*)\varpi$ into  $(3.9)$ , we see that

<span id="page-12-3"></span>
$$
\sum_{k=1}^{n} \mathcal{L}_{jk} z_k + (d - d_q^*) \left[ f_j(\beta \eta + z, \theta, d) - i \varpi \left( \beta \eta_j + z_j \right) \right] = 0, \ \ j = 1, \cdots, n. \tag{3.14}
$$

Denote the left side of [\(3.14\)](#page-12-3) by *y<sub>j</sub>* and let  $\mathbf{y} = (y_1, \dots, y_n)^T$ . Using similar arguments as in the proof of Proposition 2.5 (ii) we see that  $\mathbf{v} = \mathbf{0}$  if and only  $F_s = 0$  and  $F_s = 0$  for all  $i = 1, \dots, n$ . This Proposition [2.5](#page-5-8) (ii), we see that  $y = 0$  if and only  $F_2 = 0$  and  $F_{1,j} = 0$  for all  $j = 1, \dots, n$ . This completes the proof. the proof.

We first show that the equivalent problem [\(3.12\)](#page-12-4) has a unique solution for  $d = d_q^*$ .

<span id="page-12-6"></span>**Lemma 3.3.** *The following equation*

$$
\begin{cases} F(\beta, \varpi, \theta, z, d_q^*) = 0 \\ \beta \ge 0, \ \varpi \ge 0, \ \theta \in [0, 2\pi], \ z \in (X_1)_{\mathbb{C}} \end{cases}
$$

#### 14 *W. Liu et al.*

*has a unique solution*  $(\beta^*, \varpi^*, \theta^*, z^*)$ *. Here* 

$$
z^* = 0, \ \beta^* = 1, \ \varpi^* = \frac{\sum_{j,k=1}^n D_{jk} \eta_k \hat{\eta}_j}{\sum_{j=1}^n \eta_j \hat{\eta}_j} > 0, \ \theta^* = \frac{\pi}{2}.
$$

**Proof.** Clearly,

$$
F_{1,j}(\beta, \varpi, \theta, z, d_q^*) = 0
$$
 for all  $j = 1, \dots, n$ 

if and only if  $z = z^* = 0$ . Substituting  $z = 0$  into  $F_3(\beta, \varpi, \theta, z, d^*_q) = 0$ , we have  $\beta = \beta^* = 1$ . Then plugging  $z = 0$  and  $\beta = 1$  into  $F_2(\beta, \omega, \theta, z, d_q^*) = 0$ , we have

<span id="page-13-0"></span>
$$
\sum_{j=1}^{n} \hat{\eta}_j \left[ f_j(\boldsymbol{\eta}, \theta, d_q^*) - \mathrm{i} \varpi \eta_j \right] = 0, \tag{3.15}
$$

where  $f_i$  ( $j = 1, \dots, n$ ) are defined in [\(3.10\)](#page-12-2). By Proposition [2.5](#page-5-8) (ii), we have

$$
f_j(\boldsymbol{\eta},\theta,d_q^*)=\sum_{k=1}^n D_{jk}\eta_k-b\alpha^*\eta_j^2-e^{-i\theta}b\alpha^*\eta_j^2,
$$

where  $\alpha^*$  is defined in [\(2.11\)](#page-5-6). Then we see from [\(3.15\)](#page-13-0) that

$$
\sum_{j,k=1}^n D_{jk} \eta_k \hat{\eta}_j - \sum_{j=1}^n b \alpha^* \hat{\eta}_j \eta_j^2 - e^{-i\theta} b \alpha^* \sum_{j=1}^n \hat{\eta}_j \eta_j^2 - i \varpi \sum_{j=1}^n \eta_j \hat{\eta}_j = 0.
$$

This combined with  $(2.11)$  yields

$$
\varpi = \varpi^* = \frac{\sum_{j,k=1}^n D_{jk} \eta_k \hat{\eta}_j}{\sum_{j=1}^n \eta_j \hat{\eta}_j} > 0, \ \ \theta = \theta^* = \frac{\pi}{2}.
$$

This completes the proof.

Then we solve the equivalent problem [\(3.12\)](#page-12-4) for  $0 < d - d_q^* \ll 1$ .

<span id="page-13-2"></span>**Lemma 3.4.** Assume that  $d > d_q^*$  and  $q \in (a, na)$ . Then there exists  $\tilde{d}_1$  with  $0 < \tilde{d}_1 - d_q^* \ll 1$  and a continuously differentiable mapping  $d \mapsto (\beta_d, \varpi_d, \theta_d, z_d)$  from  $[d_q^*, \tilde{d}_1]$  to  $\mathbb{R}^3 \times (X_1)_\mathbb{C}$  such that  $(\beta_d, \varpi_d, \theta_d, z_d)$ <br>is the unique solution of the following problem: *is the unique solution of the following problem:*

<span id="page-13-1"></span>
$$
\begin{cases}\nF(\beta, \varpi, \theta, z, d) = \mathbf{0} \\
\beta \ge 0, \varpi > 0, \ \theta \in [0, 2\pi), \ z \in (X_1)_{\mathbb{C}}\n\end{cases}
$$
\n(3.16)

 $\Box$ 

*for*  $d \in (d_q^*, \tilde{d}_1].$ 

**Proof.** Let  $P\left(\check{\beta}, \check{\varpi}, \check{\theta}, \check{z}\right) = (P_{1,1}, \cdots, P_{1,n}, P_2, P_3)^T : \mathbb{R}^3 \times (X_1)_\mathbb{C} \mapsto (X_1)_\mathbb{C} \times \mathbb{C} \times \mathbb{R}$  be the Fréchet derivative of *F* with respect to  $(\beta, \varpi, \theta, z)$  at  $(\beta^*, \varpi^*, \theta^*, z^*, d_q^*)$ . It follows from [\(3.12\)](#page-12-4) and [\(3.13\)](#page-12-5) that

$$
P_{1,j}\left(\check{\beta},\check{\omega},\check{\theta},\check{z}\right) = \sum_{k=1}^{n} \mathcal{L}_{jk}\check{z}_k, \ j = 1,\cdots,n,
$$
  
\n
$$
P_2\left(\check{\beta},\check{\omega},\check{\theta},\check{z}\right) = \sum_{j=1}^{n} \hat{\eta}_j \left(\sum_{k=1}^{n} D_{jk}\check{z}_k - b\alpha^* \eta_j \check{z}_j + ib\alpha^* \eta_j \check{z}_j - i\omega^* \check{z}_j + b\alpha^* \eta_j^2 \check{\theta} - i\eta_j \check{\omega}\right)
$$
  
\n
$$
+ \sum_{j=1}^{n} \hat{\eta}_j (ib\alpha^* \eta_j^2 \check{\beta} - i\omega^* \eta_j \check{\beta}),
$$
  
\n
$$
P_3\left(\check{\beta},\check{\omega},\check{\theta},\check{z}\right) = 2\|\eta\|_2^2 \check{\beta} + \sum_{j=1}^{n} \eta_j \left(\check{z}_j + \check{z}_j\right),
$$

where we have used  $(2.11)$  to obtain  $P_2$ . A direct computation implies that **P** is a bijection from  $\mathbb{R}^3 \times (X_1)_\mathbb{C}$  to  $(X_1)_\mathbb{C} \times \mathbb{C} \times \mathbb{R}$ . It follows from the implicit function theorem that there exists  $\hat{d} > d_q^*$ and a continuously differentiable mapping  $d \mapsto (\beta_d, \varpi_d, \theta_d, z_d)$  from  $\left[ d_q^*, \hat{d} \right]$  to  $\mathbb{R}^3 \times (X_1)_\mathbb{C}$  such that  $(\beta_d, \varpi_d, \theta_d, z_d)$  solves [\(3.16\)](#page-13-1).

Then we prove the uniqueness for  $d \in [d_q^*, \tilde{d}_1]$  with  $0 < \tilde{d}_1 - d_q^* \ll 1$ . We only need to verify that if  $\left(\beta^d, \varpi^d, \theta^d, z^d\right)$  satisfies [\(3.16\)](#page-13-1), then  $\lim_{d \to d_q^*} \left(\beta^d, \varpi^d, \theta^d, z^d\right) = (\beta^*, \varpi^*, \theta^*, z^*)$ , where  $\left(\beta^*, \varpi^*, \theta^*, z^*\right)$ is defined in Lemma [3.3.](#page-12-6) Since

$$
F_3(\beta^d, \varpi^d, \theta^d, z^d, d) = 0,
$$

<span id="page-14-0"></span>we have

$$
\|\beta^d \pmb{\eta} + \pmb{z}^d\|_2^2 = \|\pmb{\eta}\|_2^2. \tag{3.17}
$$

Note from Lemma [3.1](#page-11-0) (i) that  $\varpi^d$  is bounded for  $d \in [d^*_q, \tilde{d}_1]$ . This combined with [\(3.17\)](#page-14-0) and the first equation of [\(3.13\)](#page-12-5) implies that  $\lim_{d \to d_q^*} \mathcal{L}z = 0$ , and consequently,  $\lim_{d \to d_q^*} z^d = z^* = 0$ . By the third equation of [\(3.13\)](#page-12-5), we obtain that  $\lim_{d \to d_q^*} \beta^d = \beta^* = 1$ . Then, it follows from the second equation of (3.13) that  $\lim_{d \to d_q^*} \theta^d = \theta^*$  and  $\lim_{d \to d_q^*} \varpi^d = \varpi^*$ . This completes the proof.  $\Box$ 

By Lemma [3.4,](#page-13-2) we obtain the following result.

<span id="page-14-2"></span>**Theorem 3.5.** Assume that  $q \in (a, na)$ . Then for  $d \in (d_q^*, \tilde{d}_1]$  with  $0 < \tilde{d}_1 - d_q^* \ll 1$ ,  $(v, \tau, \psi)$  satisfies *the following equation:*

<span id="page-14-1"></span>
$$
\begin{cases} \Delta(d, i\nu, \tau)\psi = 0\\ \nu > 0, \ \tau \ge 0, \ \psi(\neq 0) \in \mathbb{C}^n \end{cases}
$$
 (3.18)

*if and only if*

$$
\nu = \nu_d = (d - d_q^*) \varpi_d, \; \; \mathbf{\psi} = c_1 \mathbf{\psi}_d, \; \; \tau = \tau_{d,l} = \frac{\theta_d + 2l\pi}{\nu_d}, \; \; l = 0, 1, 2, \cdots,
$$

where  $\psi_d = \beta_d \eta + z_d$ ,  $c_1 \in \mathbb{C}$  is a non-zero constant, and  $\beta_d$ ,  $\varpi_d$ ,  $\theta_d$  and  $z_d$  are defined in Lemma [3.4.](#page-13-2)

#### 3.3. Case II *3.3. Case II*

For this case, the positive equilibrium  $u_d$  is continuously differentiable for  $d \in [0, \infty)$ . We first solve the eigenvalue problem  $(3.5)$  for  $d = 0$ .

<span id="page-15-5"></span>**Lemma 3.6.** *Let*  $\mathcal{M}(d, v, \theta)$  *be defined in* (3.6)*, and let* 

<span id="page-15-1"></span>
$$
Ker(\mathcal{M}(0,\nu,\theta)) := \{ \boldsymbol{\psi} \in \mathbb{C}^n : \mathcal{M}(0,\nu,\theta)\boldsymbol{\psi} = \boldsymbol{0} \}.
$$

*Assume that*  $q \in (0, a)$ *. Then* 

$$
\{(\nu,\theta): \nu \ge 0, \ \theta \in [0,2\pi], \ \text{Ker}(\mathcal{M}(0,\nu,\theta)) \neq \{\theta\}\} = \left\{(\nu_p^0, \theta_p^0)\right\}_{p=1}^n, \tag{3.19}
$$

<span id="page-15-3"></span>*where*

$$
\theta_p^0 = \arccos \frac{a - q - bu_{0,p}}{bu_{0,p}}, \ \nu_p^0 = \sqrt{(bu_{0,p})^2 - (a - q - bu_{0,p})^2}
$$
(3.20)

<span id="page-15-2"></span>*with*

$$
\frac{\pi}{2} = \theta_1^0 < \ldots < \theta_n^0 < \pi, \ \ a - q = \nu_1^0 < \ldots < \nu_n^0. \tag{3.21}
$$

<span id="page-15-6"></span>Moreover, for each  $p = 1, 2, \cdots, n$ ,  $Ker(\mathcal{M}(0, v_p^0, \theta_p^0)) = \{c\psi_p^0 : c \in \mathbb{C}\},\$  where  $\psi_p^0 = (\psi_{p,1}^0, \cdots, \psi_{p,n}^0)^T$ , *and*

$$
\psi_{p,j}^0 = 0 \text{ for } 1 \le j \le p-1, \quad \psi_{p,p}^0 = 1,
$$
\n
$$
\psi_{p,j}^0 = (-1)^{j-p} \prod_{k=p+1}^j \frac{q}{h_k(\theta_p^0, v_p^0)} \text{ for } p+1 \le j \le n
$$
\n(3.22)

<span id="page-15-7"></span>*with*

$$
h_k(\theta, \nu) = \left(a - q - bu_{0,k}\right) - bu_{0,k}e^{-i\theta} - i\nu, \ k = 1, \cdots, n. \tag{3.23}
$$

**Proof.** Clearly,  $det[M(d, v, \theta)] = \prod_{p=1}^{n} h_p(\theta, v)$ . For each  $p = 1, \dots, n$ , we compute that

$$
\begin{cases} h_p(\theta, \nu) = 0 \\ \theta \in [0, 2\pi], \ \nu \ge 0 \end{cases}
$$

admits a unique solution  $(\nu_p^0, \theta_p^0)$ , which yields [\(3.19\)](#page-15-1) holds. By [\(2.8\)](#page-5-3) and [\(2.9\)](#page-5-10), we see that [\(3.21\)](#page-15-2) holds, and consequently,  $h_k(\theta_p^0, v_p^0) \neq 0$  for  $k \neq p$ , which implies that  $\psi_p^0$  is well defined for  $p = 1, \dots, n$ . A direct computation implies that  $\text{Ker}(\mathcal{M}(0, v_p^0, \theta_p^0)) = \{c\psi_p^0 : c \in \mathbb{C}\}\$  for  $p = 1, \dots, n$ . This completes the proof.  $\Box$ 

The following result explores further properties of  $(v_p^0, \theta_p^0)$   $(p = 1, \dots, n)$ .

<span id="page-15-0"></span>**Lemma 3.7.** Assume that  $q \in (0, a)$ , and let  $(v_p^0, \theta_p^0)$   $(p = 1, \dots, n)$  be defined in [\(3.20\).](#page-15-3) Then the *following statements hold:*

(*i*)  $\frac{\theta_1^0}{\nu_1^0}$  $> \frac{\theta_2^0}{v_2^0}$  $> \ldots > \frac{\theta_n^0}{0}$  $v_n^0$ *; (ii)* For all  $p = 1, \dots, n, \frac{\theta_p^0}{\theta_p^0}$ ν0 *p is strictly monotone increasing in*  $q \in (0, a)$  *and satisfies*  $\lim_{q \to 0}$  $\theta_p^0$  $v_p^0$  $=\frac{\pi}{2a}$ .

**Proof.** By  $(3.20)$ , we have

<span id="page-15-4"></span>
$$
\frac{\theta_p^0}{v_p^0} = \frac{\arccos \frac{a - q - bu_{0,p}}{bu_{0,p}}}{\sqrt{(bu_{0,p})^2 - (a - q - bu_{0,p})^2}}, \ p = 1, \cdots, n,
$$
\n(3.24)

where  $u_{0,p} \ge (a-q)/b$  is defined in [\(2.8\)](#page-5-3) and depends on *q* for  $p = 1, \dots, n$ . Then we denote  $u_{0,p}$  by  $u_{0,p}(q)$  throughout the proof. We define an auxiliary function

$$
f_1(q, x) = \frac{\arccos \frac{a - q - bx}{bx}}{\sqrt{(bx)^2 - (a - q - bx)^2}} \text{ with } x \ge \frac{a - q}{b},
$$

and consequently,

$$
\frac{\theta_p^0}{\nu_p^0} = f_1 \left( q, u_{0,p}(q) \right) \text{ for } p = 1, \cdots, n. \tag{3.25}
$$

Let

<span id="page-16-0"></span>
$$
A_1 = \arccos \frac{a - q - bx}{bx}
$$
 and  $B_1 = \sqrt{(bx)^2 - (a - q - bx)^2}$ .

Noticing that *q* ∈ (0, *a*), we compute that, for  $x \ge \frac{a-q}{b}$ ,

$$
\frac{\partial A_1}{\partial x} = \frac{a - q}{x\sqrt{(bx)^2 - (a - q - bx)^2}} > 0, \quad \frac{\partial B_1}{\partial x} = \frac{b(a - q)}{\sqrt{(bx)^2 - (a - q - bx)^2}} > 0,
$$

$$
\frac{\partial A_1}{\partial q} = \frac{1}{\sqrt{(bx)^2 - (a - q - bx)^2}} > 0, \quad \frac{\partial B_1}{\partial q} = \frac{a - q - bx}{\sqrt{(bx)^2 - (a - q - bx)^2}} \le 0,
$$

which yields

$$
\frac{\partial f_1}{\partial q} = \frac{1}{B_1^2} \left( B_1 \frac{\partial A_1}{\partial q} - A_1 \frac{\partial B_1}{\partial q} \right) > 0 \text{ for } x \ge \frac{a - q}{b}.
$$
 (3.26)

By  $x \ge (a - q)/b$  and  $q \in (0, a)$  again, we have

$$
0 < \frac{\sqrt{(bx)^2 - (a - q - bx)^2}}{bx} \le 1 \text{ and } \arccos \frac{a - q - bx}{bx} \ge \frac{\pi}{2},
$$

which yields

<span id="page-16-1"></span>
$$
\frac{\partial f_1}{\partial x} = \frac{1}{B_1^2} \left( B_1 \frac{\partial A_1}{\partial x} - A_1 \frac{\partial B_1}{\partial x} \right)
$$
  
= 
$$
\frac{b(a-q)}{B_1^3} \left[ \frac{\sqrt{(bx)^2 - (a-q-bx)^2}}{bx} - \arccos \frac{a-q-bx}{bx} \right] < 0
$$
(3.27)

for  $x \geq (a - q)/b$ .

Now we prove (i). By Proposition [2.5,](#page-5-8) we have  $u_{0,1} < \cdots < u_{0,n}$ . This combined with [\(3.25\)](#page-16-0) and [\(3.27\)](#page-16-1) implies that (i) holds. Then we consider (ii). We first show that  $u_{0,p}(q) < 0$  for  $q > 0$  and  $p = 1, \dots, n$ . Here,  $\prime$  is the derivative with respect to *q*. Differentiating [\(2.8\)](#page-5-3) with respect to *q* yields

$$
\begin{cases} u'_{0,1} = -\frac{1}{b}, \\ qu'_{0,j-1} + (u_{0,j-1} - u_{0,j}) = -u'_{0,j} (a - q - 2bu_{0,j}) \text{ for } j = 2, \cdots, n. \end{cases}
$$

This combined with the fact that

$$
u_{0,n}>\cdots>u_{0,1}\geq \frac{a-q}{b},
$$

yields  $u_{0,p}(q) < 0$  for  $q > 0$  and  $p = 1, \dots, n$ . Then, by [\(3.25\)](#page-16-0)–[\(3.27\)](#page-16-1), we obtain that, for each  $p =$  $1, \cdots, n$ ,

$$
\left(\frac{\theta_p^0}{\nu_p^0}\right)' = \left[f_1(q, u_{0,p}(q))\right]' = \frac{\partial f_1(q, x)}{\partial q}\bigg|_{x=u_{0,p}(q)} + \frac{\partial f_1(q, x)}{\partial x}\bigg|_{x=u_{0,p}(q)} \cdot \frac{\partial u_{0,p}(q)}{\partial q} > 0,
$$

where ' is the derivative with respect to *q*. Note that  $\lim_{q\to 0} u_{0,p}(q) = a/b$  for  $p = 1, \dots, n$ . Then, by [\(3.24\)](#page-15-4),

we have lim *q*→0  $\theta^0_p$  $v_p^0$  $=\frac{\pi}{2a}$  for  $p = 1, \dots, n$ . This completes the proof.

Then we consider the solutions of [\(3.5\)](#page-10-4) for  $d \neq 0$ .

<span id="page-17-7"></span>**Lemma 3.8.** *Let*  $\mathcal{M}(d, v, \theta)$  *be defined in*  $(3.6)$ *, and let* 

$$
Ker(\mathcal{M}(d,\nu,\theta)) := \{ \psi \in \mathbb{C}^n : \mathcal{M}(d,\nu,\theta)\psi = 0 \}.
$$

*Then there exists*  $\tilde{d}_2 > 0$  *such that, for*  $d \in (0, \tilde{d}_2]$ *,* 

$$
W := \{ (\nu, \theta) : \nu > 0, \ \theta \in [0, 2\pi), \ Ker(\mathcal{M}(d, \nu, \theta)) \neq \{0\} \} = \left\{ (\nu_p^d, \theta_p^d) \right\}_{p=1}^n, \tag{3.28}
$$

*where*  $(v_p^d, \theta_p^d) \in (0, \infty) \times (0, \pi)$  *for each*  $p = 1, \cdots, n$ *. Moreover, for each*  $p = 1, \cdots, n$ *,* 

<span id="page-17-6"></span>
$$
Ker(\mathcal{M}(d, v_p^d, \theta_p^d)) = \{c\psi_p^d : c \in \mathbb{C}\},\qquad(3.29)
$$

<span id="page-17-5"></span> $\Box$ 

where  $(v_p^d, \theta_p^d, \psi_p^d)$  satisfies  $\lim_{d\to 0} (v_p^d, \theta_p^d, \psi_p^d) = (v_p^0, \theta_p^0, \psi_p^0)$ , and  $(v_p^0, \theta_p^0, \psi_p^0)$  is defined in Lemma [3.6.](#page-15-5)

**Proof.** Step 1. We show the existence of  $\left\{ (v_p^d, \theta_p^d) \right\}_{p=1}^n$  such that  $\left\{ (v_p^d, \theta_p^d) \right\}_{p=1}^n \subset W$ .

We only consider the existence of  $(v_2^d, \theta_2^d)$ , and  $\{(v_p^d, \theta_p^d)\}_{p\neq 2}$  can be obtained similarly. Letting  $\mathcal{M}^H(d, v, \theta)$  be the conjugate transpose of  $\mathcal{M}(d, v, \theta)$ , we compute that

$$
\text{Ker}\left(\mathcal{M}^H(0,\nu_2^0,\theta_2^0)\right) := \{\boldsymbol{\psi} \in \mathbb{C}^n : \mathcal{M}^H(0,\nu_2^0,\theta_2^0)\boldsymbol{\psi} = \mathbf{0}\} = \left\{c\boldsymbol{\varphi}_2^0 : c \in \mathbb{C}\right\},\
$$

where  $\boldsymbol{\varphi}_2^0 = (\varphi_{2,1}^0, \cdots, \varphi_{2,n}^0)^T$  with

<span id="page-17-0"></span>
$$
\varphi_{2,1}^0 = -\frac{q}{\overline{h}_1 \left(\theta_2^0, \nu_2^0\right)}, \ \varphi_{2,2}^0 = 1, \ \varphi_{2,k}^0 = 0 \ \text{for } k = 3, \cdots, n,
$$
\n(3.30)

and  $\overline{h}_1$   $(\theta_2^0, v_2^0)$  is the conjugate of  $h_1$   $(\theta_2^0, v_2^0)$ . By Lemma [3.6,](#page-15-5) we see that

<span id="page-17-1"></span>
$$
\text{Ker}(\mathcal{M}(0, v_2^0, \theta_2^0)) = \left\{c\psi_2^0 : c \in \mathbb{C}\right\}.
$$

By  $(3.22)$  and  $(3.30)$ , we see that

$$
\langle \boldsymbol{\varphi}_2^0, \boldsymbol{\psi}_2^0 \rangle = 1, \tag{3.31}
$$

and consequently

<span id="page-17-3"></span>
$$
\mathbb{C}^n = \text{Ker}(\mathcal{M}(0, \nu_2^0, \theta_2^0)) \oplus Z,
$$
\n(3.32)

where

<span id="page-17-2"></span> $Z := \{ z = (z_1, \dots, z_n)^T \in \mathbb{C}^n : \langle \varphi_2^0, z \rangle = 0 \}.$ 

Then by [\(3.30\)](#page-17-0), [\(3.31\)](#page-17-1) and [\(3.32\)](#page-17-2), we see that, for any  $y \in \mathbb{C}^n$ ,

$$
\mathbf{y} = \widehat{c}\mathbf{\psi}_2^0 + z \text{ with } \widehat{c} = \langle \mathbf{\varphi}_2^0, \mathbf{y} \rangle = \overline{\varphi}_{2,1}^0 y_1 + y_2 \text{ and } z \in \mathbb{Z}, \tag{3.33}
$$

where  $\overline{\varphi}_{2,1}^0$  is the conjugate of  $\varphi_{2,1}^0$ .

<span id="page-17-4"></span>Let

$$
\boldsymbol{H}\left(d,\nu,\theta,z\right)=(H_1,\cdots,H_n)^T:=\mathcal{M}\left(d,\nu,\theta\right)\left(\boldsymbol{\psi}_2^0+z\right): \mathbb{R}^3\times Z\to\mathbb{C}^n,\tag{3.34}
$$

where matrix  $M(d, v, \theta)$  is defined in [\(3.6\)](#page-10-5). By Lemma [3.6,](#page-15-5) we have  $H(0, v_2^0, \theta_2^0, \mathbf{0}) = \mathbf{0}$ , and then<br>we solve  $H(d, v, \theta, z) = \mathbf{0}$  for  $d \neq 0$ . By (3.32) and (3.33), we see that  $H(d, v, \theta, z) = \mathbf{0}$  if and only if we solve  $H(d, v, \theta, z) = 0$  for  $d \neq 0$ . By [\(3.32\)](#page-17-2) and [\(3.33\)](#page-17-3), we see that  $H(d, v, \theta, z) = 0$  if and only if  $H(d, v, \theta, z) = 0$ , where

$$
\widetilde{H} (d, v, \theta, \mathbf{0}) = (\widetilde{H}_0, \widetilde{H}_{1,1}, \cdots, \widetilde{H}_{1,n})^T : \mathbb{R}^3 \times Z \to \mathbb{C} \times Z,
$$

<span id="page-18-0"></span>and

$$
\widetilde{H}_0(d, v, \theta, z) = \langle \varphi_2^0, H(d, v, \theta, z) \rangle = \overline{\varphi_{2,1}^0} H_1(d, v, \theta, z) + H_2(d, v, \theta, z),
$$
\n
$$
\widetilde{H}_{1,j}(d, v, \theta, z) = H_j(d, v, \theta, z) - \widetilde{H}_0(d, v, \theta, z) \psi_{2,j}^0, \ j = 1, \cdots, n.
$$
\n(3.35)

Let

$$
\boldsymbol{P}\left(\check{\nu},\check{\theta},\check{z}\right)=(P_0,P_{1,1},\cdots,P_{1,n})^T:\mathbb{R}^2\times Z\to\mathbb{C}\times Z
$$

be the Fréchet derivative of  $\hat{H}$  with respect to  $(v, \theta, z)$  at  $(0, v_2^0, \theta_2^0, \mathbf{0})$ . By [\(3.23\)](#page-15-7) and the first equation of (3.35), we have of  $(3.35)$ , we have

$$
P_0(\check{\nu}, \check{\theta}, \check{z}) = (\overline{\varphi}_{2,1}^0 h_1(\theta_2^0, \nu_2^0) + q) \check{z}_1 + i e^{-i\theta_2^0} b u_{0,2} \check{\theta} - i \check{\nu}
$$
  
=  $i e^{-i\theta_2^0} b u_{0,2} \check{\theta} - i \check{\nu},$ 

where we have used  $(3.30)$  in the last step. By  $(3.23)$  and the second equation of  $(3.35)$ , we have

$$
(P_{1,1},\dots,P_{1,n})^T(\check{\nu},\check{\theta},\check{z})
$$
  
= $\mathcal{M}(0,\theta_2^0,\nu_2^0)\check{z}$   
+ $(0,0, (ie^{-i\theta_2^0}bu_{0,3}\check{\theta}-ie^{-i\theta_2^0}bu_{0,2}\check{\theta})\psi_{2,3}^0,\dots,(ie^{-i\theta_2^0}bu_{0,n}\check{\theta}-ie^{-i\theta_2^0}bu_{0,2}\check{\theta})\psi_{2,n}^0)^T$ .

Since  $\mathcal{M}(0, \theta_2^0, \nu_2^0)$  is a bijection from *Z* to *Z*, it follows that *P* is a bijection. Then we see from the implicit function theorem that there exists a constant  $\tilde{d} > 0$  a paighbourhood *N* of  $(\mu_0^0, \theta_$ implicit function theorem that there exists a constant  $\tilde{d} > 0$ , a neighbourhood  $N_2$  of  $(v_2^0, \theta_2^0, \mathbf{0})$  and a continuously differentiable function:

$$
(\nu_2^d, \theta_2^d, z_2^d) : [0, \tilde{d}) \mapsto N_2
$$

such that for any  $d \in [0, \tilde{d})$ ,  $(v_2^d, \theta_2^d, z_2^d)$  is the unique solution of  $\tilde{H}$   $(d, v, \theta, z) = 0$  in  $N_2$ . Therefore,  $(d, d, d, d, d)$  is also the unique solution of  $H$  (*d*,  $v, \theta$ ,  $a$ ) and  $f$  for  $d \in [0, \tilde{d})$ .  $(v_1^d, \theta_2^d, z_2^d)$  is also the unique solution of *H* (*d*, *v*,  $\theta$ , *z*) = **0** in *N*<sub>2</sub> for *d* ∈ [0,  $\tilde{d}$ ). This combined with [\(3.34\)](#page-17-4) implies that implies that

$$
\text{span}(\boldsymbol{\psi}_2^d) \subset \text{Ker}(\mathcal{M}(d, v_2^d, \theta_2^d)) \text{ for } d \in [0, \tilde{d}),\tag{3.36}
$$

<span id="page-18-1"></span>where  $\psi_2^d = \psi_2^0 + z_2^d$ . Note that the rank of  $\mathcal{M}(d, v_2^d, \theta_2^d)$  is  $n - 1$ . This combined with [\(3.36\)](#page-18-1) implies that (3.36) implies that  $(3.20)$  holds. Therefore, we show the existence of  $(v_1^d, \theta_1^d)$ . Moreove [\(3.29\)](#page-17-5) holds. Therefore, we show the existence of  $(v_2^d, \theta_2^d)$ . Moreover,  $\lim_{d\to 0} (v_2^d, \theta_2^d, \psi_2^d) = (v_2^0, \theta_2^0, \psi_2^0)$ . This combined with [\(3.21\)](#page-15-2) implies that  $(v_2^d, \theta_2^d) \in (0, \infty) \times (0, \pi)$ .

**Step 2.** We show that there exists  $\tilde{d}_2$  such that [\(3.28\)](#page-17-6) holds for  $d \in (0, \tilde{d}_2]$ .

If it is not true, then there exist sequences  $\left\{d_j\right\}_{j=1}^{\infty}$  and  $\left\{\left(v^{d_j}, \theta^{d_j}, \dot{\psi}^{d_j}\right)\right\}_{j=1}^{\infty}$  such that  $\lim_{j\to\infty}d_j=0$ , and for each *j*,

$$
\left(v^{d_j}, \theta^{d_j}\right) \notin \left\{\left(v_p^{d_j}, \theta_p^{d_j}\right)\right\}_{p=1}^n, \ \ \|\psi^{d_j}\|_2 = 1, \ \ v^{d_j} > 0, \ \ \theta^{d_j} \in [0, 2\pi),
$$

<span id="page-18-2"></span>and

$$
\mathcal{M}\left(d_{j},\nu^{d_{j}},\theta^{d_{j}}\right)\boldsymbol{\psi}^{d_{j}}=\mathbf{0}.\tag{3.37}
$$

It follows from Lemma [3.1](#page-11-0) (ii) that  $v^{d_j}$  is bounded. Then, up to a subsequence, we have

$$
\lim_{j\to\infty}\theta^{d_j}=\theta^*, \ \lim_{j\to\infty}\nu^{d_j}=\nu^*, \ \lim_{j\to\infty}\boldsymbol{\psi}^{d_j}=\boldsymbol{\psi}^*
$$

with  $\theta^* \in [0, 2\pi]$ ,  $\nu^* \ge 0$  and  $\|\psi^*\|_2 = 1$ . Taking  $j \to \infty$  on both sides of [\(3.37\)](#page-18-2), we have *M*(0,  $v^*$ ,  $\theta^*$ )  $\psi^* = 0$ . This combined with Lemma [3.6](#page-15-5) implies that there exists  $1 \le p_0 \le n$  and a constant  $c_{p_0} \neq 0$  such that  $v^* = v_{p_0}^0$  and  $\theta^* = \theta_{p_0}^0$ ,  $\boldsymbol{\psi}^* = c_{p_0} \boldsymbol{\psi}_{p_0}^0$ . Without loss of generality, we assume that  $p_0 = 2$ . Then, for sufficiently large *j*,

$$
\left(\nu^{d_j}, \theta^{d_j}, \frac{1}{c_2}\boldsymbol{\psi}^{d_j}-\boldsymbol{\psi}_2^0\right)\in N_2,
$$

where  $N_2$  (defined in step 1) is a neighbourhood of  $(v_2^0, \theta_2^0, \mathbf{0})$ . By [\(3.34\)](#page-17-4) and [\(3.37\)](#page-18-2), we see that

$$
\boldsymbol{H}\left(d_j,\nu^{d_j},\theta^{d_j},\frac{1}{c_2}\boldsymbol{\psi}^{d_j}-\boldsymbol{\psi}_2^0\right)=\mathbf{0}.
$$

Note from the proof of step 1 that  $(v_2^d, \theta_2^d, z_2^d)$  is the unique solution of *H*  $(d, v, \theta, z) = 0$  in  $N_2$  for  $d \in$  $[0, \tilde{d})$ . This implies that, for sufficiently large *j*,

$$
\left(\nu^{d_j}, \theta^{d_j}\right) = \left(\nu^{d_j}_2, \theta^{d_j}_2\right),
$$

which is a contradiction. Therefore,  $(3.28)$  holds.

By Lemmas [3.7](#page-15-0) and [3.8,](#page-17-7) we obtain the following result.

<span id="page-19-3"></span>**Theorem 3.9.** Suppose that  $q \in (0, a)$ . Then for  $d \in (0, \tilde{d}_2]$  with  $0 < \tilde{d}_2 \ll 1$ ,  $(v, \tau, \psi)$  solves  $(3.18)$  if and *only if there exists*  $1 \leq p \leq n$  *such that* 

$$
v = v_p^d
$$
,  $\mathbf{\psi} = c_2 \mathbf{\psi}_p^d$ ,  $\tau = \tau_{p,l}^d = \frac{\theta_p^d + 2l\pi}{v_p^d}$ ,  $l = 0, 1, 2, \cdots$ ,

*where*  $c_2 \in \mathbb{C}$  *is a non-zero constant, and*  $(v_p^d, \theta_p^d, \psi_p^d)$  *is defined in Lemma* [3.8.](#page-17-7) Moreover,

$$
\tau_{1,0}^d > \tau_{2,0}^d > \cdots > \tau_{n,0}^d.
$$
\n(3.38)

**Proof.** We only need to show that  $(3.38)$  holds, and other conclusions are direct results of Lemma [3.8.](#page-17-7) By Lemma [3.7](#page-15-0) (i), we have

$$
\lim_{d \to 0} \tau_{1,0}^d > \lim_{d \to 0} \tau_{2,0}^d > \cdots > \lim_{d \to 0} \tau_{n,0}^d,
$$
\n(3.39)

which implies that [\(3.38\)](#page-19-0) holds for  $d \in (0, \tilde{d}_2]$  with  $0 < \tilde{d}_2 \ll 1$ . This completes the proof.  $\Box$ 

### 3.4. Case III

For this case, we also need to find the equivalent problem of [\(3.5\)](#page-10-4). Throughout this subsection, we let  $\lambda = 1/d$  and denote  $u_d$  by  $u^{\lambda}$ . Then we rewrite [\(3.5\)](#page-10-4) as follows:

$$
\begin{cases}\nD\mathbf{\psi} + \lambda \left[ qQ + \text{diag} \left( a - bu_j^{\lambda} \right) - e^{-i\theta} \text{diag} \left( bu_j^{\lambda} \right) - i\nu I \right] \mathbf{\psi} = \mathbf{0}, \\
\nu > 0, \ \theta \in [0, 2\pi), \ \mathbf{\psi}(\neq \mathbf{0}) \in \mathbb{C}^n.\n\end{cases} \tag{3.40}
$$

By  $(2.17)$  and  $(2.18)$ , we see that

<span id="page-19-1"></span>
$$
\mathbb{C}^n = \text{span}\{\boldsymbol{\zeta}\} \oplus (\widetilde{X}_1)_{\mathbb{C}},
$$

where *ς* and  $\widetilde{X}_1$  are defined in [\(2.17\)](#page-8-5) and [\(2.18\)](#page-8-6), respectively. Ignoring a scalar factor,  $\psi \neq 0$ )  $\in \mathbb{C}^n$  in [\(3.40\)](#page-19-1) can be represented as follows:

<span id="page-19-2"></span>
$$
\begin{cases} \n\psi = \gamma \, \mathsf{s} + \mathsf{z}, \ \mathsf{z} \in (\widetilde{X}_1)_{\mathbb{C}}, \ \gamma \ge 0, \\
\|\psi\|_2^2 = \gamma^2 \|\mathsf{s}\|_2^2 + \|\mathsf{z}\|_2^2 = \|\mathsf{s}\|_2^2.\n\end{cases} \tag{3.41}
$$

Then we obtain the equivalent problem of [\(3.40\)](#page-19-1) in the following.

<span id="page-19-0"></span> $\Box$ 

**Lemma 3.10.** Assume that  $\lambda \in (0, \lambda_q^*)$  and  $q \in (0, na)$ . Then  $(\psi, v, \theta)$  solves [\(3.40\),](#page-19-1) where  $\psi$  is defined  $in (3.41)$ *, if and only if* ( $\gamma$ ,  $\nu$ ,  $\theta$ , *z*) *solves* 

<span id="page-20-3"></span>
$$
\begin{cases}\n\widetilde{F}(\gamma, \nu, \theta, z, \lambda) = \mathbf{0}, \\
\gamma \ge 0, \ \nu > 0, \ \theta \in [0, 2\pi), \ z \in (\widetilde{X}_1)_{\mathbb{C}}.\n\end{cases}
$$
\n(3.42)

*Here*

<span id="page-20-1"></span>
$$
\widetilde{F}(\gamma,\nu,\theta,z,\lambda)=(\widetilde{F}_{1,1},\cdots,\widetilde{F}_{1,n},\widetilde{F}_2,\widetilde{F}_3)^T
$$

*is a continuously differentiable mapping from*  $\mathbb{R}^3 \times (\widetilde{X}_1)_\mathbb{C} \times [0, \lambda_q^*)$  to  $(\widetilde{X}_1)_\mathbb{C} \times \mathbb{C} \times \mathbb{R}$ , and

$$
\begin{cases}\n\widetilde{F}_{1,j}(\gamma,\nu,\theta,z,\lambda) := \sum_{k=1}^{n} D_{jk}z_k + \lambda \left[ g_j(\gamma \varsigma + z,\theta) - i\nu \left( \gamma \varsigma_j + z_j \right) - \frac{1}{n} \widetilde{F}_2(\gamma,\nu,\theta,z,\lambda) \right], \\
\widetilde{F}_2(\gamma,\nu,\theta,z,\lambda) := \sum_{j=1}^{n} \left[ g_j(\gamma \varsigma + z,\theta) - i\nu \left( \gamma \varsigma_j + z_j \right) \right], \\
\widetilde{F}_3(\gamma,\nu,\theta,z,\lambda) := (\gamma^2 - 1) \| \varsigma \|_2^2 + \| z \|_2^2,\n\end{cases} \tag{3.43}
$$

*where*

$$
g_j(\gamma \boldsymbol{\varsigma} + \boldsymbol{z}, \theta) = q \sum_{k=1}^n Q_{jk} \left( \gamma \varsigma_k + z_k \right) + (a - bu_j^{\lambda}) \left( \gamma \varsigma_j + z_j \right) - e^{-i\theta} bu_j^{\lambda} \left( \gamma \varsigma_j + z_j \right).
$$

**Proof.** Plugging  $(3.41)$  into  $(3.40)$ , we see that

<span id="page-20-0"></span>
$$
\sum_{k=1}^{n} D_{jk} z_k + \lambda \left[ g_j(\gamma \mathbf{S} + \mathbf{z}, \theta) - i \nu \left( \gamma \mathbf{S}_j + z_j \right) \right] = 0, \ \ j = 1, \cdots, n. \tag{3.44}
$$

Denote the left side of [\(3.44\)](#page-20-0) by  $\tilde{y}_j$  and let  $\tilde{y} = (\tilde{y}_1, \dots, \tilde{y}_n)^T$ . Using similar arguments as in the proof of Proposition 2.5 (ii) we see that  $\tilde{y} = 0$  if and only  $\tilde{F}_i = 0$  and  $\tilde{F}_i = 0$  for all Proposition [2.5](#page-5-8) (ii), we see that  $\tilde{y} = 0$  if and only  $\tilde{F}_2 = 0$  and  $\tilde{F}_{1,j} = 0$  for all  $j = 1, \dots, n$ . This completes the proof. the proof.

We first solve the equivalent problem  $\widetilde{F}(\gamma, \nu, \theta, z, \lambda) = 0$  for  $\lambda = 0$ .

<span id="page-20-2"></span>**Lemma 3.11.** *The following equation*

$$
\begin{cases} \widetilde{F}(\gamma, \nu, \theta, z, 0) = 0 \\ \gamma \ge 0, \ \nu \ge 0, \ \theta \in [0, 2\pi], \ z \in (\widetilde{X}_1)_{\mathbb{C}} \end{cases}
$$

*has a unique solution*  $(\gamma_*, \nu_*, \theta_*, z_*)$ *, where* 

$$
z_* = 0
$$
,  $\gamma_* = 1$ ,  $\nu_* = a - \frac{q}{n}$ ,  $\theta_* = \frac{\pi}{2}$ .

**Proof.** Clearly,

$$
\big(\widetilde{F}_{1,1}(\gamma,\nu,\theta,z,0),\cdots,\widetilde{F}_{1,n}(\gamma,\nu,\theta,z,0)\big)=\mathbf{0},
$$

if and only if  $z = z_* = 0$ . Plugging  $z = 0$  into  $\widetilde{F}_3(\gamma, \nu, \theta, z, 0) = 0$ , we have  $\gamma = \gamma_* = 1$ . Note from Proposition [2.6](#page-8-0) that  $u_j^0 = \frac{na - q}{nb}$  for  $j = 1, \dots, n$ . Then plugging  $z = 0$  and  $\gamma = 1$  into  $\widetilde{F}_2(\gamma, \nu, \theta, z, 0) =$ 0, we have

$$
\sum_{j=1}^{n} \left[ q \sum_{k=1}^{n} Q_{jk} \zeta_k + \frac{q}{n} \zeta_j - e^{-i\theta} \zeta_j \frac{na - q}{n} \right] - i\nu = 0,
$$

#### 22 *W. Liu et al.*

which yields

$$
v = v_* = a - \frac{q}{n}, \ \theta = \theta_* = \frac{\pi}{2}.
$$

This completes the proof.

Now we solve  $\mathbf{F}(\gamma, \nu, \theta, z, \lambda) = \mathbf{0}$  for  $0 < \lambda \ll 1$ .

<span id="page-21-3"></span><span id="page-21-1"></span>**Lemma 3.12.** *There exists*  $\tilde{\lambda}$  with  $0 < \tilde{\lambda} \ll 1$  *and a continuously differentiable mapping*  $\lambda \mapsto$  $(\gamma^{\lambda}, v^{\lambda}, \theta^{\lambda}, z^{\lambda})$  *from* [0,  $\tilde{\lambda}$ ] *to*  $\mathbb{R}^{3} \times (\tilde{X}_{1})_{\mathbb{C}}$  *such that*  $(\gamma^{\lambda}, v^{\lambda}, \theta^{\lambda}, z^{\lambda})$  *is the unique solution of the following problem*: *problem:*

$$
\begin{cases} \widetilde{F}(\gamma, \nu, \theta, z, \lambda) = 0 \\ \gamma \ge 0, \ \nu > 0, \ \theta \in [0, 2\pi), \ z \in (\widetilde{X}_1)_{\mathbb{C}} \end{cases}
$$
 (3.45)

*for*  $\lambda \in [0, \tilde{\lambda}]$ *.* 

**Proof.** Let  $\widetilde{P}(\check{\gamma}, \check{\nu}, \check{\theta}, \check{z}) = (\widetilde{P}_{1,1}, \cdots, \widetilde{P}_{1,n}, \widetilde{P}_2, \widetilde{P}_3)^T : \mathbb{R}^3 \times (\widetilde{X}_1)_\mathbb{C} \to (\widetilde{X}_1)_\mathbb{C} \times \mathbb{C} \times \mathbb{R}$  be the Fréchet derivative of  $\widetilde{F}$  with respect to  $(\gamma, \nu, \theta, z)$  at  $(\gamma_*, \nu_*, \theta_*, z_*, 0)$ . Then we compute that

$$
\widetilde{P}_{1,j}\left(\check{\gamma},\check{\nu},\check{\theta},\check{z}\right) = \sum_{k=1}^{n} D_{jk}\check{z}_k,
$$
\n
$$
\widetilde{P}_2\left(\check{\gamma},\check{\nu},\check{\theta},\check{z}\right) = -q\check{z}_n + \left(a - \frac{q}{n}\right)\check{\theta} - i\check{\nu},
$$
\n
$$
\widetilde{P}_3\left(\check{\gamma},\check{\nu},\check{\theta},\check{z}\right) = 2\|\mathbf{S}\|_2^2\check{\gamma},
$$

where we have used  $\sum_{i=1}^{n} \check{\zeta}_i = 0$  to obtain  $\widetilde{P}_2$ . Clearly,  $\widetilde{P}$  is a bijection from  $\mathbb{R}^3 \times (\widetilde{X}_1)_C$  to  $\mathbb{R} \times \mathbb{C} \times (\widetilde{X}_1)_C$ . It follows from the implicit function theorem that there exists  $\tilde{\lambda} > 0$  with  $0 < \tilde{\lambda} \ll 1$  and a continuously differentiable mapping  $\lambda \mapsto (\gamma^{\lambda}, \nu^{\lambda}, \theta^{\lambda}, z^{\lambda})$  from [0,  $\tilde{\lambda}$ ] to  $\mathbb{R}^3 \times (\tilde{X}_1)_{\mathbb{C}}$  such that  $(\gamma^{\lambda}, \nu^{\lambda}, \theta^{\lambda}, z^{\lambda})$  satisfies  $(3.45).$  $(3.45).$ 

Then we prove the uniqueness of the solution of [\(3.45\)](#page-21-1). Actually, we only need to verify that if  $(\gamma_{\lambda}, \nu_{\lambda}, \theta_{\lambda}, z_{\lambda})$  satisfies [\(3.45\)](#page-21-1), then  $\lim_{\lambda \to 0} (\gamma_{\lambda}, \nu_{\lambda}, \theta_{\lambda}, z_{\lambda}) = (\gamma_*, \nu_*, \theta_*, z_*)$ . Since  $\tilde{F}_3$  ( $\gamma_{\lambda}, \nu_{\lambda}, \theta_{\lambda}, z_{\lambda}, \lambda$ ) = 0, we see that

$$
\|\gamma_{\lambda}\mathbf{S} + \mathbf{Z}_{\lambda}\|_{2}^{2} = \|\mathbf{S}\|_{2}^{2},\tag{3.46}
$$

<span id="page-21-2"></span>which implies that  $\gamma_{\lambda}$  is bounded for  $\lambda \in (0, \tilde{\lambda})$ . Note from Lemma [3.1](#page-11-0) (iii) that  $\nu_{\lambda}$  is bounded for  $\lambda \in (0, \tilde{\lambda}]$ . This combined with [\(3.46\)](#page-21-2) and the first equation of [\(3.43\)](#page-20-1) implies that  $\lim_{\lambda \to 0} Dz_{\lambda} = 0$ , and consequently,  $\lim_{\lambda \to 0} z_{\lambda} = 0$ . By the third equation of [\(3.43\)](#page-20-1), we obtain that  $\lim_{\lambda \to 0} \gamma_{\lambda} = 1$ . Then, it follows the second equation of [\(3.43\)](#page-20-1) that  $\lim_{\lambda \to 0} \theta_{\lambda} = \theta_{*}$  and  $\lim_{\lambda \to 0} \nu_{\lambda} = \nu_{*}$ . Therefore,  $(\beta_{\lambda}, \nu_{\lambda}, \theta_{\lambda}, z_{\lambda}) \to (\gamma_{*}, \nu_{*}, \theta_{*}, z_{*})$  as  $\lambda \rightarrow 0$ . This completes the proof.

By Lemma  $3.12$ , we obtain the following result.

<span id="page-21-4"></span>**Theorem 3.13.** Assume that  $q \in (0, na)$  and let  $\lambda = 1/d$ . Then for  $d \in (\tilde{d}_3, \infty)$  with  $\tilde{d}_3 \gg 1$ ,  $(\nu, \tau, \psi)$ *satisfies [\(3.18\)](#page-14-1) if and only if*

$$
\nu=\nu^{\lambda}, \ \ \boldsymbol{\psi}=c_3\boldsymbol{\psi}^{\lambda}, \ \ \tau=\tau_l^{\lambda}=\frac{\theta^{\lambda}+2l\pi}{\nu_{\lambda}}, \ \ l=0, 1, 2, \cdots,
$$

<span id="page-21-0"></span>*where*  $\mathbf{\psi}^{\lambda} = \gamma^{\lambda} \mathbf{g} + \mathbf{z}^{\lambda}$ ,  $c_3 \in \mathbb{C}$  *is a non-zero constant, and*  $(\gamma^{\lambda}, v^{\lambda}, \theta^{\lambda}, \mathbf{z}^{\lambda})$  *is defined in Lemma [3.12.](#page-21-3)* 

 $\Box$ 

#### **4. Local dynamics**

In this section, we obtain the local dynamics of model [\(1.5\)](#page-1-2) and show the existence of Hopf bifurcations. We first show that the purely imaginary eigenvalues obtained in Theorems [3.5,](#page-14-2) [3.9](#page-19-3) and [3.13](#page-21-4) are simple.

## <span id="page-22-3"></span>**Lemma 4.1.** *Let*  $\lambda = 1/d$ *, and define*

$$
(\widehat{\tau}_*, \widehat{\nu}_*) = \begin{cases} (\tau_{d,l}, \nu_d), & \text{if } a < q < na \text{ and } d \in (d^*_q, \widetilde{d}_1] \\ (\tau^d_{p,l}, \nu^d_p), & \text{if } 0 < q < a \text{ and } d \in (0, \widetilde{d}_2] \\ (\tau^{\lambda}_l, \nu^{\lambda}), & \text{if } 0 < q < na \text{ and } d \in [\widetilde{d}_3, \infty) \end{cases}
$$

with  $p=1,\dots, n$  and  $l=1,2,\dots,$  where  $(\tau_{d,l},\nu_d,\tilde{d}_1), (\tau_{p,l}^d,\nu_p^d,\tilde{d}_2)$  and  $(\tau_l^{\lambda},\nu^{\lambda},\tilde{d}_3)$  are defined in *Theorems* [3.5,](#page-14-2) [3.9](#page-19-3) *and* [3.13,](#page-21-4) *respectively. Then,*  $\mu = i\hat{v}_*$  *is an algebraically simple eigenvalue of*  $A_{\hat{\tau}_*}(d)$ .

**Proof.** For simplicity, we only consider the case that  $0 < q < a$  and  $d \in (0, \tilde{d}_2]$ , and other cases can be proved similarly. For this case,  $\hat{\tau}_* = \tau_{p,l}^d$  and  $\hat{\nu}_* = \nu_p^d$  with  $l = 0, 1, \dots$  and  $p = 1, \dots, n$ . It from Theorem [3.9](#page-19-3) that  $\mathcal{N}\left[A_{\tau_{p,l}^{d}}(d) - i\nu_{p}^{d}\right] = \text{span}\left[e^{i\nu_{p}^{d}\theta}\psi_{p}^{d}\right]$ , where  $\theta \in [-\tau_{p,l}^{d}, 0]$  and  $\psi_{p}^{d}$  is defined in Theorem [3.9.](#page-19-3) Now we show that

$$
\mathcal{N}\bigg[A_{\tau_{p,l}^d}(d) - \mathrm{i}\nu_p^d\bigg]^2 = \mathcal{N}\bigg[A_{\tau_{p,l}^d}(d) - \mathrm{i}\nu_p^d\bigg].
$$

If  $\boldsymbol{\phi} \in \mathcal{N} \Big[ A_{\tau_{p,l}^d}(d) - i v_p^d \Big]^2$ , then  $\left[A_{\tau^d_{p,l}}(d) - i\nu^d_p\right]\boldsymbol{\phi} \in \mathcal{N}\left[A_{\tau^d_{p,l}}(d) - i\nu^d_p\right] = \text{span}\left[e^{i\nu^d_p\boldsymbol{\phi}}\boldsymbol{\psi}^d_p\right],$ 

and consequently, there is a constant  $\sigma$  such that

<span id="page-22-0"></span>
$$
\left[A_{\tau_{p,l}^d}(d)-{\rm i}\nu_p^d\right]\boldsymbol{\phi}=\sigma e^{{\rm i}\nu_p^d\theta}\boldsymbol{\psi}_p^d.
$$

This together with  $(3.2)$  and  $(3.3)$  yields

$$
\dot{\boldsymbol{\phi}}(\theta) = i v_p^d \boldsymbol{\phi}(\theta) + \sigma e^{i v_p^d \theta} \boldsymbol{\psi}_p^d, \quad \theta \in \left[ -\tau_{p,l}^d, 0 \right],
$$
\n
$$
\dot{\boldsymbol{\phi}}(0) = dD\boldsymbol{\phi}(0) + qQ\boldsymbol{\phi}(0) + \text{diag}\left( a - bu_{dj} \right) \boldsymbol{\phi}(0) - \text{diag}\left( bu_{dj} \right) \boldsymbol{\phi}(-\tau_{p,l}^d). \tag{4.1}
$$

By the first equation of  $(4.1)$ , we have

<span id="page-22-1"></span>
$$
\boldsymbol{\phi}(\theta) = \boldsymbol{\phi}(0)e^{i\nu_p^d\theta} + \sigma \theta e^{i\nu_p^d\theta} \boldsymbol{\psi}_p^d,
$$
\n
$$
\boldsymbol{\dot{\phi}}(0) = i\nu_p^d \boldsymbol{\phi}(0) + \sigma \boldsymbol{\psi}_p^d.
$$
\n(4.2)

Note from Theorem [3.9](#page-19-3) that  $e^{-i\tau_{p,l}^d v_p^d} = e^{-i\theta_p^d}$  with  $\theta_p^d$  defined in Lemma [3.8.](#page-17-7) Then, substituting [\(4.2\)](#page-22-1) into the second equation of  $(4.1)$ , we have

<span id="page-22-2"></span>
$$
\mathcal{M}\left(d,\theta_p^d,\nu_p^d\right)\boldsymbol{\phi}(0) = \sigma\left(\boldsymbol{\psi}_p^d - \tau_{p,l}^d e^{-i\theta_p^d} \text{diag}\left(bu_{d,j}\right)\boldsymbol{\psi}_p^d\right),\tag{4.3}
$$

where  $\mathcal{M}(d, v, \theta)$  is defined in [\(3.6\)](#page-10-5).

Let  $\mathcal{M}^H(d, v, \theta)$  be the conjugate transpose of  $\mathcal{M}(d, v, \theta)$ . Using similar arguments as in the proof of Lemma [3.8,](#page-17-7) we obtain that for  $d \in (0, d_2]$ ,

$$
\left\{ \boldsymbol{\varphi} \in \mathbb{C}^n : \mathcal{M}^H(d, v_p^d, \theta_p^d) \boldsymbol{\varphi} = \mathbf{0} \right\} = \{ c \boldsymbol{\varphi}_p^d : c \in \mathbb{C} \},\
$$

and  $\lim_{d\to 0} \boldsymbol{\varphi}_p^d = \boldsymbol{\varphi}_p^0$ . Here,  $\boldsymbol{\varphi}_p^0 = (\varphi_{p,1}^0, \cdots, \varphi_{p,n}^0)^T$  satisfies

$$
\varphi_{p,j}^0 = 0
$$
 for  $p + 1 \le j \le n$ ,  $\varphi_{p,p}^0 = 1$ ,  
\n
$$
\varphi_{p,j}^0 = (-1)^{p-j} \prod_{k=j}^{p-1} \frac{q}{\overline{h}_k(\theta_p^0, v_p^0)}
$$
 for  $1 \le j \le p - 1$ ,

where  $\overline{h}_k$   $(\theta_p^0, v_p^0)$  is the conjugate of  $h_k(\theta_p^0, v_p^0)$ , and  $h_k(\theta, v)$  is defined in [\(3.23\)](#page-15-7). One can also refer to  $(3.30)$  for  $p = 2$ . Then by  $(4.3)$ , we have

$$
0 = \langle \mathcal{M}^H \left( d, \theta_p^d, v_p^d \right) \boldsymbol{\varphi}_p^d, \boldsymbol{\phi}(0) \rangle = \langle \boldsymbol{\varphi}_p^d, \mathcal{M} \left( d, \theta_p^d, v_p^d \right) \boldsymbol{\phi}(0) \rangle
$$
  
\n
$$
= \sigma \left[ \langle \boldsymbol{\varphi}_p^d, \boldsymbol{\psi}_p^d \rangle - \langle \boldsymbol{\varphi}^d, \tau_{p,l}^d e^{-i\theta_p^d} \text{diag} \left( b u_{dj} \right) \boldsymbol{\psi}_p^d \rangle \right]
$$
  
\n
$$
= \sigma \left[ \sum_{j=1}^n \overline{\varphi}_{p,j}^d \psi_{p,j}^d - \tau_{p,l}^d e^{-i\theta_p^d} \sum_{j=1}^n b u_{dj} \overline{\varphi}_{p,j}^d \psi_{p,j}^d \right] = \sigma S_{p,l}(d).
$$

where

$$
S_{p,l}(d) := \sum_{j=1}^n \overline{\varphi}_{p,j}^d \psi_{p,j}^d - \tau_{p,l}^d e^{-i\theta_p^d} \sum_{j=1}^n bu_{d,j} \overline{\varphi}_{p,j}^d \psi_{p,j}^d \text{ for } p = 1, \cdots, n \text{ and } l = 0, 1, \cdots.
$$

By Lemmas [3.6,](#page-15-5) [3.8](#page-17-7) and Theorem [3.9,](#page-19-3) we obtain that

$$
\lim_{d \to 0} S_{p,l}(d) = 1 - \frac{\theta_p^0 + 2l\pi}{\nu_p^0} bu_{0,p} \cos \theta_p^0 + i \frac{\theta_p^0 + 2l\pi}{\nu_p^0} bu_{0,p} \sin \theta_p^0 \neq 0,
$$

which implies  $\sigma = 0$ . Therefore, for any  $l = 1, 2, \dots$ ,

$$
\mathcal{N}\left[A_{\tau_{p,l}^d}(d) - \mathrm{i}\nu_p^d\right]^2 = \mathcal{N}\left[A_{\tau_{p,l}^d}(d) - \mathrm{i}\nu_p^d\right],
$$

and consequently,  $i\nu_p^d$  is a simple eigenvalue of  $A_{\tau_{p,l}^d}$  for  $l = 0, 1, 2, \cdots$ .

It follows from Lemma [4.1](#page-22-3) that  $\mu = i\hat{v}_*$  is an algebraically simple eigenvalue of  $A_{\hat{\tau}_*}$ . Then, by the implicit function theorem, we see that there exists a neighbourhood  $\hat{O}_* \times \hat{D}_* \times \hat{H}_*$  of  $(\hat{\tau}_*, \hat{\psi}_*, \hat{\psi}_*)$  and a continuously differentiable function  $(\mu(\tau), \psi(\tau)) : \hat{O}_* \to \hat{D}_* \times \hat{H}_*$  such that  $\mu(\hat{\tau}_*) = i\hat{v}_*, \psi(\hat{\tau}_*) = \hat{\psi}_*,$ <br>and for  $\tau \in \hat{O}$ and for  $\tau \in \widehat{O}_*$ ,

<span id="page-23-0"></span>
$$
\Delta(d, \mu(\tau), \tau)\psi(\tau) = dD\psi(\tau) + qQ\psi(\tau) + \text{diag}(a - bu_{d,j})\psi(\tau)
$$
  

$$
-e^{-\mu(\tau)\tau}\text{diag}(bu_{d,j})\psi(\tau) - \mu(\tau)\psi(\tau) = 0.
$$
 (4.4)

 $\Box$ 

Here

$$
\widehat{\boldsymbol{\psi}}_{*} = \begin{cases} \psi_{d}, & \text{if } a < q < na \text{ and } d \in (d_{q}^{*}, \widetilde{d}_{1}] \\ \psi_{p}^{d}, & \text{if } 0 < q < a \text{ and } d \in (0, \widetilde{d}_{2}] \\ \psi^{\lambda}, & \text{if } 0 < q < na \text{ and } d \in [\widetilde{d}_{3}, \infty) \end{cases}
$$

with  $\lambda = 1/d$  and  $p = 1, \dots, n$ , where  $(\psi_d, \tilde{d}_1), (\psi_p^d, \tilde{d}_2)$  and  $(\psi^{\lambda}, \tilde{d}_3)$  are defined in Theorems [3.5,](#page-14-2) [3.9](#page-19-3) and [3.13,](#page-21-4) respectively. By a direct calculation, we obtain the following transversality condition.

<span id="page-23-1"></span>**Lemma 4.2.** *Let*  $\hat{\tau}$ , *be defined in Lemma [4.1.](#page-22-3) Then* 

$$
\frac{d\mathcal{R}e\left[\mu\left(\widehat{\tau}_{*}\right)\right]}{d\tau}>0.
$$

**Proof.** Similar to Lemma [4.1,](#page-22-3) we only consider the case that  $0 < q < a$  and  $d \in (0, \overline{d_2}]$ , and other cases can be proved similarly. For this case,  $\hat{\tau}_* = \tau^d_{p,l}$ ,  $\hat{\nu}_* = \nu^d_p$  and  $\hat{\psi}_* = \psi^d_p$  with  $p = 1, \dots, n$  and  $l = 0, 1, \dots$ .

Differentiating [\(4.4\)](#page-23-0) with respect to  $\tau$  at  $\tau = \tau_{p,l}^d$ , we have

<span id="page-24-2"></span>
$$
\frac{d\mu\left(\tau_{p,l}^{d}\right)}{d\tau}\left[-\boldsymbol{\psi}_{p}^{d}+\tau_{p,l}^{d}e^{-i\theta_{p}^{d}}\text{diag}\left(bu_{dj}\right)\boldsymbol{\psi}_{p}^{d}\right] \n+\mathcal{M}\left(d,\theta_{p}^{d},v_{p}^{d}\right)\frac{d\boldsymbol{\psi}\left(\tau_{p,l}^{d}\right)}{d\tau}+i\nu_{p}^{d}e^{-i\theta_{p}^{d}}\text{diag}\left(bu_{dj}\right)\boldsymbol{\psi}_{p}^{d}=\mathbf{0},
$$
\n(4.5)

where  $\mathcal{M}(d, v, \theta)$  is defined in [\(3.6\)](#page-10-5). Note that, for  $l = 0, 1, \dots$ ,

$$
0 = \left\langle \mathcal{M}^H \left( d, \theta_p^d, v_p^d \right) \boldsymbol{\varphi}_p^d, \frac{d \boldsymbol{\psi} \left( \tau_{p,l}^d \right)}{d \tau} \right\rangle = \left\langle \boldsymbol{\varphi}_p^d, \mathcal{M} \left( d, \theta_p^d, v_p^d \right) \frac{d \boldsymbol{\psi} \left( \tau_{p,l}^d \right)}{d \tau} \right\rangle,
$$

where  $\mathcal{M}^H(d, v, \theta)$  and  $\varphi_p^d$  are defined in the proof of Lemma [4.1.](#page-22-3) This combined with [\(4.5\)](#page-24-2) implies that

$$
\frac{d\mu(\tau_{p,l}^d)}{d\tau} = \frac{1}{|S_l(d)|^2} \left[ i\nu_p^d e^{-i\theta_p^d} \left( \sum_{j=1}^n b u_{d,j} \overline{\varphi}_{p,j}^d \psi_{p,j}^d \right) \left( \sum_{j=1}^n \varphi_{p,j}^d \overline{\psi}_{p,j}^d \right) \right.\left. - i\nu_p^d \tau_{p,l}^d \left( \sum_{j=1}^n b u_{d,j} \varphi_{p,j}^d \overline{\psi}_{p,j}^d \right) \left( \sum_{j=1}^n b u_{d,j} \overline{\varphi}_{p,j}^d \psi_{p,j}^d \right) \right].
$$

By Lemmas [3.6](#page-15-5) and [3.8](#page-17-7) and Theorem [3.9,](#page-19-3) we obtain that

$$
\lim_{d\to 0}\frac{d\mathcal{R}e\left[\mu\left(\tau_{p,l}^d\right)\right]}{d\tau}>0.
$$

This completes the proof.

By Theorems [3.5,](#page-14-2) [3.9](#page-19-3) and [3.13](#page-21-4) and Lemmas [4.1](#page-22-3) and [4.2,](#page-23-1) we obtain the following result.

<span id="page-24-3"></span>**Theorem 4.3.** Let  $u_d$  be the unique positive equilibrium of model [\(1.5\)](#page-1-2) obtained in Proposition [2.4.](#page-4-4) *Then the following statements hold:*

- (i) For  $q \in (a, na)$  and  $d \in (d^*_q, \tilde{d}_1]$  with  $0 < \tilde{d}_1 d^*_q \ll 1$ ,  $u_d$  is locally asymptotically stable for  $\tau \in (0, \tau_1)$  and unstable for  $\tau \in (0, \infty)$ . Moreover model (1.5) undergoes a Hopf bifurcation when  $[0, \tau_{d,0})$  and unstable for  $\tau \in (\tau_{d,0}, \infty)$ *. Moreover, model* [\(1.5\)](#page-1-2) undergoes a Hopf bifurcation when  $\tau = \tau_{d,0}$
- *(ii)* For  $q \in (0, a)$  and  $d \in (0, \tilde{d}_2]$  with  $0 < \tilde{d}_2 \ll 1$ ,  $u_d$  is locally asymptotically stable for  $\tau \in [0, \tau_{n,0}^d]$ <br>and unstable for  $\tau \in (\tau^d \cap \infty)$ . Moreover, model (1.5) undergoes a Hopf bifurcation when  $\tau$  $a$ nd unstable for  $\tau\in(\tau^d_{n,0},\infty)$ . Moreover, model [\(1.5\)](#page-1-2) undergoes a Hopf bifurcation when  $\tau=\tau^d_{n,0}$ .
- (iii) For  $q \in (0, na)$  and  $d \in [\tilde{d}_3, \infty)$  with  $\tilde{d}_3 \gg 1$ ,  $u_d$  is locally asymptotically stable for  $\tau \in [0, \tau_0^{\lambda}]$ <br>and unstable for  $\tau \in (\tau^{\lambda}, \infty)$  with  $\lambda = 1/d$  Moreover model (1.5) undergoes a Hopf bifurcation *and unstable for*  $\tau \in (\tau_0^{\lambda}, \infty)$  *with*  $\lambda = 1/d$ *. Moreover, model* [\(1.5\)](#page-1-2) *undergoes a Hopf bifurcation* when  $\tau = \tau_0^{\lambda}$ .

<span id="page-24-1"></span>*Here*  $\tau_{d,0}$ ,  $\tau_{n,0}^d$  *and*  $\tau = \tau_0^{\lambda}$  *are defined in Theorems* [3.5,](#page-14-2) [3.9](#page-19-3) *and* [3.13,](#page-21-4) *respectively.* 

#### **5. The effect of drift rate and numerical simulations**

In this section, we show the effect of drift rate and give some numerical simulations. Throughout this section, we define the minimum Hopf bifurcation value by the first Hopf bifurcation value.

If the directed drift rate  $q = 0$  (non-advective case), then model [\(1.5\)](#page-1-2) admits a unique positive equilibrium  $u_d = (a/b, \dots, a/b)^T$  for all  $d > 0$ . By the framework of [\[43,](#page-30-26) [48\]](#page-31-8), we can show the existence of a Hopf bifurcation as follows. Here, we omit the proof a Hopf bifurcation as follows. Here, we omit the proof.

<span id="page-24-0"></span>**Proposition 5.1.** *Let*  $q = 0$ *. Then the first Hopf bifurcation value of model* [\(1.5\)](#page-1-2) *is*  $\tau_{non} = \pi/2a$ *. Moreover, the unique positive equilibrium*  $u_d$  *of model* [\(1.5\)](#page-1-2) *is stable for*  $\tau < \tau_{non}$  *and unstable for*  $\tau > \tau_{\text{non}}$ , and model [\(1.5\)](#page-1-2) undergoes a Hopf bifurcation when  $\tau = \tau_{\text{non}}$ .

 $\Box$ 

Therefore, the first Hopf bifurcation value  $\tau_{non}$  for  $q=0$  is independent of the random diffusion rate *d*. By Theorems [3.5,](#page-14-2) [3.9,](#page-19-3) [3.13](#page-21-4) and [4.3,](#page-24-3) we see that the first Hopf bifurcation value τ*adv* depends on the diffusion rate *d* for  $q \neq 0$ . Actually, we show that it can be strictly monotone decreasing in *d* when *d* is large.

<span id="page-25-0"></span>**Proposition 5.2.** Assume that  $q \in (0, na)$ , and let  $\lambda = 1/d$ . Then, for  $d \in [\tilde{d}_3, \infty)$  with  $\tilde{d}_3 \gg 1$ , the first *Hopf bifurcation value of model* [\(1.5\)](#page-1-2) *is*  $\tau_{adv} = \tau_0^{\lambda}$ , where  $\tau_0^{\lambda}$  *and*  $\tilde{d}_3$  *are defined in Theorem* [3.13.](#page-21-4) *Moreover, the following statements hold:*

*(i)*

<span id="page-25-1"></span>
$$
\left(\tau_0^{\lambda}\right)' \big|_{\lambda=0} = \frac{\pi q^2 (n+1)(n-1)}{12(na-q)^2} > 0,\tag{5.1}
$$

*where is the derivative with respect to* λ*;*

*(ii)* There exists  $\hat{d}_3 > \tilde{d}_3$  such that  $\tau_{adv} > \tau_{non}$  for  $d \in [\hat{d}_3, \infty)$ , and  $\tau_{adv}$  is strictly monotone decreasing *in*  $d \in [\hat{d}_3, \infty)$ .

**Proof.** By Theorems [3.13](#page-21-4) and [4.3,](#page-24-3) we see that for  $d \in [\tilde{d}_3, \infty)$  with  $\tilde{d}_3 \gg 1$ , the first Hopf bifurcation value of model [\(1.5\)](#page-1-2) is  $\tau_{adv} = \tau_0^{\lambda}$ . We first show that (i) holds. Note from Lemmas [3.11](#page-20-2) and [3.12](#page-21-3) that  $(\gamma^{\lambda}, \nu^{\lambda}, \theta^{\lambda}, z^{\lambda})$  is the unique solution of [\(3.42\)](#page-20-3) and  $(\gamma^0, \nu^0, \theta^0, z^0) = (1, a - \frac{q}{n}, \frac{\pi}{2}, \mathbf{0})$ . Differentiating the first equation of [\(3.42\)](#page-20-3) with respect to  $\lambda$  at  $\lambda = 0$  and noticing that  $\overline{z}^0 = 0$ , we have

$$
0 = \sum_{k=1}^{n} D_{jk} (z_k^{\lambda})' \Big|_{\lambda=0} + q \sum_{k=1}^{n} Q_{jk} \gamma^0 \zeta_k + (a - bu_j^0) \gamma^0 \zeta_j - e^{-i\theta^0} bu_j^0 \gamma^0 \zeta_j - i\nu^0 \gamma^0 \zeta_j - \frac{1}{n} \widetilde{F}_2 (\gamma^0, \nu^0, \theta^0, z^0, 0),\n0 = q \sum_{j=1}^{n} \sum_{k=1}^{n} Q_{jk} [(\gamma^{\lambda})' \zeta_k + (z_k^{\lambda})'] \Big|_{\lambda=0} - \sum_{j=1}^{n} (bu_j^{\lambda})' \Big|_{\lambda=0} \gamma^0 \zeta_j + \sum_{j=1}^{n} \Big[ (a - bu_j^0) - e^{-i\theta^0} bu_j^0 - i\nu^0 \Big] [(\gamma^{\lambda})' \zeta_j + (z_j^{\lambda})'] \Big|_{\lambda=0} + \sum_{j=1}^{n} \Big[ (\theta^{\lambda})' \Big|_{\lambda=0} bu_j^0 - e^{-i\theta^0} (bu_j^{\lambda})' \Big|_{\lambda=0} - i (\nu^{\lambda})' \Big|_{\lambda=0} \Big] \gamma^0 \zeta_j, \n0 = 2\gamma^0 (\gamma^{\lambda})' \Big|_{\lambda=0} \| \zeta \|_2^2.
$$
\n(5.2)

By the third equation of [\(5.2\)](#page-25-1), we have  $(\gamma^{\lambda})' \big|_{\lambda=0} = 0$ . Then plugging  $(\gamma^{\lambda})' \big|_{\lambda=0} = 0$  into the first and second equations of [\(5.2\)](#page-25-1), and noticing that  $(\gamma^0, \nu^0, \theta^0) = (1, a - \frac{q}{n}, \frac{\pi}{2})$ 2  $\left( \sum_{i=1}^{n} z_i \in (\widetilde{X}_1)_{\mathbb{C}} \right)$  and  $\widetilde{F}_2(\gamma^0, \nu^0, \theta^0, z^0, 0) = 0$ , we have

$$
\begin{cases}\n\sum_{j=1}^{n} D_{1k}(z_k^{\lambda})'\Big|_{\lambda=0} - \frac{q}{n} + \frac{q}{n} \cdot \frac{1}{n} = 0, \\
\sum_{j=1}^{n} D_{jk}(z_k^{\lambda})'\Big|_{\lambda=0} + \frac{q}{n} \cdot \frac{1}{n} = 0, \ j = 2, \cdots, n, \\
\left[ -q(z_n^{\lambda})' - \frac{1}{n} \sum_{j=1}^{n} (bu_j^{\lambda})' + \left(a - \frac{q}{n}\right) (\theta^{\lambda})' + i \frac{1}{n} \sum_{j=1}^{n} (bu_j^{\lambda})' - i (\nu^{\lambda})' \right] \right]_{\lambda=0} = 0.\n\end{cases}
$$

<https://doi.org/10.1017/S0956792524000342>Published online by Cambridge University Press

This combined with  $(2.20)$  and Proposition [5.4](#page-31-9) in the appendix implies that

$$
(z_n^{\lambda})'|_{\lambda=0} = \frac{q(n+1)(n-1)}{6n}, (v^{\lambda})'|_{\lambda=0} = -\frac{q^2(n+1)(n-1)}{6n}, (\theta^{\lambda})'|_{\lambda=0} = 0.
$$

Then, by Theorem [3.13,](#page-21-4) we have

$$
\left(\tau_0^{\lambda}\right)' \big|_{\lambda=0} = \left(\frac{\theta^{\lambda}}{\nu^{\lambda}}\right)' \bigg|_{\lambda=0} = \frac{\pi q^2 n(n+1)(n-1)}{12(na-q)^2}.
$$

Now we consider (ii). By Lemmas [3.11,](#page-20-2) [3.12](#page-21-3) and Theorem [3.13,](#page-21-4) we see that

$$
\lim_{\lambda \to 0} \tau_{adv} = \lim_{\lambda \to 0} \tau_0^{\lambda} = \frac{\pi}{2(a - \frac{q}{n})} > \tau_{non}.
$$

This, combined with (i), implies that (ii) holds. This completes the proof.

Then we consider the case of small diffusion rate.

<span id="page-26-0"></span>**Proposition 5.3.** Assume that  $q \in (0, a)$ . Then, for  $d \in (0, \tilde{d}_2]$  with  $\tilde{d}_2 \ll 1$ , the first Hopf bifurcation *value of model* [\(1.5\)](#page-1-2) *is*  $\tau_{adv} = \tau_{n,0}^d$ , where  $\tau_{n,0}^d$  *and*  $\tilde{d}_2$  *are defined in Theorem* [3.9.](#page-19-3) Moreover, there exists  $\hat{d}_2 \in (0, \tilde{d}_2]$  *such that*  $\tau_{adv} > \tau_{non}$  *for*  $d \in (0, \hat{d}_2]$ *.* 

**Proof.** By Theorems [3.9](#page-19-3) and [4.3,](#page-24-3) we see that the first Hopf bifurcation value is  $\tau_{adv} = \tau_{n,0}^d$ , and

$$
\lim_{d \to 0} \tau_{adv} = \lim_{d \to 0} \tau_{n,0}^d = \frac{\theta_n^0}{\nu_n^0}.
$$
\n(5.3)

By Lemma [3.7](#page-15-0) (ii), we see that  $\frac{\theta_n^0}{\theta_n^0}$  $> \frac{\pi}{2}$  $\frac{\pi}{2a}$ . This, combined with [\(5.3\)](#page-26-1), implies that there exists  $\hat{d}_2 \in (0, \tilde{d}_2]$  $v_n^0$ such that  $\tau_{adv} > \tau_{non}$  for  $d \in (0, d_2]$ . П

It follows from Propositions [5.2](#page-25-0) and [5.3](#page-26-0) that the first Hopf bifurcation value in advective environments is larger than that in non-advective environments if  $d \gg 1$  or  $d \ll 1$ , see Figure [2.](#page-3-0) This result suggests that directed movements of the individuals inhibit the occurrence of Hopf bifurcation.

Now we give some numerical simulations. We choose three patches, that is,  $n = 3$ , and set  $a = 1$  and  $b = 1$ . Then we can numerically compute the first Hopf bifurcation  $\tau_{adv}$  for a wider range of parameters. For the case  $q \in (0, a)$ , we prove that large delay can also induce Hopf bifurcations for model [\(1.5\)](#page-1-2) if  $0 < d \ll 1$  or  $d \gg 1$  (Theorem [4.3](#page-24-3) (ii) and (iii)). Then we compute that there exist three families of Hopf bifurcation curves  $\left\{\tau_{j,l}^d\right\}_{l=1}^{\infty}$  (*j* = 1, 2, 3). For simplicity, we only plot the first one for each family  $\left\{\tau_{j,l}^d\right\}_{l=1}^{\infty}$   $(j = 1, 2, 3)$  in Figure [4.](#page-27-0)

Then  $\tau_{adv} = \min_{1 \leq j \leq 3} \tau_{j,0}^d$ , and it exists for  $d \in (0, \infty)$ , which implies that delay-induced Hopf bifurcation may occur for  $d \in (0, \infty)$ . Actually, we choose  $d = 0.06, 1.5, 20, 150$  and numerically show that there exist periodic solutions; see Figure [5.](#page-27-1)

For the case  $q \in (a, na)$ , we prove that large delay can induce Hopf bifurcations for model [\(1.5\)](#page-1-2) if  $0 < d - d_q^* \ll 1$  or  $d \gg 1$  (Theorem [4.3](#page-24-3) (i) and (iii)). In Figure [6,](#page-28-0) we plot  $\tau_{adv}$  for this case, and it exists for  $d \in (d^*_q, \infty)$ , which implies that delay-induced Hopf bifurcation may occur for  $d \in (d^*_q, \infty)$ .

By Figures [4](#page-27-0) and [6,](#page-28-0) we conjecture that  $\tau_{adv}$  change monotonicity once with respect to *d*. In fact, by Proposition [5.2,](#page-25-0) τ*adv* is decreasing in *d* when *d* is sufficiently large. Moreover, in Propositions [5.2](#page-25-0) and [5.3,](#page-26-0) we show that  $\tau_{adv} > \tau_{non}$ , which suggests that directed movements of the individuals inhibit the occurrence of Hopf bifurcation. To illustrate this phenomenon, we fix  $\tau = 1.6$ ,  $d = 1.14$  and choose the same initial values for  $q = 0$  and  $q = 2$ . As is shown in Figure [7,](#page-28-1) periodic oscillations disappear for  $q \neq 0$ .

<span id="page-26-1"></span> $\Box$ 

<span id="page-27-0"></span>

*Figure 4. The relation between Hopf bifurcation values and dispersal rate d for the case*  $q \in (0, a)$  *with*  $a = 1, b = 1$  *and*  $q = 0.6$ *.* (*a*)  $d \in (0, 1]$ *;* (*b*)  $d \in [5, 150]$ *.* 

<span id="page-27-1"></span>

*Figure 5.* Periodic solutions induced by a Hopf bifurcation with  $a = 1$ ,  $b = 1$  and  $q = 0.6$ . (a)  $d = 0.06$ *and*  $\tau = 3.1$ *; (b)*  $d = 1.5$  *and*  $\tau = 2.1$ *; (c)*  $d = 20$  *and*  $\tau = 2.1$ *; (d)*  $d = 150$  *and*  $\tau = 2.0$ *.* 

<span id="page-28-0"></span>

*<i>Figure 6.* The relation between Hopf bifurcation value and dispersal rate d for the case  $q \in (a, na)$ *with*  $a = 1$ *,*  $b = 1$  *and*  $q = 2$ *.* 

<span id="page-28-1"></span>

*Figure 7. Directed drift rate q inhibit the occurrence of Hopf bifurcation. Here,*  $a = 1$ *,*  $b = 1$ *,*  $d = 1.14$ *and*  $\tau = 1.6$ *.* (*a*)  $q = 0$ *;* (*b*)  $q = 2$ *.* 

We remark that model  $(1.5)$  is a discrete form of model  $(1.6)$ , where  $D_{jk}$  is defined in  $(1.3)$  and

<span id="page-28-2"></span>
$$
Q_{jk} = \begin{cases} 1, & j = k + 1, \\ -1, & j = k = 1, \dots, n - 1, \\ -\beta, & j = k = n, \\ 0, & \text{otherwise.} \end{cases}
$$
(5.4)

In this paper, we consider the case  $\beta = 1$ , and it is natural to ask whether  $\beta$  (the population loss rate at the downstream end) affects Hopf bifurcations. Then we consider this problem from the point view of numerical simulations and choose

 $n = 3$  (three patches),  $a = 1$ ,  $b = 1$ ,  $q = 0.6$ ,  $d = 2$ .

<span id="page-29-10"></span>

*<i>Figure 8. The effect of*  $\beta$  *on the dynamics of model [\(1.5\)](#page-1-2)* with  $D_{ik}$  *and*  $Q_{ik}$  *defined in [\(1.3\)](#page-0-2) and* [\(5.4\)](#page-28-2)*, respectively.* (a)  $\beta = 0.9$ ,  $\tau = 2.1$ ; (b)  $\beta = 1.5$ ,  $\tau = 2.1$ .

If  $\beta = 1$ , we compute that  $\tau_{adv} \approx 2.03$ . Then set  $\tau = 2.1$ , we show that there also exists periodic solutions for  $\beta = 0.9$ , and periodic oscillations disappear for  $\beta = 1.5$ , see Figure [8.](#page-29-10) Then, we conjecture that if the positive equilibrium of [\(1.5\)](#page-1-2) exists, the minimum Hopf bifurcation value for the case  $\beta > 1$  (resp.  $\beta$  < 1) is larger (resp. smaller) than that for the case  $\beta = 1$ .

**Ethical standards.** Not applicable.

**Competing interests.** The authors declare that they have no conflict of interest.

**Author contributions.** All authors contributed to the study conception and design. SC developed the idea for the study. The manuscript was written by LW, SZ and SC, and LW and SZ prepared Figures [1](#page-1-0)[–7.](#page-28-1) All authors read and approved the final manuscript.

**Financial support.** This research is Taishan Scholars Program of Shandong Province (No. tsqn 202306137), National Natural Science Foundation of China (Nos. 12171117 and 12101161) and Heilongjiang Provincial Natural Science Foundation of China (No. YQ2021A007).

**Data availability statement.** All data generated or analysed during this study are included in this published article.

#### **References**

- <span id="page-29-3"></span>[1] An, Q., Wang, C. & Wang, H. (2020) Analysis of a spatial memory model with nonlocal maturation delay and hostile boundary condition. *Discrete Contin. Dyn. Syst.* **40**(10), 5845–5868.
- <span id="page-29-4"></span>[2] Busenberg, S. & Huang, W. (1996) Stability and Hopf bifurcation for a population delay model with diffusion effects. *J. Differt. Eq.* **124**(1), 80–107.
- <span id="page-29-1"></span>[3] Cantrell, R. S., Cosner, C. & Lou, Y. (2006) Movement toward better environments and the evolution of rapid diffusion. *Math. Biosci.* **204**(2), 199–214.
- <span id="page-29-8"></span>[4] Chang, L., Duan, M., Sun, G. & Jin, Z. (2020) Cross-diffusion-induced patterns in an SIR epidemic model on complex networks. *Chaos* **30**(1), 013147.
- <span id="page-29-9"></span>[5] Chang, L., Liu, C., Sun, G., Wang, Z. & Jin, Z. (2019) Delay-induced patterns in a predator-prey model on complex networks with diffusion. *New J. Phys.* **21**(7), 073035.
- <span id="page-29-2"></span>[6] Chen, S., Liu, J. & Wu, Y. (2022) Invasion analysis of a two-species Lotka-Volterra competition model in an advective patchy environment. *Stud. Appl. Math.* **149**(3), 762–797.
- <span id="page-29-6"></span>[7] Chen, S., Lou, Y. & Wei, J. (2018) Hopf bifurcation in a delayed reaction-diffusion-advection population model. *J. Differ. Eq.* **264**(8), 5333–5359.
- <span id="page-29-0"></span>[8] Chen, S., Shen, Z. & Wei, J. (2023) Hopf bifurcation of a delayed single population model with patch structure. *J. Dynam. Differ. Eq.* **35**(2), 1457–1487.
- <span id="page-29-7"></span><span id="page-29-5"></span>[9] Chen, S. & Shi, J. (2012) Stability and Hopf bifurcation in a diffusive logistic population model with nonlocal delay effect. *J. Differ. Eq.* **253**(12), 3440–3470.
- [10] Chen, S., Shi, J., Shuai, Z. & Wu, Y. (2023) Evolution of dispersal in advective patchy environments. *J. Nonlinear Sci.* **33**(3), 35.
- <span id="page-30-14"></span>[11] Chen, S., Wei, J. & Zhang, X. (2020) Bifurcation analysis for a delayed diffusive logistic population model in the advective heterogeneous environment. *J. Dynam. Differ. Eq.* **32**(2), 823–847.
- <span id="page-30-10"></span>[12] Chen, S. & Yu, J. (2016) Stability and bifurcations in a nonlocal delayed reaction-diffusion population model. *J. Differ. Eq.* **260**(1), 218–240.
- <span id="page-30-5"></span>[13] Chen, X., Lam, K.-Y. & Lou, Y. (2012) Dynamics of a reaction-diffusion-advection model for two competing species. *Discrete Contin. Dyn. Syst.* **32**(11), 3841–3859.
- <span id="page-30-29"></span>[14] Cosner, C. (1996) Variability, vagueness and comparison methods for ecological models. *Bull. Math. Biol.* **58**(2), 207–246.
- <span id="page-30-23"></span>[15] Duan, M., Chang, L. & Jin, Z. (2019) Turing patterns of an SI epidemic model with cross-diffusion on complex networks. *Physica A* **533**, 122023.
- <span id="page-30-21"></span>[16] Fernandes, L. D. & Aguiar, M. D. (2012) Turing patterns and apparent competition in predator-prey food webs on networks. *Phy. Rev. E* **86**(5), 056203.
- [17] Gou, W., Jin, Z. & Wang, H. (2023) Hopf bifurcation for general network-organized reaction-diffusion systems and its application in a multi-patch predator-prey system. *J. Differ. Eq.* **346**, 64–107.
- <span id="page-30-22"></span>[18] Gou, W., Song, Y. & Jin, Z. (2023) The steady state bifurcation for general network-organized reaction-diffusion systems and its application in a metapopulation epidemic model. *SIAM J. Appl. Dyn. Syst.* **22**(2), 559–602.
- <span id="page-30-11"></span>[19] Guo, S. (2015) Stability and bifurcation in a reaction-diffusion model with nonlocal delay effect. *J. Differ. Eq.* **259**(4), 1409–1448.
- <span id="page-30-12"></span>[20] Guo, S. & Yan, S. (2016) Hopf bifurcation in a diffusive Lotka-Volterra type system with nonlocal delay effect. *J. Differ. Eq.* **260**(1), 781–817.
- <span id="page-30-32"></span>[21] Hale, J. (1977). *Theory of functional differential equations. Applied Mathematical Sciences*. 2nd ed. Springer-Verlag, New York-Heidelberg.
- <span id="page-30-6"></span>[22] Hamida, Y. (2017). The evolution of dispersal for the case of two-patches and two-species with travel loss [*PhD thesis*). The Ohio State University
- <span id="page-30-13"></span>[23] Hu, R. & Yuan, Y. (2011) Spatially nonhomogeneous equilibrium in a reaction-diffusion system with distributed delay. *J. Differ. Eq.* **250**(6), 2779–2806.
- <span id="page-30-0"></span>[24] Huang, D., Chen, S. & Zou, X. (2023) Hopf bifurcation in a delayed population model over patches with general dispersion matrix and nonlocal interactions. *J. Dynam. Differ. Eq.* **35**, 3521–3543.
- <span id="page-30-7"></span>[25] Jiang, H., Lam, K.-Y. & Lou, Y. (2020) Are two-patch models sufficient? The evolution of dispersal and topology of river network modules. *Bull. Math. Biol.* **82**(10), 42.
- <span id="page-30-8"></span>[26] Jiang, H., Lam, K.-Y. & Lou, Y. (2021) Three-patch models for the evolution of dispersal in advective environments: Varying drift and network topology. *Bull. Math. Biol.* **83**(10), 46.
- <span id="page-30-15"></span>[27] Jin, Z. & Yuan, R. (2021) Hopf bifurcation in a reaction-diffusion-advection equation with nonlocal delay effect. *J. Differ. Eq.* **271**, 533–562.
- <span id="page-30-28"></span>[28] Li, C.-K. & Schneider, H. (2002) Applications of Perron-Frobenius theory to population dynamics. *J. Math. Biol.* **44**(5), 450–462.
- <span id="page-30-30"></span>[29] Li, M. Y. & Shuai, Z. (2010) Global-stability problem for coupled systems of differential equations on networks. *J. Differ. Eq.* **248**(1), 1–20.
- <span id="page-30-16"></span>[30] Li, Z. & Dai, B. (2021) Stability and Hopf bifurcation analysis in a Lotka-Volterra competition-diffusion-advection model with time delay effect. *Nonlinearity* **34**(5), 3271–3313.
- <span id="page-30-1"></span>[31] Liao, K.-L. & Lou, Y. (2014) The effect of time delay in a two-patch model with random dispersal. *Bull. Math. Biol.* **76**(2), 335–376.
- <span id="page-30-24"></span>[32] Liu, H., Cong, Y. .& Su, Y. (2022) Dynamics of a two-patch Nicholson's blowflies model with random dispersal. *J. Appl. Anal. Comput.* **12**(2), 692–711.
- <span id="page-30-17"></span>[33] Liu, J. & Chen, S. (2022) Delay-induced instability in a reaction-diffusion model with a general advection term. *J. Math. Anal. Appl.* **512**(2), 20.
- <span id="page-30-2"></span>[34] Lou, Y. (2019) Ideal free distribution in two patches. *J. Nonlinear Model Anal.* **2**, 151–167.
- <span id="page-30-3"></span>[35] Lou, Y. & Lutscher, F. (2014) Evolution of dispersal in open advective environments. *J. Math. Biol.* **69**(6-7), 1319–1342.
- <span id="page-30-4"></span>[36] Lou, Y. & Zhou, P. (2015) Evolution of dispersal in advective homogeneous environment: The effect of boundary conditions. *J. Differ. Eq.* **259**(1), 141–171.
- <span id="page-30-31"></span>[37] Lu, Z. & Takeuchi, Y. (1993) Global asymptotic behavior in single-species discrete diffusion systems. *J. Math. Biol.* **32**(1), 67–77.
- <span id="page-30-18"></span>[38] Ma, L. & Feng, Z. (2021) Stability and bifurcation in a two-species reaction-diffusion-advection competition model with time delay. *Nonlinear Anal. Real World Appl.* **61**(103327), 32.
- <span id="page-30-25"></span>[39] Madras, N., Wu, J. & Zou, X. (1996) Local-nonlocal interaction and spatial-temporal patterns in single species population over a patchy environment. *Can. Appl. Math. Q.* **4**(1), 109–134.
- <span id="page-30-20"></span>[40] Memory, M. C. (1989) Bifurcation and asymptotic behavior of solutions of a delay-differential equation with diffusion. *SIAM J. Math. Anal.* **20**(3), 533–546.
- <span id="page-30-19"></span>[41] Meng, Q., Liu, G. & Jin, Z. (2021) Hopf bifurcation in a reaction-diffusive-advection two-species competition model with one delay. *Electron. J. Qual. Theory Differ. Eq.* **72**(72), 24–24.
- <span id="page-30-9"></span>[42] Noble, L. (2015). Evolution of dispersal in patchy habitats [*PhD thesis*). The Ohio State University.
- <span id="page-30-27"></span><span id="page-30-26"></span>[43] Petit, J., Asllani, M., Fanelli, D., Lauwens, B. & Carletti, T. (2016) Pattern formation in a two-component reaction-diffusion system with delayed processes on a network. *Physica A* **462**, 230–249.
- [44] So, J. W.-H., Wu, J. & Zou, X. (2001) Structured population on two patches: Modeling dispersal and delay. *J. Math. Biol.* **43**(1), 37–51.
- <span id="page-31-1"></span>[45] Speirs, D. C. & Gurney, W. S. C. (2001) Population persistence in rivers and estuaries. *Ecology* **82**(5), 1219–1237.
- <span id="page-31-3"></span>[46] Su, Y., Wei, J. & Shi, J. (2009) Hopf bifurcations in a reaction-diffusion population model with delay effect. *J. Differ. Eq.* **247**(4), 1156–1184.
- <span id="page-31-6"></span>[47] Sun, X. & Yuan, R. (2022) Hopf bifurcation in a diffusive population system with nonlocal delay effect. *Nonlinear Anal.* **214**(112544), 21.
- <span id="page-31-8"></span>[48] Tian, C. & Ruan, S. (2019) Pattern formation and synchronism in an allelopathic plankton model with delay in a network. *SIAM J. Appl. Dyn. Syst.* **18**(1), 531–557.
- <span id="page-31-0"></span>[49] Vasilyeva, O. & Lutscher, F. (2010) Population dynamics in rivers: Analysis of steady states. *Can. Appl. Math. Q.* **18**(4), 439–469.
- <span id="page-31-4"></span>[50] Yan, X.-P. & Li, W.-T. (2010) Stability of bifurcating periodic solutions in a delayed reaction-diffusion population model. *Nonlinearity* **23**(6), 1413–1431.
- <span id="page-31-5"></span>[51] Yan, X.-P. & Li, W.-T. (2012) Stability and Hopf bifurcations for a delayed diffusion system in population dynamics. *Discrete Contin. Dyn. Syst. Ser. B* **17**(1), 367–399.
- <span id="page-31-7"></span>[52] Yoshida, K. (1982) The Hopf bifurcation and its stability for semilinear diffusion equations with time delay arising in ecology. *Hiroshima Math. J.* **12**(2), 321–348.
- <span id="page-31-2"></span>[53] Zhou, P. (2016) On a Lotka-Volterra competition system: Diffusion vs advection. *Calc. Var. Partial Differ. Eq.* **55**(6), 137.

## **Appendix**

In the appendix, we show the following result by linear algebraic techniques.

<span id="page-31-9"></span>**Proposition 5.4.** *Let*  $D = (D_{jk})$  *with*  $D_{jk}$  *defined in* [\(1.3\),](#page-0-2) *and let*  $\widetilde{X}_1$  *be defined in* [\(2.18\).](#page-8-6) Assume that  $D\mathbf{y} = \mathbf{a}$  *with*  $\mathbf{a}, \mathbf{y} \in \widetilde{X}_1$ *. Then* 

<span id="page-31-13"></span>
$$
y_n = \frac{1}{n} \sum_{k=1}^{n-1} k \left( \sum_{j=1}^k a_j \right).
$$
 (5.5)

*Especially, if*  $a_2 = \cdots = a_n$ *, then* 

<span id="page-31-14"></span>
$$
y_n = \frac{n(n-1)}{2}a_1 + \frac{(n-2)(n-1)n}{3}a_2 \tag{5.6}
$$

**Proof.** Since  $Dy = a$  and  $y \in \widetilde{X}_1$ , we have

<span id="page-31-10"></span>
$$
-y_1 + y_2 = a_1,\tag{5.7a}
$$

$$
y_{j-1} - 2y_j + y_{j+1} = a_j, \ j = 2, \cdots, n-1,
$$
 (5.7b)

<span id="page-31-12"></span>and

<span id="page-31-11"></span>
$$
\sum_{j=1}^{n} y_j = 0.
$$
\n(5.8)

Summing the first  $k$  equations in  $(5.7)$ , we find

$$
-y_k + y_{k+1} = \sum_{j=1}^{k} a_j, \ k = 1, \cdots, n-1.
$$
 (5.9)

Multiplying [\(5.9\)](#page-31-11) by *k* and summing these over all *k* yields

$$
-\sum_{j=1}^{n-1} y_j + (n-1)y_n = \sum_{k=1}^{n-1} k\left(\sum_{j=1}^k a_j\right).
$$

This combined with  $(5.8)$  implies that  $(5.5)$  holds.

 $\Box$ 

Now we consider [\(5.6\)](#page-31-14). A direct computation yields

$$
y_n = \left(\sum_{j=1}^{n-1} j\right) a_1 + \left(\sum_{j=1}^{n-2} j(j+1)\right) a_2
$$
  
= 
$$
\frac{n(n-1)}{2} a_1 + \frac{(n-2)(n-1)n}{3} a_2,
$$

where we have used  $\sum_{j=1}^{n} j(j+1) = \frac{n(n+1)(n+2)}{3}$  in the last step. This completes the proof.

**Cite this article:** Liu W., Shen Z. and Chen S. Hopf bifurcations for a delayed discrete single population patch model in advective environments. *European Journal of Applied Mathematics*, <https://doi.org/10.1017/S0956792524000342>