# Supernova Nucleosynthesis with Neutrino Processes: Dependence of Fluorine abundance on Stellar Mass, Explosion Energy and Metallicity

## Natsuko Izutani, Hideyuki Umeda, and Takashi Yoshida

Department of Astronomy, School of Science, University of Tokyo, Bunkyo-ku, Tokyo, 113-0033 email: izutani@astron.s.u-tokyo.ac.jp

**Abstract.** We investigate the effects of neutrino-nucleus interactions on the production of Fluorine during normal supernovae and hypernovae, and discuss stellar mass, metallicity and explosion energy dependence of [F/Fe,Ne,O]. We find the clear trend of [F/Fe,O,Ne] with stellar mass and explosion energy, while no clear trend with metallicity. This trend of [F/O] can be used to constrain the contributed stellar mass by comparing with the observational abundance.

 $\textbf{Keywords.} \ \ \text{neutrinos, nuclear reactions, nucleosynthesis, abundances, supernova: general}$ 

#### 1. Introduction

The interaction of the neutrinos with matter and the effects on the nucleosynthesis have only been discussed for a few models (e.g., Woosley et al. 1990; Woosley & Weaver 1995; Yoshida et al. 2004; Heger et al. 2005; Yoshida et al. 2008; Nakamura et al. 2010). The  $\nu$ -process does not affect the yields of major elements such as Fe and  $\alpha$  elements, but it will increase those of some elements such as B, F, K, Sc, V, and Mn. In this paper, we focus on the effect of the  $\nu$ -process on F during normal supernova (SN) and hypernova (HN) explosions, and discuss stellar mass, metallicity, and explosion energy dependence of [F/Fe,O,Ne].

#### 2. Model & Method

We calculate the nucleosynthesis for core-collapse SNe with progenitor masses of  $M=15, 25, \, {\rm and} \, 50 \, M_{\odot}$  and initial metallicities of  $Z=0, \, 0.004, \, {\rm and} \, 0.02$  for normal SNe and HNe. The explosion energy is set to be  $1\times 10^{51}$  ergs for normal SNe,  $10\times 10^{51}$  and  $40\times 10^{51}$  ergs for HNe with  $M_{\rm MS}=25$  and  $50\, M_{\odot}$ , respectively. For normal SNe, the mass cut is set to meet the observed iron mass of  $0.07\, M_{\odot}$ . For HNe, the parameters of mixing fallback models are determined to get  ${\rm [O/Fe]}=0.5$ . The nuclear network includes 809 species up to  $^{121}{\rm Pd}$  (Izutani et al. 2009, Izutani & Umeda 2010). We adopt the  $\nu$ -process up to  $^{80}{\rm Kr}$  as in Yoshida et al. (2008). The neutrino luminosity is assumed to be uniformly partitioned among the neutrino flavors, and decrease exponentially in time with a timescale of 3 s. The total neutrino energy is set to be  $E_{\nu}=3$  and  $9\times 10^{53}$  ergs. The neutrino energy spectra are assumed to be Fermi-Dirac distributions with zero chemical potentials. The temperatures of  $\nu_{\mu,\tau}, \, \bar{\nu}_{\mu,\tau}$  and  $\nu_{\rm e}, \, \bar{\nu}_{\rm e}$  are set to be  $T_{\nu}=6$  and  $4\,{\rm MeV}/k$ , respectively.

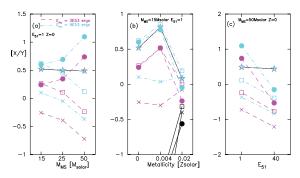


Figure 1. [F/Fe,O,Ne] (filled circles, open squares, and crosses) and [Ne/O] (stars) in the models without  $\nu$ -processes (black solid lines), with  $\nu$ -processes of  $E_{\nu}=3$  and  $9\times 10^{53}$  ergs (magenta dashed lines and cyan dot-dashed lines). (a) Stellar mass dependence ([X/Y] in the models with  $M_{\rm MS}=15,\,25,\,50\,M_{\odot},\,Z=0$  and  $E_{51}=1$ ) (b) Metallicity dependence ([X/Y] in the models with  $Z=0,\,0.004,\,0.02,\,M_{\rm MS}=15\,M_{\odot}$  and  $E_{51}=1$ ) (c)  $E_{51}$  dependence ([X/Y] in the models with  $M_{\rm MS}=50\,M_{\odot},\,Z=0,\,E_{51}=1$  and 40.

### 3. Results and Discussion

With the  $\nu$ -process, <sup>19</sup>F is produced in the O/Ne-enriched region through <sup>20</sup>Ne( $\nu$ ,  $\nu$ 'p)<sup>19</sup>F. Figure 1 (a) shows mass dependence of [F/Fe,O,Ne] and [Ne/O] in Z=0 star SNe. [Ne/O] is about 0.5 in these models. Without the  $\nu$ -processes, [F/Fe,O,Ne] are  $\sim -5$ . With the  $\nu$ -processes, [F/Fe,O,Ne] range from -1 to 1. [F/Fe] is higher for more massive stars because of the larger O/Ne-enriched region. By contrast, [F/O,Ne] are lower for more massive stars because the radius of the O/Ne-enriched region is larger, and the neutrino flux becomes smaller. Figure 1(b) shows metallicity dependence of [F/Fe,O,Ne] and [Ne/O] in 15  $M_{\odot}$  SNe. [Ne/O] is different between these models, though it is not clear whether this trend of [Ne/O] is due to metallicity or not. With the  $\nu$ -processes, the F yield is increased by a factor of  $\sim 10$  and 1000 for Z = 0.02 and 0, respectively. There is no clear trend of [F/Fe,O,Ne] with metallicity. Figure 1 (c) shows  $E_{\text{exp}}$  dependence of [F/Fe,O,Ne]and [Ne/O] in  $Z=0.50M_{\odot}$  explosions. [Ne/O] is about 0.5 in these models. With the  $\nu$ processes, [F/Fe,O,Ne] range from -1 to 1. [F/O,Ne] are lower in the HN model. In the HN model, the shock wave reaches the O/Ne-enriched region earlier, and the region expands earlier, which causes smaller neutrino flux. [F/Fe] is also lower in the HN model, which is caused by both the larger mass of Fe and the smaller mass of F in the HN model.

It is true that  $\nu$ -cross-sections contain some uncertainties. Nevertheless, the trend of [F/Fe,O,Ne] discussed above is robust for these uncertainties. For the galactic chemical evolution calculation using these yields, see Kobayashi *et al.* (2011).

#### References

Heger, A., Kolbe, E., Haxton, W. C., et al. 2005, Phys. Lett. B. 606, 258

Izutani, N., Umeda, H., & Tominaga, N. 2009, ApJ, 692, 1517

Izutani, N. & Umeda, H. 2010, ApJ, 720, L1

Kobayashi, C., Izutani, N., Karakas, A. I., Yoshida, T., Yong, D., & Umeda, H. 2011, ApJ, 739, L57

Nakamura, K., Yoshida, T., Shigeyama, T., & Kajino, T. 2010, ApJ, 718, L137

Umeda, H. & Nomoto, K., 2002, ApJ, 565, 385

Woosley, S. E., Hartmann, D. H., Hoffman, R. D., & Haxton, W. C. 1990, ApJ, 356, 272

Woosley, S. E. & Weaver, T. A. 1995, ApJS, 101, 181

Yoshida, T., Terasawa, M., Kajino, T., & Sumiyoshi, K. 2004 ApJ, 600, 204

Yoshida, T., Umeda, H., & Nomoto, K. 2008, ApJ, 672, 1043