

Limits to Seeing High-Redshift Galaxies Due to Planck-Scale-Induced Blurring

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Keywords. galaxies, gravitation, cosmology: theory

If spacetime is “foamy” travel along a lightpath must be subject to continual, random distance fluctuations $\pm\delta l$ proportional to Planck length $l_P \sim 10^{-35}$ m (Lieu & Hillman 2003). Although each “kick” by itself is tiny, these may accumulate. Accounting for redshifted (bluer) emitted photons, over a cosmological distance $L = (1+z)L_C$ for co-moving distance L_C , the resultant phase perturbations $\Delta\phi = 2\pi\delta l/\lambda$ at observed wavelength λ could grow independently of telescope diameter D to a maximum of $\Delta\phi_{\max} = (1+z)\Delta\phi_0$ (Steinbring 2007) where $\Delta\phi_0 = 2\pi a_0 (l_P^2/\lambda)L^{1-\alpha}$ follows Ng *et al.* (2003). Here $a_0 \sim 1$ and α specifies the quantum-gravity model: 1/2 implies a random walk and 2/3 is consistent with the holographic principle; a vanishingly small $\Delta\phi_P = \Delta\phi_{\max}/[(1+z)a_0(L/l_P)^{1-\alpha}] = 2\pi l_P/\lambda$ is approached when $\alpha = 1$.

The highest-resolution Nyquist-sampled images of $z \approx 6$ active galactic nuclei (AGN) ever obtained are suggestive although do not show clear evidence of $\Delta\phi_{\max}$ for $\alpha = 2/3$ (Steinbring 2007; Tamburini *et al.* 2011). That level of blurring, however, would imply an image full-width at half maximum (FWHM) just at the diffraction limit of the 2.4-m diameter *Hubble Space Telescope* (*HST*). It might be that α is larger or that these phase errors become invisible as they approach the wavelength of observed light (e.g. Perlman *et al.* 2015). But it can be shown by taking a linear superposition of all phase-error amplitudes $\Delta\phi \sigma(\Delta\phi) = 1 - A \log(\Delta\phi/\Delta\phi_P)$ that $(1/A) \int \Delta\phi \sigma(\Delta\phi) d\Delta\phi = (1+z)\Delta\phi_0$ is recovered for $A = 1/\log[(1+z)(L/l_P)^{1-\alpha}]$. If correct, any pointlike source viewed by a diffraction-limited telescope for which $\Delta\phi_{\max}$ is visible must be inflated to a lesser mean width of

$$\Phi = 1.22 \frac{\lambda}{D} + \int \Delta\phi \sigma(\Delta\phi) d\Delta\phi \approx 1.22 \frac{\lambda}{D} + 2\pi \frac{l_P}{\lambda} A e^{1/A}.$$

That is consistent with the *HST* results as well as *Fermi* observations of $z < 4$ gamma-ray bursts (GRBs; Steinbring 2015). Unfortunately, those and current ground-based telescopes employing adaptive optics (AO) are unable to reliably resolve structures much below 0.3 kpc wide at $z = 1$ to 4 (assuming a cosmology with $\Omega_\Lambda = 0.7$, $\Omega_M = 0.3$, and $H_0 = 70$ km s $^{-1}$ Mpc $^{-1}$). Future optical/near-infrared AO and space telescopes viewing out to $z = 6$ to 8 will do better. In a long exposure, the minimal angular extent of known pointlike galaxy structures would be noticeably enlarged - a FWHM expanded by more than a few percent - and possibly also embedded within an extra halo of scattered “fuzz.” For example, with $z = 6.0$, $\alpha = 0.667$, and $D = 6.5$ m no AGN or GRB, despite being physically much smaller, would be found with an apparent size less than about 0.1 kpc in the restframe, 0.025 arcsec across, or consistently 8% larger than diffraction at $0.6 \mu\text{m}$, the shortest wavelength to be visible with the *James Webb Space Telescope*.

References

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