

A LARGE DEVIATION PRINCIPLE FOR A BROWNIAN IMMIGRATION PARTICLE SYSTEM

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Abstract

We derive a large deviation principle for a Brownian immigration branching particle system, where the immigration is governed by a Poisson random measure with a Lebesgue intensity measure.

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1. Introduction

Consider a particle system in \mathbb{R}^d . Particles are initially distributed according to a Poisson random field with intensity measure μ . Each of these particles undergoes Brownian motion until it either splits into two particles or disappears at an exponential rate. For a bounded measurable set $A \subset \mathbb{R}^d$, let $M_t(A)$ denote the number of particles in A at time t . Define

$$\langle M_t, f \rangle = \sum_{x \in \text{supp } M_t} f(x), \quad f \in L^1(\mathbb{R}^d).$$

We call $(M_t)_{t \geq 0}$ a *Brownian critical binary branching particle system*; see Dawson (1993). Consider a situation in which there are additional sources of particles from which immigration occurs during the evolution. The immigration time and sites are determined by a Poisson random field on $[0, \infty) \times \mathbb{R}^d$ with a Lebesgue intensity measure. After arriving, each of these particles propagates and moves in \mathbb{R}^d in the same way as the other particles. Let N_t denote the empirical measure of the immigration particle system at t . The process $(N_t)_{t \geq 0}$ is called a *Brownian immigration branching particle system*; see Li (1998).

The large and moderate deviation principles (LDPs and MDPs) for Brownian particle systems and super-Brownian motion have been studied by several authors; see, for example, Cox and Griffeath (1985), Deuschel and Wang (1994), Deuschel and Rosen (1998), Iscoe and Lee (1993), Lee (1993), and Hong (2003). In particular, Deuschel and Wang (1994) studied the LDP for the occupation time process of a Poisson system of independent Brownian particles without branching. The LDP for the occupation time process of branching Brownian motion was studied by Cox and Griffeath (1985). Iscoe and Lee (1993) and Lee (1993) obtained the LDPs for occupation processes of both a Brownian branching particle system and its measure-valued version. In Zhang (2004a), (2004b), the author studied the LDP and MDP for super-Brownian motion with immigration, where the speed function is $t^{1/2}$ for $d = 1$, $t/\log t$ for $d = 2$, and t for $d \geq 3$.

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In this paper, we are interested in the LDP for the Brownian immigration branching particle system $(N_t)_{t \geq 0}$. Suppose that $N_0 = 0$, i.e. there are no particles at the initial time. For a bounded integrable function $f \geq 0$, we have

$$\frac{1}{T} \langle N_T, f \rangle \rightarrow \int_{\mathbb{R}^d} f(x) dx, \quad T \rightarrow \infty.$$

We shall study the LDP based on this central tendency for $d = 1$. In this case, the speed function is $t^{1/2}$, as in the case of super-Brownian motion with immigration (see Zhang (2004a)). However, in the present paper we obtain the complete LDP, while in Zhang (2004a) only a local LDP (in a neighborhood of $\int f(x) dx$) was proved. These discussions can also be applied to super-Brownian motion with immigration, so the results of Zhang (2004a) can be modified and the complete LDP holds.

We introduce some notation before we state our results. If u is a Borel measurable function on $[0, \infty) \times \mathbb{R}^d$, we shall often suppress a variable and write $u(t)$ for the function whose value at x is $u(t, x)$. If u is differentiable, we simply write $\partial u(t)/\partial t$ for $\partial u(t, x)/\partial t$, and denote by

$$\Delta u(t) \equiv \Delta u(t, x) := \sum_{i=1}^d \frac{\partial^2}{\partial x_i^2} u(t, x)$$

the Laplacian of u .

Let $C(\mathbb{R}^d)$ denote the space of bounded continuous functions on \mathbb{R}^d . We fix a constant $p > d$ and let $\phi_p(x) := (1 + |x|^2)^{-p/2}$ for $x \in \mathbb{R}^d$. Let

$$C_p(\mathbb{R}^d) := \{f \in C(\mathbb{R}^d) : |f(x)| \leq \text{const.} \times \phi_p(x)\}.$$

The nonnegative subset of $C_p(\mathbb{R}^d)$ will be denoted by $C_p^+(\mathbb{R}^d)$. We denote by $M(\mathbb{R}^d)$ the set of all positive Radon measures on the Borel σ -algebra of \mathbb{R}^d . Let $M_p(\mathbb{R}^d) \subset M(\mathbb{R}^d)$ be the set of μ such that

$$\langle \mu, f \rangle := \int f(x) \mu(dx) < \infty \quad \text{for all } f \in C_p(\mathbb{R}^d).$$

We endow $M_p(\mathbb{R}^d)$ with the p -vague topology: a sequence $\{\mu_k\} \subset M_p(\mathbb{R}^d)$ converges in this topology to $\mu \in M_p(\mathbb{R}^d)$ if and only if $\langle \mu_k, f \rangle \rightarrow \langle \mu, f \rangle$ for all $f \in C_p(\mathbb{R}^d)$. Note that the Lebesgue measure, which will always be denoted by λ , belongs to $M_p(\mathbb{R}^d)$ for $p > d$. Let $\|\cdot\|$ denote the usual supremum norm on \mathbb{R}^d . For any Lebesgue square-integrable function φ on $[0, 1] \times \mathbb{R}^d$, define its L^2 -norm by

$$\|\varphi(\cdot, \cdot)\|_{L^2([0,1] \times \mathbb{R}^d)} := \left(\int_0^1 dt \int_{\mathbb{R}^d} \varphi^2(t, x) dx \right)^{1/2}.$$

If $(X_t)_{t \geq 0}$ is a Markov process with state space $M_p(\mathbb{R}^d)$, we will denote by P_μ the probability measure such that $P_\mu(X_0 = \mu) = 1$. Expectation with respect to P_μ will be denoted by E_μ , and P_0 and E_0 will be abbreviated to P and E , respectively.

Suppose that $(P_t)_{t \geq 0}$ is the transition semigroup of Brownian motion in \mathbb{R}^d and $p_t(x)$ is its density function: $p_t(x) = (4\pi t)^{-d/2} \exp\{-|x|^2/4t\}$. Suppose that $(M_t)_{t \geq 0}$ is the Brownian critical binary branching particle system introduced at the beginning of the paper. For $\mu \in$

$M_p(\mathbb{R}^d)$, let $E_{(\mu)}$ denote the conditional law of $(M_t)_{t \geq 0}$ given that M_0 is a Poisson random measure with intensity μ . Then

$$E_{(\mu)} \exp\{-\langle M_t, f \rangle\} = \exp\{-\langle \mu, v(t, \cdot) \rangle\}, \quad f \in C_p^+(\mathbb{R}^d), \tag{1.1}$$

where $v \equiv v(t, x)$ is the unique mild solution to

$$\frac{\partial v(t)}{\partial t} = \Delta v(t) - v^2(t), \quad v(0) = 1 - e^{-f}; \tag{1.2}$$

see Dawson (1993). A Brownian immigration branching particle system $(N_t)_{t \geq 0}$ is an $M_p(\mathbb{R}^d)$ -valued Markov process with $N_0 = 0$ and Laplace transform given by

$$E \exp\{-\langle N_t, f \rangle\} = \exp\left\{-\int_0^t \langle \lambda, v(s, \cdot) \rangle ds\right\}, \quad f \in C_p^+(\mathbb{R}^d);$$

see Li (1998).

Suppose that $f \in C_p^+(\mathbb{R}^d)$ and $\langle \lambda, f \rangle = 1$. For $d = 1$ and $T > 0$, define

$$\Lambda(T, \theta) = T^{-1/2} \log E \exp\{\theta T^{-1/2} \langle N_T, f \rangle\}. \tag{1.3}$$

We will prove, for some $\theta \in (-\infty, \theta_0)$, $\theta_0 > 0$, that

$$\Lambda(\theta) := \lim_{T \rightarrow \infty} \Lambda(T, \theta) = \theta + \int_0^1 ds \int_0^s \langle \lambda, [V(r, \cdot; \theta)]^2 \rangle dr, \tag{1.4}$$

where $V(\cdot, \cdot; \theta)$ is the unique solution to the singular PDE

$$\frac{\partial V(s)}{\partial s} = \Delta V(s) + V^2(s), \quad 0 \leq s \leq 1, \quad V(0) = \theta \delta_0.$$

Here δ_0 is the Dirac mass at 0. Moreover, $\lim_{\theta \uparrow \theta_0} \Lambda'(\theta) = \infty$, where a prime denotes differentiation. Let $I(a)$ be the Legendre transform of $\Lambda(\theta)$:

$$I(a) = \sup_{\theta \in (-\infty, \theta_0)} [a\theta - \Lambda(\theta)], \quad a \in \mathbb{R}.$$

Theorem 1.1. *For $d = 1$, if $U \subset \mathbb{R}$ is open and $C \subset \mathbb{R}$ is closed, then*

$$\liminf_{T \rightarrow \infty} T^{-1/2} \log P\left\{\frac{N_T}{T} \in U\right\} \geq - \inf_{a \in U} I(a)$$

and

$$\limsup_{T \rightarrow \infty} T^{-1/2} \log P\left\{\frac{N_T}{T} \in C\right\} \leq - \inf_{a \in C} I(a).$$

2. Proofs

In this section we prove Theorem 1.1. We need several supporting lemmas.

Lemma 2.1. *There exists a $\theta_0 > 0$ such that, for $\theta \in (-\infty, \theta_0)$,*

$$\frac{\partial V(t)}{\partial t} = \Delta V(t) + V^2(t), \quad 0 \leq t \leq 1, \quad V(0) = \theta \delta_0, \tag{2.1}$$

has a unique solution $V \equiv V(t, x; \theta)$ satisfying

$$\lim_{\theta \uparrow \theta_0} \|V(\cdot, \cdot; \theta)\|_{L^2([0,1] \times \mathbb{R})} = \infty. \tag{2.2}$$

Proof. For $\theta \leq 0$, the assertion follows from Kamin and Peletier (1985, p. 205). We discuss the case $\theta > 0$. In Zhang (2004a) (see Corollary 3.1 thereof), the author proved that the unique solution of (2.1) exists in $L^2([0, 1] \times \mathbb{R})$ for $\theta \in (0, \sqrt{\pi}/4)$. Hence, it remains to determine the boundary point θ_0 . For $l > 0$, consider the following equation:

$$\frac{\partial w(s)}{\partial s} = \Delta w(s) + w^2(s), \quad s < l, \quad w(0) = \delta_0. \tag{2.3}$$

Clearly, the solutions to (2.3) (which we denote by $w \equiv w(s, x; \delta_0)$) and (2.1) are related by

$$w(s, x; \delta_0) = \theta^{-2} V(\theta^{-2}s, \theta^{-1}x; \theta).$$

Thus, (2.3) has a unique solution in $L^2([0, l] \times \mathbb{R})$ if $l < \pi/16$. Define

$$c = \sup\{t : (2.3) \text{ has a unique solution (still denoted by } w(\cdot, \cdot; \delta_0)) \text{ in } L^2([0, t] \times \mathbb{R})\}.$$

Therefore, $c \geq \pi/16$. In the following, we prove that $c < \infty$. Fix $t < c$ and define

$$\bar{w}(s, x) = \int p_{t-s}(x, y) w(s, y; \delta_0) dy, \quad s < t.$$

By (2.3), we have

$$w(s, x; \delta_0) = p_s(x) + \int_0^s P_{s-r}[w^2(r, \cdot; \delta_0)](x) dr$$

and, so, with $p_t(x, y) := p_t(x - y)$,

$$\begin{aligned} \bar{w}(s, x) &= p_t(x) + \int p_{t-s}(x, y) dy \int_0^s P_{s-r}[w^2(r, \cdot; \delta_0)](y) dr \\ &= p_t(x) + \int_0^s dr \int p_{t-s}(x, y) P_{s-r}[w^2(r, \cdot; \delta_0)](y) dy \\ &= p_t(x) + \int_0^s P_{t-r}[w^2(r, \cdot; \delta_0)](x) dr. \end{aligned} \tag{2.4}$$

For $s < t < c$ and $x \in \mathbb{R}$, and noting the definition of c , we have

$$\begin{aligned} \bar{w}(s, x) &\leq p_t(x) + \int_0^s [4\pi(t - r)]^{-1/2} dr \int w^2(r, x; \delta_0) dx \\ &\leq p_t(x) + [4\pi(t - s)]^{-1/2} \int_0^s dr \int w^2(r, x; \delta_0) dx \\ &< \infty. \end{aligned}$$

On the other hand, by (2.4) and the Schwarz inequality, we obtain

$$\begin{aligned} \bar{w}(s, x) &\geq p_t(x) + \int_0^s (P_{t-r}[w(r, \cdot; \delta_0)](x))^2 dr \\ &= p_t(x) + \int_0^s \bar{w}^2(r, x) dr, \end{aligned}$$

which implies that $\bar{w}(s, x)$ is a super solution to

$$\frac{\partial W(s, x)}{\partial s} = W^2(s, x), \quad s < t, \quad W(0, x) = p_t(x).$$

(That is, $\bar{w}(s, x) \geq W(s, x)$ for any $s < t$ and $x \in \mathbb{R}$.) Clearly,

$$W(s, x) = \frac{p_t(x)}{1 - sp_t(x)}.$$

Choose $t_0 = 64\pi$ and $x = 16\pi\sqrt{\log 2}$, meaning that $\frac{1}{2}t_0 p_{t_0}(x) = 1$. Let $s = \frac{1}{2}t_0$. Then $\bar{w}(s, x) \geq W(s, x) = \infty$, which implies that $c \leq t_0 < \infty$. Let

$$V(t, x; \theta) := \theta^2 w(\theta^2 t, \theta x; \delta_0), \quad 0 < t \leq 1, x \in \mathbb{R}. \tag{2.5}$$

Then $V(t, x; \theta)$ is the solution of (2.1) for each $\theta < \theta_0 := \sqrt{c}$. To see (2.2), note that

$$\lim_{\theta \uparrow \theta_0} \|V(\cdot, \cdot; \theta)\|_{L^2([0,1] \times \mathbb{R})} = \theta_0 \int_0^c ds \int w^2(s, y; \delta_0) dy.$$

By the definition of c , we conclude that $\int_0^c ds \int w^2(s, y; \delta_0) dy = \infty$, and obtain (2.2).

Otherwise, if

$$\int_0^c ds \int w^2(s, y; \delta_0) dy < \infty$$

then let $0 < \Delta c < 1$. For $t \in [c, c + \Delta c]$ and $n = 0, 1, 2, \dots$, define

$$\begin{aligned} \bar{w}_{n+1}(t, x; \delta_0) &= P_{t-c}[w(c, \cdot; \delta_0)](x) + \int_c^t P_{t-r}[\bar{w}_n^2(r, \cdot; \delta_0)](x) dr, \\ \bar{w}_0(t, x; \delta_0) &= P_{t-c}[w(c, \cdot; \delta_0)](x). \end{aligned} \tag{2.6}$$

Define the L^2 -norm of a function $g(\cdot, \cdot)$ on $[c, c + \Delta c] \times \mathbb{R}$ by

$$\|g\|_{L^2}^2 = \int_c^{c+\Delta c} dt \int g^2(t, x) dx,$$

if this integral is finite. In analogy with (2.4), we have

$$\bar{w}_0(t, x) = p_t(x) + \int_0^c P_{t-r}[w^2(r, \cdot; \delta_0)](x) dr, \quad t \in [c, c + \Delta c].$$

By the C_r -inequality, this implies that

$$\begin{aligned} \|\bar{w}_0\|_{L^2}^2 &\leq 2 \int_c^{c+\Delta c} dt \int p_t^2(x) dx + 2 \int_c^{c+\Delta c} dt \int \left(\int_0^c P_{t-r}[w^2(r, \cdot; \delta_0)](x) dr \right)^2 dx \\ &= 2 \int_c^{c+\Delta c} p_{2t}(0) dt + 2 \int_c^{c+\Delta c} dt \int_0^c dr \int_0^c du \\ &\quad \times \iint p_{2t-r-u}(y, z) w^2(r, y; \delta_0) w^2(u, z; \delta_0) dy dz \\ &\leq \sqrt{\frac{2}{\pi}} (\sqrt{c + \Delta c} - \sqrt{c}) + 2 \int_c^{c+\Delta c} [4\pi(2t - 2c)]^{-1/2} dt \\ &\quad \times \int_0^c dr \int w^2(r, y; \delta_0) dy \int_0^c du \int w^2(u, z; \delta_0) dz. \end{aligned}$$

Let $A := \int_0^c ds \int w^2(s, y; \delta_0) dy < \infty$. We then find that

$$\|\bar{w}_0\|_{L^2}^2 \leq \sqrt{\frac{2}{\pi}}(\sqrt{c + \Delta c} - \sqrt{c}) + \sqrt{\frac{2}{\pi}}A^2\sqrt{\Delta c} \leq k_1\sqrt{\Delta c},$$

where k_1 is a constant depending only on c and A . In the following, we suppose that $k_1 > 1$. By the C_r -inequality and (2.6), we similarly find that

$$\begin{aligned} & \|\bar{w}_{n+1}\|_{L^2}^2 \\ & \leq 2\|w_0\|_{L^2}^2 + 2 \int_c^{c+\Delta c} dt \int \left(\int_c^t P_{t-r}[\bar{w}_n^2(r, \cdot; \delta_0)](x) dr \right)^2 dx \\ & \leq 2k_1\sqrt{\Delta c} + 2 \int_c^{c+\Delta c} dt \int_c^t dr \int_c^t du \iint p_{2t-r-u}(y, z) \bar{w}_n^2(r, y; \delta_0) \bar{w}_n^2(u, z; \delta_0) dy dz \\ & \leq 2k_1\sqrt{\Delta c} + 2 \int_c^{c+\Delta c} dt \int_c^t dr \int_c^t [4\pi(2t - r - u)]^{-1/2} du \\ & \quad \times \iint \bar{w}_n^2(r, y; \delta_0) \bar{w}_n^2(u, z; \delta_0) dy dz \\ & \leq 2k_1\sqrt{\Delta c} + \frac{1}{\sqrt{\pi}} \int_c^{c+\Delta c} dt \int_c^t (t - r)^{-1/2} dr \int \bar{w}_n^2(r, y; \delta_0) dy \int_c^t du \\ & \quad \times \int \bar{w}_n^2(u, z; \delta_0) dz \\ & \leq 2k_1\sqrt{\Delta c} + \frac{1}{\sqrt{\pi}} \|\bar{w}_n\|_{L^2}^2 \int_c^{c+\Delta c} dt \int_c^t (t - r)^{-1/2} dr \int \bar{w}_n^2(r, y; \delta_0) dy \\ & \leq 2k_1\sqrt{\Delta c} + \frac{1}{\sqrt{\pi}} \|\bar{w}_n\|_{L^2}^2 \int_c^{c+\Delta c} dr \int_r^{c+\Delta c} (t - r)^{-1/2} dt \int \bar{w}_n^2(r, y; \delta_0) dy \\ & \leq 2k_1\sqrt{\Delta c} + \frac{1}{\sqrt{\pi}} \|\bar{w}_n\|_{L^2}^2 \int_c^{c+\Delta c} 2\sqrt{c + \Delta c - r} dr \int \bar{w}_n^2(r, y; \delta_0) dy \\ & \leq 2k_1\sqrt{\Delta c} + \frac{2}{\sqrt{\pi}} \sqrt{\Delta c} \|\bar{w}_n\|_{L^2}^4 \\ & \leq k_2\sqrt{\Delta c}(1 + \|\bar{w}_n\|_{L^2}^4), \tag{2.7} \end{aligned}$$

where $k_2 > 0$ is a constant depending only on k_1 and c . Choose Δc , $0 < \Delta c < 1$, such that $(k_1 + 1)k_2\sqrt{\Delta c} \leq 1\sqrt{1 + \Delta c}$, and note that

$$\|\bar{w}_0\|_{L^2}^2 \leq k_1\sqrt{\Delta c} \leq k_1\sqrt{1 + \Delta c}.$$

Suppose that $\|\bar{w}_n\|_{L^2}^2 \leq k_1\sqrt{1 + \Delta c}$. Recalling that we have assumed $k_1 > 1$, by (2.7) we have

$$\begin{aligned} \|\bar{w}_{n+1}\|_{L^2}^2 & \leq k_2\sqrt{\Delta c}[1 + k_1^2(1 + \Delta c)] \\ & \leq k_2(k_1^2 + 1)\sqrt{\Delta c}(1 + \Delta c) \\ & \leq k_2(k_1^2 + k_1)\sqrt{\Delta c}(1 + \Delta c) \\ & \leq k_1\sqrt{1 + \Delta c} \\ & < \infty. \end{aligned}$$

From (2.6), we know that \bar{w}_n is increasing in n . Therefore, $\bar{w} := \lim_{n \rightarrow \infty} \bar{w}_n$ exists in $L^2([c, c + \Delta c] \times \mathbb{R})$ and satisfies (2.3) on the interval $[c, c + \Delta c]$.

It is easy to check that

$$\tilde{w}(t, x; \delta_0) = \begin{cases} w(t, x; \delta_0), & t \in [0, c], \\ \bar{w}(t, x; \delta_0), & t \in [c, c + \Delta c], \end{cases} \tag{2.8}$$

is a solution to (2.3) that remains in $L^2([c, c + \Delta c] \times \mathbb{R})$. Let us see the uniqueness of this solution. Suppose that \tilde{w}_1 is another solution to (2.3) in $L^2([c, c + \Delta c] \times \mathbb{R})$. Then

$$\begin{aligned} \tilde{w}(t, x; \delta_0) &= P_{t-c}[w(c, \cdot; \delta_0)](x) + \int_c^t P_{t-r}[\tilde{w}^2(r, \cdot; \delta_0)](x) \, dr, & t \in [c, c + \Delta c], \\ \tilde{w}_1(t, x; \delta_0) &= P_{t-c}[w(c, \cdot; \delta_0)](x) + \int_c^t P_{t-r}[\tilde{w}_1^2(r, \cdot; \delta_0)](x) \, dr, & t \in [c, c + \Delta c]. \end{aligned}$$

By using the Schwarz inequality and arguing as in (2.7), we find that

$$\begin{aligned} J &:= \int_c^{c+\Delta c} dt \int |\tilde{w}(t, x; \delta_0) - \tilde{w}_1(t, x; \delta_0)|^2 \, dx \\ &= \int_c^{c+\Delta c} dt \int \left(\int_c^t P_{t-r} [|\tilde{w}^2(r, \cdot; \delta_0) - \tilde{w}_1^2(r, \cdot; \delta_0)|](x) \, dr \right)^2 \, dx \\ &\leq \frac{2}{\sqrt{\pi}} \|\tilde{w}\|_{L^2}^2 \|\tilde{w}_1\|_{L^2}^2 \int_c^{c+\Delta c} \sqrt{c + \Delta c - r} \, dr \int |\tilde{w}(r, y; \delta_0) - \tilde{w}_1(r, y; \delta_0)|^2 \, dy \\ &\leq \frac{2}{\sqrt{\pi}} \|\tilde{w}\|_{L^2}^2 \|\tilde{w}_1\|_{L^2}^2 \sqrt{\Delta c} J. \end{aligned}$$

Since both \tilde{w} and \tilde{w}_1 belong to $L^2([c, c + \Delta c] \times \mathbb{R})$, we can choose $\Delta c > 0$ such that

$$\frac{2}{\sqrt{\pi}} \|\tilde{w}\|_{L^2}^2 \|\tilde{w}_1\|_{L^2}^2 \sqrt{\Delta c} < 1.$$

Then we have

$$J = \int_c^{c+\Delta c} dt \int |\tilde{w}(t, x; \delta_0) - \tilde{w}_1(t, x; \delta_0)|^2 \, dx = 0.$$

Therefore, if

$$\int_0^c ds \int w^2(s, y; \delta_0) \, dy < \infty,$$

then by (2.8) we can construct the unique solution of (2.3) in $L^2([0, c + \Delta c] \times \mathbb{R})$. This contradicts the definition of c , and the proof is complete.

Corollary 2.1. *For each $\theta \in (-\infty, \theta_0)$, there exists an $a > 1$ such that*

$$\frac{\partial U(t)}{\partial t} = \Delta U(t) + aU^2(t), \quad 0 < t \leq 1, \quad U(0) = \theta\delta_0, \tag{2.9}$$

has a unique solution $U(\cdot, \cdot; \theta) \in L^2([0, 1] \times \mathbb{R})$.

Proof. Set $\tilde{U} = aU$, where U is the solution to (2.9). Then \tilde{U} is the solution to

$$\frac{\partial \tilde{U}(t)}{\partial t} = \Delta \tilde{U}(t) + \tilde{U}^2(t), \quad 0 < t \leq 1, \quad \tilde{U}(0) = a\theta\delta_0. \tag{2.10}$$

For each $\theta \in (-\infty, \theta_0)$, we can choose an $a > 1$ such that $a\theta \in (-\infty, \theta_0)$. By Lemma 2.1, (2.10) has a unique solution \tilde{U} , and the assertion follows.

Lemma 2.2. *Suppose that $f \in C_p^+(\mathbb{R}^d)$ with $\langle \lambda, f \rangle = 1$. Let*

$$f_T(x, \theta) = T(\exp\{\theta T^{-1/2} f(T^{1/2}x)\} - 1).$$

For $\theta \in (-\infty, \theta_0)$, the equation

$$\frac{\partial V_T(t)}{\partial t} = \Delta V_T(t) + V_T^2(t), \quad 0 \leq t \leq 1, \quad V_T(0) = f_T(\cdot, \theta), \tag{2.11}$$

has a unique solution $V_T(\cdot, \cdot; \theta)$, which converges to $V(\cdot, \cdot; \theta)$ in $L^2([0, 1] \times \mathbb{R})$ as $T \rightarrow \infty$. If we allow θ to be a complex number such that $|\theta| < \theta_0$, then $V_T(\cdot, \cdot; \theta)$ is analytic in $|\theta| < \theta_0$.

To prove Lemma 2.2 we need the following result.

Lemma 2.3. *As $T \rightarrow \infty$,*

$$\int_0^1 dt \int (P_t[f_T(\cdot, \theta)](x) - \theta p_t(x))^2 dx \rightarrow 0.$$

Proof. If $\theta = 0$ the assertion is obvious. For $\theta \neq 0$,

$$\begin{aligned} \int_0^1 dt \int (P_t[f_T(\cdot, \theta)](x) - \theta p_t(x))^2 dx &= \int_0^1 dt \iint p_{2t}(x, y) f_T(x, \theta) f_T(y, \theta) dx dy \\ &\quad - 2\theta \int_0^1 dt \int p_{2t}(x, 0) f_T(x, \theta) dx \\ &\quad + \theta^2 \int_0^1 p_{2t}(0) dt. \end{aligned} \tag{2.12}$$

For the first term on the right-hand side of (2.12), we have

$$\begin{aligned} I_1 &:= \int_0^1 dt \iint p_{2t}(x, y) f_T(x, \theta) f_T(y, \theta) dx dy \\ &= \int_0^1 dt \iint p_{2t}(T^{-1/2}(w - u)) T(e^{\theta T^{-1/2} f(w)} - 1)(e^{\theta T^{-1/2} f(u)} - 1) dw du. \end{aligned}$$

After simple calculation, we find that

$$|e^x - 1| < \begin{cases} a|x|, & 0 < x < \log a, \ a > 1, \\ |x|, & x \leq 0. \end{cases} \tag{2.13}$$

For $T > T_0(\theta, f) := (\theta f)^2(\log 2)^{-2}$ we have

$$|\theta T^{-1/2} f(x)| \leq |\theta| \|f\| T^{-1/2} < \log 2.$$

Then

$$T^{1/2}|(e^{\theta T^{-1/2} f(w)} - 1)| \leq 2|\theta f(w)|.$$

By applying the dominated convergence theorem to I_1 , and recalling that $\langle \lambda, f \rangle = 1$, we find that

$$I_1 \rightarrow \theta^2 \int_0^1 p_{2t}(0) dt < \infty.$$

Similarly, we obtain

$$I_2 := 2\theta \int_0^1 dt \int p_{2t}(x, 0) f_T(x, \theta) dx \rightarrow 2\theta^2 \int_0^1 p_{2t}(0) dt.$$

This completes the proof.

Proof of Lemma 2.2. For $\theta \leq 0$, by replacing f by $-\theta T^{-1/2} f(T^{1/2} \cdot)$ in (1.2) and making the change of variable $V(t) \leftrightarrow -V(t)$, we obtain (2.11). For $0 < \theta < \theta_0$, define

$$\bar{V}_T(s, x, \theta) = \int \frac{1}{\theta} f_T(x - y, \theta) U(s, y; \theta) dy,$$

where $U(\cdot, \cdot; \theta)$ is the solution to (2.9). By (2.1) and the Schwarz inequality, we have

$$\begin{aligned} \bar{V}_T(s, x, \theta) &= P_s[f_T(\cdot, \theta)](x) + a \int_0^s dr \int \frac{1}{\theta} f_T(x - y, \theta) P_{s-r}[U^2(r, \cdot; \theta)](y) dy \\ &= P_s[f_T(\cdot, \theta)](x) + \frac{a}{\theta} \int_0^s dr \iint f_T(x - y, \theta) p_{s-r}(y, z) U^2(r, z; \theta) dy dz \\ &= P_s[f_T(\cdot, \theta)](x) + \frac{a}{\theta} \int_0^s dr \iint p_{s-r}(w) f_T(x - w - z, \theta) U^2(r, z; \theta) dw dz \\ &\geq P_s[f_T(\cdot, \theta)](x) + \frac{a}{\theta} \int_0^s dr \int p_{s-r}(w) dw \left(\int f_T(x - w - z, \theta) U(r, z; \theta) dz \right)^2 \\ &\quad \times \left(\int f_T(x - w - z, \theta) dz \right)^{-1} \\ &\geq P_s[f_T(\cdot, \theta)](x) + \frac{1}{\theta^2} \int_0^s dr \int p_{s-r}(w) dw \left(\int f_T(x - w - z, \theta) U(r, z; \theta) dz \right)^2 \\ &= P_s[f_T(\cdot, \theta)](x) + \int_0^s dr \int P_{s-r}[\bar{V}_T^2(r, \cdot, \theta)] dr. \end{aligned}$$

In the fifth step, we have used (2.13), which yields

$$\begin{aligned} \left(\int f_T(x - w - z, \theta) dz \right)^{-1} &= \left(T \int (\exp\{\theta T^{-1/2} f(T^{1/2}(x - w - z))\} - 1) dz \right)^{-1} \\ &\geq \left(a\theta T^{1/2} \int f(T^{1/2}(x - w - z)) dz \right)^{-1} \\ &= (a\theta)^{-1} \end{aligned}$$

when T is sufficiently large.

Thus, \bar{V}_T is a super solution of (2.11), i.e.

$$\bar{V}_T(s, x; \theta) \geq V_T(s, x; \theta)$$

for any $(s, x) \in [0, 1] \times \mathbb{R}$. By standard results on differential equations, we conclude that (2.11) has a unique solution V_T for $\theta \in (0, \theta_0)$.

In the following we discuss the limit of V_T . By (2.11) and (2.1) we have, for $0 < \varepsilon < 1$,

$$\begin{aligned} I(T) &:= \int_0^\varepsilon dt \int |V_T(t, x; \theta) - V(t, x; \theta)|^2 dx \\ &\leq 2 \int_0^\varepsilon dt \int (P_t[f_T(x, \theta)] - \theta p_t(x))^2 dx \\ &\quad + 2 \int_0^\varepsilon dt \int \left(\int_0^t P_{t-r}[|V_T^2(r) - V^2(r)|](x) dr \right)^2 dx \\ &=: 2I_1(T) + 2I_2(T), \end{aligned} \tag{2.14}$$

where

$$I_1(T) = \int_0^\varepsilon dt \int (P_t[f_T(x, \theta)] - \theta p_t(x))^2 dx \rightarrow 0$$

according to Lemma 2.3. By the Schwarz inequality, in analogy with (2.7) we have

$$\begin{aligned} I_2(T) &= \int_0^\varepsilon dt \int \left(\int_0^t P_{t-r}[|V_T^2(r) - V^2(r)|](x) dr \right)^2 dx \\ &= \int_0^\varepsilon dt \int_0^t ds \int_0^t dr \iint p_{2t-s-r}(x, y) |V_T(r) - V(r)|^2(x) \\ &\quad \times |V_T(s) + V(s)|^2(y) dx dy \\ &\leq (4\pi)^{-1/2} \int_0^\varepsilon dt \int_0^t dr \int |V_T(r, x) - V(r, x)|^2 dx \\ &\quad \times \int_0^t (2t - s - r)^{-1/2} ds \int (V_T^2(s, y) + V^2(s, y)) dy \\ &\leq (4\pi)^{-1/2} \int_0^\varepsilon dt \int_0^t (t - r)^{-1/2} dr \int |V_T(r, x) - V(r, x)|^2 dx \\ &\quad \times \int_0^t ds \int (\bar{V}_T^2(s, y) + V^2(s, y)) dy. \end{aligned}$$

For any $x \in \mathbb{R}$, using (2.13) we obtain

$$\frac{1}{\theta} \int |f_T(x - y, \theta)| dy = \frac{1}{\theta} \int |f_T(x - y, \theta)| dx = \frac{1}{\theta} \int |f_T(z, \theta)| dz \leq M(\theta, f)$$

for $T > T_0$, where $M(\theta, f) > 0$ is a constant depending only on θ and f . Then, by the Schwarz inequality, we have

$$\begin{aligned} & \int_0^1 \int \bar{V}_T^2(s, x) \, dx \, ds \\ &= \int_0^1 \, ds \int \left(\frac{1}{\theta} \int f_T(x - y, \theta) U(s, y; \theta) \, dy \right)^2 \, dx \\ &\leq \frac{1}{\theta^2} \int_0^1 \, ds \int \, dx \left(\int |f_T(x - y, \theta)| U^2(s, y; \theta) \, dy \right) \left(\int |f_T(x - y)| \, dy \right) \\ &\leq \frac{M(\theta, f)}{\theta} \int_0^1 \, ds \int U^2(s, y; \theta) \, dy \int |f_T(x - y, \theta)| \, dx \\ &\leq M^2(\theta, f) \int_0^1 \, ds \int U^2(s, y; \theta) \, dy \\ &< \infty. \end{aligned}$$

By Lemma 2.1, we have

$$\int_0^1 \, ds \int V^2(s, y; \theta) \, dy < \infty.$$

By combining these results, for $0 < \varepsilon < 1$ we find that

$$\begin{aligned} I_2(T) &\leq C \int_0^\varepsilon \, dt \int_0^t (t - r)^{-1/2} \, dr \int |V_T(r, x) - V(r, x)|^2 \, dx \\ &\leq C \int_0^\varepsilon \, dr \int_r^\varepsilon (t - r)^{-1/2} \, dt \int |V_T(r, x) - V(r, x)|^2 \, dx \\ &\leq C \int_0^\varepsilon 2(\varepsilon - r)^{1/2} \, dr \int |V_T(r, x) - V(r, x)|^2 \, dx \\ &\leq 2C\varepsilon^{1/2} \int_0^\varepsilon \, dr \int |V_T(r, x) - V(r, x)|^2 \, dx \\ &= 2C\varepsilon^{1/2} I(T), \end{aligned}$$

where $C > 0$ is a constant. Recalling (2.14), we have

$$I(T) \leq 2I_1(T) + 4C\varepsilon^{1/2} I(T).$$

By choosing an $\varepsilon > 0$ such that $4C\varepsilon^{1/2} < 1$, we see that

$$I(T) \leq 2(1 - 4C\varepsilon^{1/2})^{-1} I_1(T) \rightarrow 0$$

as $T \rightarrow \infty$. Now consider dividing $[0, 1]$ into intervals $[0, \varepsilon], [\varepsilon, 2\varepsilon], \dots$. By similar methods we can prove that

$$\int_\varepsilon^{2\varepsilon} \, dr \int (V_T(r, x) - V(r, x))^2 \, dx \rightarrow 0, \quad n = 0, 1, 2, \dots \tag{2.15}$$

In fact, for $t \in [\varepsilon, 2\varepsilon]$ we have

$$\begin{aligned} V_T(t, x; \theta) &= P_{t-\varepsilon}[V_T(\varepsilon, \cdot; \theta)](x) + \int_\varepsilon^t P_{t-r}[V_T^2(r, \cdot; \theta)] \, dr, \\ V(t, x; \theta) &= P_{t-\varepsilon}[V(\varepsilon, \cdot; \theta)](x) + \int_\varepsilon^t P_{t-r}[V^2(r, \cdot; \theta)] \, dr. \end{aligned}$$

In analogy with (2.14), define

$$I'(T) = \int_{\varepsilon}^{2\varepsilon} dt \int |V_T(t, x; \theta) - V(t, x; \theta)|^2 dx.$$

As $T \rightarrow \infty$,

$$\begin{aligned} I'_1(T) &:= \int_{\varepsilon}^{2\varepsilon} dt \int (P_{t-\varepsilon}[V_T(\varepsilon, \cdot; \theta)](x) - P_{t-\varepsilon}[V(\varepsilon, \cdot; \theta)](x))^2 dx \\ &\leq \int_{\varepsilon}^{2\varepsilon} dt \int (P_{\varepsilon}[V_T(t - \varepsilon, \cdot; \theta)](x) - V(t - \varepsilon, \cdot; \theta)(x))^2 dx \\ &= \int_{\varepsilon}^{2\varepsilon} dt \int (V_T(t - \varepsilon, \cdot; \theta)(x) - V(t - \varepsilon, \cdot; \theta)(x))^2 dx \\ &= I_1(T) \\ &\rightarrow 0. \end{aligned}$$

The remainder of the proof of (2.15) is analogous to our previous discussions on $I(T)$. Clearly there are no more than $\lfloor 1/\varepsilon \rfloor + 1$ such intervals of the form $[n\varepsilon, (n + 1)\varepsilon]$, where $\lfloor \cdot \rfloor$ is the least-integer function. Therefore we arrive at

$$\int_0^1 dr \int (V_T(r, x) - V(r, x))^2 dx \rightarrow 0.$$

Now let us see the analyticity of $V_T(t, x; \theta)$. For each $T > 0$, define $\{v_T^{(n)}, n \in \mathbb{N}\}$ recursively by

$$\begin{aligned} v_T^{(0)}(t, x; \theta) &= 0, \\ v_T^{(n+1)}(t, x; \theta) &= P_t[f_T(\cdot, \theta)](x) + \int_0^t P_{t-s}[v_T^{(n)}(s, \cdot; \theta)]^2(x) ds. \end{aligned} \tag{2.16}$$

By induction, for each real number $\theta \in (0, \theta_0)$, we obtain $0 \leq v_T^{(n)}(t, x; \theta) \leq \bar{V}_T(t, x; \theta)$ for all n, t , and x , and $v_T^{(n)}$ is increasing in n . By applying the monotone convergence theorem to (2.16), for each $\theta \in (0, \theta_0)$ we find that

$$v_T^{(n)}(t, x; \theta) \rightarrow V_T(t, x; \theta) \quad \text{as } n \rightarrow \infty.$$

Now allow θ to be a complex number. Again by induction, it is easy to see that, for $n \in \mathbb{N}$, $v_T^{(n)}(t, x; \theta)$ is analytic in $|\theta| < \theta_0$ and

$$\sup_n |v_T^{(n)}(t, x; \theta)| \leq V_T(t, x; |\theta|), \quad |\theta| < \theta_0.$$

Thus, $\{v_T^{(n)}(t, x; \theta), n \in \mathbb{N}\}$ is a normal family of analytic functions on the disc $|\theta| < \theta_0$, and $v_T^{(n)}(t, x; \theta) \rightarrow V_T(t, x; \theta)$ for θ in the real interval $(0, \theta_0)$. Therefore, by Conway (1978, p. 151, Theorem 2.1) and Vitali's theorem (see Conway (1978, p. 154)), the sequence

$$\{v_T^{(n)}(t, x; \theta), n \in \mathbb{N}\}$$

converges to an analytic function $V_T(t, x; \theta)$ on the disc $|\theta| < \theta_0$.

Corollary 2.2. *Suppose that $\Lambda(\theta)$ is defined by (1.4). As $T \rightarrow \infty$, we have*

$$\Lambda(\theta) = \theta + \int_0^1 ds \int_0^s \langle \lambda, [V(r, \cdot; \theta)]^2 \rangle dr, \tag{2.17}$$

and $\lim_{\theta \uparrow \theta_0} \Lambda'(\theta) = \infty$.

Proof. For $\theta \leq 0$, by (1.3) and (1.1) we have

$$\Lambda(T, \theta) = T^{-1/2} \log E \exp\{\langle N_T, \theta T^{-1/2} f \rangle\} = T^{-1/2} \int_0^T dr \int \tilde{v}(r, x; \theta T^{-1/2}) dx,$$

where $\tilde{v}(\cdot, \cdot; \theta T^{-1/2})$ is the solution to

$$\frac{\partial \tilde{v}(t)}{\partial t} = \Delta \tilde{v}(t) + \tilde{v}^2(t), \quad 0 < t \leq 1, \quad \tilde{v}(0) = e^{\theta T^{-1/2} f} - 1.$$

By changing variables according to $r = uT$ and $x = T^{1/2}y$, we obtain

$$\Lambda(T, \theta) = \int_0^1 \langle \lambda, V_T(u, \cdot; \theta) \rangle du = \langle \lambda, f_T(\cdot, \theta) \rangle + \int_0^1 ds \int_0^s \langle \lambda, [V_T(r, \cdot; \theta)]^2 \rangle dr \tag{2.18}$$

for $\theta \leq 0$, where f_T is as given in Lemma 2.2 and $V_T(\cdot, \cdot; \theta)$ is the solution to (2.11). Now allow θ to be a complex variable. For $|\theta| < \theta_0$, the analyticity of $f_T(x, \theta)$ and $V_T(t, x; \theta)$ implies that

$$\Gamma(T, \theta) := \langle \lambda, f_T(\cdot, \theta) \rangle + \int_0^1 ds \int_0^s \langle \lambda, [V_T(r, \cdot; \theta)]^2 \rangle dr$$

is an analytic function on the disc $|\theta| < \theta_0$. For a fixed f and T , denote the law of $\langle N_T, T^{-1/2} f \rangle$ by μ_T . Then $E \exp\{\langle N_T, \theta T^{-1/2} f \rangle\}$ is the Laplace transform of the probability law μ_T on $[0, \infty)$. By Widder (1941, p. 57, Theorem 5a), $E \exp\{\langle N_T, \theta T^{-1/2} f \rangle\}$ and, thus, $\Lambda(T, \theta)$ are analytic in the half-plane $\{\theta: \theta = \sigma + i\tau, \sigma < 0\}$. For each real number $\theta < 0$, we then have

$$\Lambda(T, \theta) = \Gamma(T, \theta).$$

Therefore, by the uniqueness of analytic extension, $\Lambda(T, \theta)$ can be uniquely extended to the real line-segment $[0, \theta_0)$, upon which it coincides with $\Gamma(T, \theta)$. Thus, (2.18) holds for the real number θ , $-\infty < \theta < \theta_0$. Let $T \rightarrow \infty$. Then, noting the two formulae below (2.13), we have $\langle \lambda, f_T(\cdot, \theta) \rangle \rightarrow \theta$, and by applying Lemma 2.2 we recover (2.17). By (2.5), for $\theta \in (0, \theta_0)$ we have

$$\begin{aligned} \Lambda(\theta) &= \theta + \theta^4 \int_0^1 ds \int_0^s dr \int w^2(\theta^2 r, \theta x; \delta_0) dx \\ &= \theta + \theta^{-1} \int_0^{\theta^2} ds \int_0^s dr \int w^2(r, y; \delta_0) dy, \end{aligned}$$

and, so,

$$\begin{aligned}\Lambda'(\theta) &= 1 - \theta^{-2} \int_0^{\theta^2} ds \int_0^s dr \int w^2(r, y; \delta_0) dy + 2 \int_0^{\theta^2} dr \int w^2(r, y; \delta_0) dy \\ &\geq 1 - \theta^{-2} \int_0^{\theta^2} ds \int_0^{\theta^2} dr \int w^2(r, y; \delta_0) dy + 2 \int_0^{\theta^2} dr \int w^2(r, y; \delta_0) dy \\ &= 1 + \int_0^{\theta^2} dr \int w^2(r, y; \delta_0) dy \\ &= 1 + \|V(\cdot, \cdot; \theta)\|_{L^2([0,1] \times \mathbb{R})}^2.\end{aligned}$$

Thus, the second assertion follows from Lemma 2.1, and the proof is complete.

Proof of Theorem 1.1. This is immediate from Corollary 2.2 and the Gärtner–Ellis theorem (see Dembo and Zeitouni (1998)).

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