

EXISTENCE THEORY FOR NONRESONANT SINGULAR BOUNDARY VALUE PROBLEMS

by DONAL O'REGAN

(Received 22nd December 1993)

We present some existence results for the “nonresonant” singular boundary value problem $\frac{1}{pq}(py')' + \mu y = f(t, y)$ a.e. on $[0, 1]$ with $\lim_{t \rightarrow 0^+} p(t)y'(t) = y(1) = 0$. Here μ is such that $\frac{1}{pq}(py')' + \mu u = 0$ a.e. on $[0, 1]$ with $\lim_{t \rightarrow 0^+} p(t)u'(t) = u(1) = 0$ has only the trivial solution.

1991 Mathematics subject classification: 34B15.

1. Introduction

This paper establishes existence results for the “nonresonant” singular boundary value problem

$$\begin{cases} \frac{1}{p(t)q(t)}(p(t)y'(t))' + \mu y(t) = f(t, y(t)) & \text{a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{cases} \quad (1.1)$$

where μ is such that

$$\begin{cases} \frac{1}{pq}(py')' + \mu y(t) = 0 & \text{a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{cases} \quad (1.2)$$

has only the trivial solution. Throughout the paper $p \in C[0, 1] \cap C^1(0, 1)$ together with $p > 0$ on $(0, 1)$; also q is measurable with $q > 0$ a.e. on $[0, 1]$ and $\int_0^1 p(x)q(x) dx < \infty$.

Remarks. (i). Throughout the condition $y(1) = 0$ could be replaced by the more general condition $ay(1) + b \lim_{t \rightarrow 1^-} p(t)y'(t) = 0, a > 0, b \geq 0$.

(ii). We do **not** assume $\int_0^1 \frac{ds}{p(s)} < \infty$.

In addition $f: [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ will be a *Carathéodory* function. By this we mean:

(i). $t \rightarrow f(t, y)$ is measurable for all $y \in \mathbf{R}$

(ii). $y \rightarrow f(t, y)$ is a continuous for a.e. $t \in [0, 1]$.

For notional purposes let w be a weight function. By $L_w^r[0, 1], r \geq 1$ we mean the space of functions u such that $\int_0^1 w(t)|u(t)|^r dt < \infty$. In particular $L_w^2[0, 1]$ denotes the space of functions u such that $\int_0^1 w(t)|u(t)|^2 dt < \infty$; also for $u, v \in L_w^2[0, 1]$ define $\langle u, v \rangle =$

$\int_0^1 w(t)u(t)\overline{v(t)} dt$. Let $AC[0, 1]$ be the space of functions which are absolutely continuous on $[0, 1]$.

This paper will be divided into three main sections. Section 2 discusses the linear problem i.e. (1.1) with $f \equiv 0$. In Section 3 fixed point methods (in particular a nonlinear alternative of Leray–Schauder type) is used to obtain an existence principle. The final section establishes some existence results for (1.1); these results extend and complement the theory in [4, 6, 21].

Finally we remark here that problems of type (1.1) occur in many applications in the physical sciences, for example in radially symmetric nonlinear diffusion [20, 22] in the n -dimensional sphere we have $p(t) = t^{n-1}$; these problems involve a homogeneous Neumann condition at zero i.e. $\lim_{t \rightarrow 0^+} t^{n-1} y'(t) = 0$. Another important example is the Poisson–Boltzmann equation

$$\begin{cases} y'' + \frac{\alpha}{t}y' = f(t, y), 0 < t < 1 \\ y'(0^+) = y(1) = 0, \alpha \geq 1 \end{cases} \tag{1.3}$$

which occurs in the theory of thermal explosions [3] and in the theory of electrohydrodynamics [11]. The results related to (1.3) in the literature [4] usually consider the situation when $\inf \frac{\partial f}{\partial y}, \sup \frac{\partial f}{\partial y}$ are bounded and satisfy a “nonresonant” condition; here the infimum and supremum are taken over $\{(t, y) : 0 \leq t \leq 1, -\infty < y < \infty\}$. In this paper we improve the above existence result; in fact in our theory the existence of $\frac{\partial f}{\partial y}$ is not assumed.

2. Linear problem

Theorem 2.1. *Suppose*

$$p \in C[0, 1] \cap C^1(0, 1) \text{ with } p > 0 \text{ on } (0, 1) \tag{2.1}$$

$$q \in L^1_p[0, 1] \text{ with } q > 0 \text{ a.e. on } [0, 1] \tag{2.2}$$

and

$$\int_0^1 \frac{1}{p(s)} \left(\int_0^s p(x)q(x) dx \right)^{1/\alpha} ds < \infty \text{ for some constant } \alpha > 1 \tag{2.3}$$

are satisfied.

(i) *Then*

$$\begin{cases} \frac{1}{p}(py')' + \mu qy = 0 \text{ a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(0) = a_0 \neq 0 \end{cases} \tag{2.4}$$

has a solution $y_1 \in C[0, 1] \cap C^1(0, 1)$ with $py'_1 \in AC[0, 1]$. (By a solution to (2.4) we mean a function $y \in C[0, 1] \cap C^1(0, 1)$, $py' \in AC[0, 1]$ which satisfies the differential equation a.e. on $[0, 1]$ and the stated conditions.)

(ii) Then

$$\begin{cases} \frac{1}{p}(py)' + \mu qy = 0 & \text{a.e. on } [0, 1] \\ y(1) = 0 \\ \lim_{t \rightarrow 1^-} p(t)y'(t) = 1 \end{cases} \tag{2.5}$$

has a solution $y_2 \in L^{\alpha}_{pq}[0, 1]$ with $y_2 \in C(0, 1] \cap C^1(0, 1)$ and $py'_2 \in AC[0, 1]$.

Proof. (i). Let $C[0, 1]$ denote the Banach space of continuous functions on $[0, 1]$ endowed with the norm

$$\|u\|_K = \sup_{t \in [0, 1]} |e^{-KR(t)}u(t)| \quad \text{where} \quad R(t) = \int_0^t p(x)q(x) dx$$

and

$$K = \frac{1}{\beta} \left(|\mu| \int_0^1 \frac{1}{p(s)} \left(\int_0^s p(x)q(x) dx \right)^{1/\alpha} ds \right)^{\beta}.$$

Remark. Here $\beta = \frac{\alpha}{\alpha-1}$ i.e. β and α are conjugate exponents.

Solving (2.4) is equivalent to finding $y \in C[0, 1]$ which satisfies

$$y(t) = a_0 - \mu \int_0^t \frac{1}{p(s)} \int_0^s p(x)q(x)y(x) dx ds.$$

Define the operator $N : C[0, 1] \rightarrow C[0, 1]$ by

$$Ny(t) = a_0 - \mu \int_0^t \frac{1}{p(s)} \int_0^s p(x)q(x)y(x) dx ds.$$

Now N is a contraction since

$$\begin{aligned} \|Nu - Nv\|_K &\leq |\mu| \|u - v\|_K \sup_{t \in [0, 1]} e^{-KR(t)} \int_0^t \frac{1}{p(s)} \int_0^s p(x)q(x)e^{KR(x)} dx ds \\ &\leq |\mu| \|u - v\|_K \sup_{t \in [0, 1]} e^{-KR(t)} \int_0^t \frac{1}{p(s)} \left(\int_0^s pq dx \right)^{1/\alpha} \left(\int_0^s pqe^{\beta KR(x)} dx \right)^{1/\beta} ds \\ &\leq |\mu| \|u - v\|_K \sup_{t \in [0, 1]} e^{-KR(t)} \int_0^t \frac{1}{p(s)} \left(\int_0^s p(x)q(x) dx \right)^{1/\alpha} \left(\frac{e^{\beta KR(s)}}{\beta K} - \frac{1}{\beta K} \right)^{1/\beta} ds \\ &\leq \frac{|\mu|}{(\beta K)^{1/\beta}} \|u - v\|_K \sup_{t \in [0, 1]} e^{-KR(t)} (e^{\beta KR(t)} - 1)^{1/\beta} \int_0^t \left(\int_0^s p(x)q(x) dx \right)^{1/\alpha} ds \\ &\leq (1 - e^{-\beta KR(1)})^{1/\beta} \|u - v\|_K \end{aligned}$$

using Hölder's integral inequality. The Banach contraction principle now establishes the result.

(ii). Let $L^{\alpha}_{pq}[0, 1]$ denote the Banach space of functions u , with $\int_0^1 pq|u|^{\alpha} dt < \infty$, endowed with the norm

$$\|u\|_K = \left(\int_0^1 p(t)q(t)e^{-KQ(t)} |u(t)|^{\alpha} dt \right)^{1/\alpha} \quad \text{where } Q(t) = \int_t^1 p(x)q(x) dx$$

and

$$K = \frac{\alpha}{\beta} \left(|\mu|^{\alpha} \int_0^1 p(t)q(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{\alpha} dt \right)^{\beta/\alpha} \quad \text{where } \beta = \frac{\alpha}{\alpha - 1}.$$

Remarks. (i). Notice for example that $\int_{1/2}^1 \frac{ds}{p(s)} < \infty$ since

$$\int_{1/2}^1 \frac{ds}{p(s)} = \int_{1/2}^1 \frac{(\int_0^s p(x)q(x) dx)^{1/\alpha}}{p(s)(\int_0^s p(x)q(x) dx)^{1/\alpha}} ds \leq \frac{1}{(\int_0^{1/2} p(x)q(x) dx)^{1/\alpha}} \int_{1/2}^1 \frac{1}{p(s)} \left(\int_0^s p(x)q(x) dx \right)^{1/\alpha} ds.$$

(ii). Notice (2.3) implies

$$\int_0^1 p(t)q(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{\alpha} dt < \infty. \tag{2.6}$$

To see this let

$$g(t) = \left(\int_t^1 \frac{ds}{p(s)} \right)^{\alpha - 1}$$

and fix $\varepsilon, 0 < \varepsilon < 1$. Interchange the order of integration and use Hölder's inequality to obtain

$$\begin{aligned} \int_{\varepsilon}^1 p(t)q(t)g(t) \int_t^1 \frac{ds}{p(s)} dt &= \int_{\varepsilon}^1 \frac{1}{p(s)} \int_{\varepsilon}^s p(t)q(t)g(t) dt \\ &\leq \left(\int_{\varepsilon}^1 p(t)q(t)g^{\beta}(t) dt \right)^{1/\beta} \int_{\varepsilon}^1 \frac{1}{p(s)} \left(\int_{\varepsilon}^s p(t)q(t) dt \right)^{1/\alpha} ds. \end{aligned}$$

Consequently

$$\int_{\varepsilon}^1 p(t)q(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{\alpha} dt \leq \left(\int_{\varepsilon}^1 p(t)q(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{\alpha} dt \right)^{1/\beta} \int_{\varepsilon}^1 \frac{1}{p(s)} \left(\int_{\varepsilon}^s p(t)q(t) dt \right)^{1/\alpha} ds.$$

We will show that

$$y(t) = -\int_t^1 \frac{ds}{p(s)} - \mu \int_t^1 \frac{1}{p(s)} \int_s^1 p(x)q(x)y(x) dx ds \tag{2.7}$$

has a solution $y_2 \in L_{pq}^\alpha [0, 1]$. Also we will show $y_2 \in C(0, 1] \cap C^1(0, 1)$ and $py_2' \in AC[0, 1]$ and consequently y_2 will be a solution of (2.5).

Define the operator: $L_{pq}^\alpha [0, 1] \rightarrow L_{pq}^\alpha [0, 1]$ by

$$My(t) = -\int_t^1 \frac{ds}{p(s)} - \mu \int_t^1 \frac{1}{p(s)} \int_s^1 p(x)q(x)y(x) dx ds.$$

Remark. M is well defined because of (2.6) and

$$\begin{aligned} \int_0^1 pq \left(\int_t^1 \frac{1}{p} \int_s^1 pq|y| dx ds \right)^\alpha dt &\leq \left(\int_0^1 pq|y| dx \right)^\alpha \int_0^1 pq \left(\int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \\ &\leq \left(\int_0^1 pq|y|^\alpha dx \right) \left(\int_0^1 pq dx \right)^{\alpha/\beta} \int_0^1 pq \left(\int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \end{aligned}$$

for any $y \in L_{pq}^\alpha [0, 1]$.

Now M is a contraction since

$$\begin{aligned} \|Mu - Mv\|_K^\alpha &\leq |\mu|^\alpha \int_0^1 pq e^{-KQ(t)} \left(\int_t^1 \frac{1}{p(s)} \int_s^1 p(x)q(x)|u(x) - v(x)| dx ds \right)^\alpha dt \\ &\leq |\mu|^\alpha \int_0^1 pq e^{-KQ(t)} \left(\int_t^1 pq e^{-KQ(x)/\alpha} e^{KQ(x)/\alpha} |u(x) - v(x)| dx \int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \\ &\leq |\mu|^\alpha \|u - v\|_K^\alpha \int_0^1 pq e^{-KQ(t)} \left(\int_t^1 pq e^{\beta KQ(x)/\alpha} dx \right)^{\alpha/\beta} \left(\int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \\ &\leq |\mu|^\alpha \|u - v\|_K^\alpha \int_0^1 pq e^{-KQ(t)} \left(\frac{\alpha}{\beta K} e^{\beta KQ(t)/\alpha} - \frac{\alpha}{\beta K} \right)^{\alpha/\beta} \left(\int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \\ &\leq \left(\frac{\alpha}{\beta K} \right)^{\alpha/\beta} |\mu|^\alpha \|u - v\|_K^\alpha \int_0^1 pq \left(1 - e^{-\beta KQ(t)/\alpha} \right)^{\alpha/\beta} \left(\int_t^1 \frac{ds}{p(s)} \right)^\alpha dt \\ &\leq \left(1 - e^{-\beta KQ(0)/\alpha} \right)^{\alpha/\beta} \|u - v\|_K^\alpha. \end{aligned}$$

The Banach contraction principle now establishes that (2.7) has a solution $y_2 \in L^{\alpha}_{pq}[0, 1]$. Also

$$p(t)y'_2(t) = 1 + \mu \int_0^1 p(x)q(x)y_2(x) dx$$

so $py'_2 \in AC[0, 1]$ since $y_2 \in L^{\alpha}_{pq}[0, 1]$ implies $pqy_2 \in L^1[0, 1]$. Thus y_2 is a solution of (2.5). □

Consider now

$$\frac{1}{pq}(py')' + \mu y = h(t) \quad \text{a.e. on } [0, 1] \tag{2.8}$$

where (2.1), (2.2), (2.3) and

$$h \in L^{\beta}_{pq}[0, 1]; \text{ here } \beta = \frac{\alpha}{\alpha - 1} \tag{2.9}$$

hold.

Theorem 2.2. *Suppose (2.1), (2.2), (2.3) and (2.9) are satisfied. In addition μ is such that (1.2) has only the trivial solution. Then*

$$\begin{cases} \frac{1}{pq}(py')' + \mu y = h(t) & \text{a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{cases} \tag{2.10}$$

has exactly one solution y (note $y \in L^{\alpha}_{pq}[0, 1]$ with $y \in C(0, 1) \cap C^1(0, 1)$ and $py' \in AC[0, 1]$) given by

$$y(t) = \int_0^1 G(t, s)q(s)h(s) ds \tag{2.11}$$

where $G(t, s)$ is the Green's function i.e.

$$G(t, s) = \begin{cases} \frac{y_1(s)y_2(t)}{W(s)} = c_0 p(s)y_1(s)y_2(t), & 0 < s \leq t \\ \frac{y_1(t)y_2(s)}{W(s)} = c_0 p(s)y_2(s)y_1(t), & 0 \leq s < t. \end{cases}$$

Here y_1 and y_2 are as described in Theorem 2.1 and $W(s)$ is the Wronskian of y_1 and y_2 at s and notice $p(s)W(s) = (1/c_0) \neq 0$ for $s \in [0, 1]$.

Proof. This follows the standard construction of the Green's function; see [22, 24] for example. We will just justify that $p(s)W(s) \neq 0$ for $s \in [0, 1]$. To see this all one needs to show is that $y_1(1) \neq 0$. If $y_1(1) = 0$ then y_1 satisfies (1.2) and consequently $y_1 \equiv 0$. This contradicts the fact that $y(0) = a_0 \neq 0$. \square

Remark. Notice y in (2.11) is in $L^{\alpha}_{pq}[0, 1]$ since

$$\int_0^1 p(t)q(t) \left(\int_t^1 p(s)q|y_2(s)y_1(t)h(s)| ds \right)^{\alpha} dt \leq \int_0^1 pq|y_1|^{\alpha} \left(\int_t^1 pq|y_2|^{\alpha} ds \right) \left(\int_t^1 pq|h|^{\beta} ds \right)^{\alpha/\beta} ds$$

and so

$$\int_0^1 p(t)q(t) \left(\int_0^t p(s)q(s)|y_1(s)y_2(t)h(s)| ds \right)^{\alpha} dt < \infty.$$

3. Existence principle

We use a nonlinear alternative of Leray–Schauder type [9] to establish our existence principle. By a map being *compact* we mean it is continuous with relatively compact range. A map is *completely continuous* if it is continuous and the image of every bounded set in the domain is contained in a compact set in the range.

Theorem 3.1. *Assume U is a relatively open subset of a convex set K in a Banach space E . Let $N : \bar{U} \rightarrow K$ be a compact map with $0 \in U$. Then either*

- (i) N has a fixed point in \bar{U} ; or
- (ii) there is a $u \in \partial U$ and a $\lambda \in (0, 1)$ such that $u = \lambda Nu$.

Next we gather some well known results [12] from the theory of nonlinear integral equations.

Theorem 3.2. *Let $\alpha > 1$ be a constant and $f : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ be a Carathéodory function. Define the operator*

$$Fy(t) = f(t, y(t))$$

and suppose $F : L^{\alpha}_{pq}[0, 1] \rightarrow L^{\beta}_{pq}[0, 1]$; here $\beta = \frac{\alpha}{\alpha-1}$. Then F is continuous and bounded.

Theorem 3.3. *Consider the linear integral operator*

$$Ay(t) = \int_0^1 p(s)q(s)k(t, s)y(s) ds$$

with

$$\int_0^1 p(t)q(t) \int_0^1 p(s)q(s)|k(t,s)|^\alpha ds dt < \infty \text{ for some } \alpha > 1. \tag{3.1}$$

Then $A : L_{pq}^\beta[0, 1] \rightarrow L_{pq}^\alpha[0, 1]$, $\beta = \frac{\alpha}{\alpha-1}$ is completely continuous.

We next prove an existence principle for (1.1).

Theorem 3.4. *Let $f : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ be a Carathéodory function and suppose (2.1), (2.2) and (2.3) are satisfied. Also suppose*

$$f(t, y(t)) \in L_{pq}^\beta[0, 1] \text{ whenever } y \in L_{pq}^\alpha[0, 1]; \text{ here } \beta = \frac{\alpha}{\alpha-1}. \tag{3.2}$$

In addition μ is such that (1.2) has only the **trivial** solution. Now suppose there is a constant M_0 , independent of λ , with

$$\|y\| = \left(\int_0^1 p(t)q(t)|y(t)|^\alpha dt \right)^{1/\alpha} \leq M_0$$

for any solution y (here $y \in L_{pq}^\alpha[0, 1]$ with $y \in C(0, 1] \cap C^1(0, 1)$ and $py' \in AC[0, 1]$) to

$$\begin{cases} \frac{1}{pq}(py')' + \mu y = \lambda f(t, y) & \text{a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{cases} \tag{3.3}_\lambda$$

for each $\lambda \in (0, 1)$. Then (1.1) has at least one solution.

Proof. Solving (3.3) $_\lambda$ is equivalent to finding $y \in L_{pq}^\alpha[0, 1]$ which satisfies

$$y(t) = \lambda \int_0^1 p(s)q(s)k(t,s)f(s,y(s)) ds \tag{3.4}$$

where

$$k(t,s) = \begin{cases} c_0 y_1(s)y_2(t), & 0 < s \leq t \\ c_0 y_2(s)y_1(t), & t \leq s < 1, \end{cases}$$

and y_1, y_2, c_0 are described in Theorem 2.2. Define the operator $N : L_{pq}^\alpha[0, 1] \rightarrow L_{pq}^\alpha[0, 1]$ by

$$Ny(t) = \int_0^1 p(s)q(s)k(t,s)f(s,y(s)) ds.$$

Remark. N is well defined since

$$\int_0^1 pq \left(\int_t^1 pq |y_1(t)y_2(s)f(s,y)| ds \right)^\alpha dt \leq \int_0^1 pq |y_1|^\alpha \left(\int_t^1 pq |y_2|^\alpha ds \right) \left(\int_t^1 pq |f(s,y)|^\beta ds \right)^{\alpha/\beta} dt$$

and so

$$\int_0^1 p(t)q(t) \left(\int_0^t p(s)q(s) |y_1(s)y_2(t)f(s,y(s))| ds \right)^\alpha dt < \infty.$$

Next define $F : L_{pq}^\alpha[0, 1] \rightarrow L_{pq}^\beta[0, 1]$ by

$$Fy(t) = f(t, y(t))$$

and $A : L_{pq}^\beta[0, 1] \rightarrow L_{pq}^\alpha[0, 1]$ by

$$Ay(t) = \int_0^1 p(s)q(s)k(t,s)y(s) ds.$$

Notice (3.2) and Theorem 3.2 implies F is bounded and continuous. A is completely continuous by Theorem 3.3.

Remark. Notice $\int_0^1 p(t)q(t) \int_0^1 p(s)q(s) |k(t,s)|^\alpha ds dt < \infty$ since

$$\int_0^1 p(t)q(t) \left(\int_0^t p(s)q(s) |y_1(s)y_2(t)|^\alpha ds dt \leq \int_0^1 p(t)q(t) |y_2(t)|^\alpha \int_0^1 p(s)q(s) |y_1(s)|^\alpha ds dt < \infty$$

and so

$$\int_0^1 p(t)q(t) \left(\int_t^1 p(s)q(s) |y_2(s)y_1(t)|^\alpha ds dt < \infty.$$

Consequently $N = AF : L_{pq}^\alpha[0, 1] \rightarrow L_{pq}^\alpha[0, 1]$ is completely continuous. Set

$$U = \{u \in L_{pq}^\alpha[0, 1] : \|u\| < M_0 + 1\}, K = E = L_{pq}^\alpha[0, 1].$$

Then Theorem 3.1 implies that N has a fixed point i.e. (1.1) has a solution $y \in L_{pq}^\alpha[0, 1]$. The fact that $y \in C(0, 1) \cap C^1(0, 1)$ with $py' \in AC[0, 1]$ follows from (3.4) with $\lambda = 1$. \square

4. Existence theory

Theorem 4.1. Let $f : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ be a Carathéodory function and suppose (2.1), (2.2)

and (2.3) are satisfied. In addition μ is such that (1.2) has only the **trivial** solution. Let $\beta = \frac{\alpha}{\alpha-1}$. Now assume

$$\left\{ \begin{aligned} &|f(t, u)| \leq \phi_1(t) + \phi_2(t)\psi(|u|) \text{ a.e. on } [0, 1] \text{ where } \phi_1^\beta, \phi_2 \in L_{pq}^1[0, 1] \\ &\text{and } \psi : [0, \infty) \rightarrow [0, \infty) \text{ is a continuous function} \end{aligned} \right. \tag{4.1}$$

$$\left\{ \begin{aligned} &\text{there exists } Q_0 \geq 0 \text{ and a continuous function } \theta : [0, \infty) \rightarrow [0, \infty) \text{ with} \\ &\int_0^1 p(s)q(s)\phi_2^\beta(s)\psi^\beta(|y(s)|) ds \leq Q_0\theta(\|y\|) \text{ for any } y \in L_{pq}^\alpha[0, 1]; \\ &\text{here } \|y\| = (\int_0^1 p(t)q(t)|y(t)|^\alpha dt)^{1/\alpha} \end{aligned} \right. \tag{4.2}$$

and

$$\left\{ \begin{aligned} &A_0 \equiv 2^{q_0} c_0^\alpha Q_0^{\alpha/\beta} \|y_2\|^\alpha \|y_1\|^\alpha \limsup_{x \rightarrow \infty} \frac{(\theta(x))^{\alpha/\beta}}{x^\alpha} < 1 \text{ where } y_1, y_2, c_0 \\ &\text{are as described in Theorem 2.2 and } q_0 = \frac{2\alpha^2\beta - \beta^2 - \alpha^2 + \alpha\beta}{\alpha\beta} \end{aligned} \right. \tag{4.3}$$

are satisfied. Then (1.1) has at least one solution.

Remarks. (i). Notice (3.2) is automatically satisfied since (4.2) holds and also since $\phi_1^\beta \in L_{pq}^1[0, 1]$.

(ii). If $\psi(|u|) = |u|^\gamma$, $0 \leq \gamma < \min\{\frac{\alpha}{\beta}, 1\}$ and $\phi_2^{\beta\alpha/(\alpha-\beta\gamma)} \in L_{pq}^1[0, 1]$ then (4.2) and (4.3) are satisfied since

$$\int_0^1 p(s)q(s)\phi_2^\beta(s)|y(s)|^{\beta\gamma} ds \leq \|y\|^{\beta\gamma} \left(\int_0^1 pq\phi_2^{\beta\alpha/(\alpha-\beta\gamma)} ds \right)^{(\alpha-\beta\gamma)/\alpha} \text{ for any } y \in L_{pq}^\alpha[0, 1]$$

and so with $\theta(x) = x^{\beta\gamma}$ we have

$$\limsup_{x \rightarrow \infty} \frac{(\theta(x))^{\alpha/\beta}}{x^\alpha} = \limsup_{x \rightarrow \infty} x^{\alpha(\gamma-1)} = 0.$$

Proof. Let y be a solution to (3.3) $_\lambda$ for $0 < \lambda < 1$. Then

$$y(t) = \lambda c_0 y_2(t) \int_0^t p(s)q(s)y_1(s)f(s, y(s)) ds + \lambda c_0 y_1(t) \int_t^1 p(s)q(s)f(s, y(s)) ds$$

where y_1, y_2, c_0 are as described in Theorem 2.2. Recall $(a_0 + b_0)^{r_0} \leq 2^{r_0-1}(a_0^{r_0} + b_0^{r_0})$, $a_0 \geq 0, b_0 \geq 0, r_0 \geq 1$ so

$$\begin{aligned} \|y\|^\alpha &\leq 2^{\alpha-1} c_0^\alpha \int_0^1 p(t)q(t)|y_2(t)|^\alpha \left(\int_0^t p(s)q(s)|y_1(s)||f(s, y(s))| ds \right)^\alpha dt \\ &\quad + 2^{\alpha-1} c_0^\alpha \int_0^1 p(t)q(t)|y_1(t)|^\alpha \left(\int_t^1 p(s)q(s)|y_2(s)||f(s, y(s))| ds \right)^\alpha dt. \end{aligned}$$

This together with Hölder’s inequality implies

$$\|y\|^\alpha \leq 2^\alpha c_0^\alpha \|y_2\|^\alpha \|y_1\|^\alpha \left(\int_0^1 p(s)q(s)|f(s, y(s))|^\beta ds \right)^{\alpha/\beta}. \tag{4.4}$$

In addition

$$\begin{aligned} \int_0^1 p(s)q(s)|f(s, y(s))|^\beta ds &\leq 2^{\beta-1} \int_0^1 p(s)q(s)\phi_1^\beta(s) ds + 2^{\beta-1} \int_0^1 p(s)q(s)\phi_2^\beta(s)\psi^\beta(|y(s)|) ds \\ &\leq 2^{\beta-1} \int_0^1 p(s)q(s)\phi_1^\beta(s) ds + 2^{\beta-1} Q_0\theta(\|y\|). \end{aligned}$$

This inequality together with $(a_0 + b_0)^{1/r_0} \leq 2^{(r_0-1)/r_0}(a_0^{1/r_0} + b_0^{1/r_0})$, $a_0 \geq 0, b_0 \geq 0, r_0 \geq 1$ or $(a_0 + b_0)^{s_0} \leq 2^{s_0-1}(a_0^{s_0} + b_0^{s_0})$, $s_0 \geq 1$ and (4.4) implies

$$\|y\|^\alpha \leq 2^\alpha c_0^\alpha \|y_2\|^\alpha \|y_1\|^\alpha 2^{\alpha(\alpha-\beta)/\alpha} \left(2^{\alpha(\beta-1)/\beta} \left(\int_0^1 pq\phi_1^\beta ds \right)^{\alpha/\beta} + 2^{\alpha(\beta-1)/\beta} Q_0^{\alpha/\beta} (\theta(\|y\|))^{\alpha/\beta} \right). \tag{4.5}$$

Consequently

$$1 \leq 2^{2\alpha} c_0^\alpha \|y_2\|^\alpha \|y_1\|^\alpha \left(\frac{(\int_0^1 pq\phi_1^\beta ds)^{\alpha/\beta}}{\|y\|^\alpha} + \frac{Q_0^{\alpha/\beta} (\theta(\|y\|))^{\alpha/\beta}}{\|y\|^\alpha} \right). \tag{4.6}$$

Thus there exists a constant M_0 , independent of λ , with $\|y\| \leq M_0$ for any solution y satisfying (3.3) $_\lambda$ i.e. $y = \lambda Ny$ where N is as described in Theorem 3.4. If this was not true then there exists $u_n = \lambda_n Nu_n$ with $\|u_n\| \rightarrow \infty$ as $n \rightarrow \infty$ and since $\limsup(s_n + t_n) \leq \limsup s_n + \limsup t_n$ for any sequences $s_n \geq 0, t_n \geq 0$ we have from (4.6) that $1 \leq A_0$, a contradiction (see (4.3)). Thus there exists a constant M_0 , independent of λ , with $\|y\| \leq M_0$ and the result now follows from Theorem 3.4. \square

The next two existence results extend in a “particular direction” Theorem 4.1 if certain criteria are fulfilled. To discuss the first result we begin by gathering together some facts on the singular eigenvalue problem

$$\begin{cases} Lu = \lambda u \text{ a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)u'(t) = 0 \\ u(1) = 0 \end{cases} \tag{4.7}$$

where $Lu = -\frac{1}{pq}(pu)'$. Assume (2.1), (2.2) and

$$\int_0^1 \frac{1}{p(s)} \left(\int_0^s p(x)q(x) dx \right)^{1/2} ds < \infty \tag{4.8}$$

hold.

- Remarks.** (i). In this case $\alpha = 2$ in (2.3).
 (ii). Here $t = 0$ is a singular point in the limit circle case [18, 19, 24].

Let

$$D(L) = \left\{ \omega \in C[0, 1] : w, pw' \in AC[0, 1] \text{ with } \frac{1}{pq} (pw')' \in L^2_{pq}[0, 1] \right. \\ \left. \text{and } \lim_{t \rightarrow 0^+} p(t)w'(t) = w(1) = 0 \right\}.$$

In [18, 19] it was shown that $L^{-1} : L^2_{pq}[0, 1] \rightarrow D(L)$ and L^{-1} is completely continuous with $\langle L^{-1}u, v \rangle = \langle u, L^{-1}v \rangle$ for $u, v \in L^2_{pq}[0, 1]$. Consequently the spectral theorem for compact self adjoint operators [24] implies that L has a countably infinite number of real eigenvalues λ_i with corresponding eigenfunctions $\psi_i \in D(L)$. The eigenfunctions ψ_i may be chosen so that they form an orthonormal set and we may arrange the eigenvalues so that

$$\lambda_0 < \lambda_1 < \lambda_2 < \dots$$

The following Rayleigh–Ritz minimization theorem [18, 19] also holds.

Theorem 4.2. *Suppose (2.1), (2.2) and (4.8) hold. Then*

$$\lambda_0 \int_0^1 p(t)q(t)y^2(t) dt \leq \int_0^1 p(t)[y'(t)]^2 dt$$

for all functions $y \in D(L)$.

We can improve the result in Theorem 4.1 if (4.8) holds and if $\mu < \lambda_0$; here λ_0 is the first eigenvalue of (4.7). In particular consider

$$\begin{cases} \frac{1}{pq}(py')' = f(t, y) \text{ a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0. \end{cases} \tag{4.9}$$

Theorem 4.3. *Let $f : [0, 1] \times \mathbb{R} \rightarrow \mathbb{R}$ be a Carathéodory function and suppose (2.1), (2.2) and (4.8) are satisfied. Also assume*

$$f(t, y(t)) \in L^2_{pq}[0, 1] \text{ whenever } y \in L^2_{pq}[0, 1]. \tag{4.10}$$

In addition suppose $f(t, u) = g(t, u) + h(t, u)$ with $g, h: [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ Carathéodory functions and

$$\left\{ \begin{array}{l} |uh(t, u)| \leq \phi_1(t)|u| + \phi_2(t)\rho(|u|) \text{ a.e. on } [0, 1] \text{ where } \rho: [0, \infty) \rightarrow [0, \infty) \\ \text{is a nondecreasing continuous function} \end{array} \right. \tag{4.11}$$

$$ug(t, u) \geq -\mu_0 u^2 \text{ for a.e. } t \in [0, 1] \text{ and } u \in \mathbf{R}; \text{ here } \mu_0 < \lambda_0 \tag{4.12}$$

$$\int_0^1 p(t)q(t)\phi_1(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} dt < \infty \text{ and } \int_0^1 p(t)q(t)\phi_2(t)\rho \left(\left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} \right) dt < \infty \tag{4.13}$$

$$\left\{ \begin{array}{l} \text{there exist constants } Q_1 \text{ (independent of } a_0 \text{ and } b_0) \text{ and } Q_2 \text{ such that for any} \\ a_0 \geq 0, b_0 \geq 0 \text{ we have } \rho(a_0 b_0) \leq Q_1 \rho(a_0)\rho(b_0) + Q_2 \rho(b_0) \end{array} \right. \tag{4.14}$$

and

$$\left\{ \begin{array}{l} A_1 \equiv Q_1 \left(\int_0^1 p(t)q(t)\phi_2(t)\rho \left(\left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} \right) dt \right) \limsup_{x \rightarrow \infty} \frac{\rho(x)}{x^2} < \eta_0 \\ \text{with } \eta_0 = 1 \text{ if } \mu_0 < 0 \text{ whereas } \eta_0 = 1 - \frac{\mu_0}{\lambda_0} \text{ if } 0 \leq \mu_0 < \lambda_0 \end{array} \right. \tag{4.15}$$

are satisfied. Then (4.9) has at least one solution.

Remark. If $\rho(|u|) = |u|^{\gamma+1}, 0 \leq \gamma < 1$ and $\int_0^1 p(t)q(t)\phi_2(t) \left(\int_t^1 \frac{ds}{p(s)} \right)^{(\gamma+1)/2} dt < \infty$ then (4.13), (4.14) and (4.15) are satisfied since if $Q_1 = 1, Q_2 = 0$ we have $\rho(a_0 b_0) = |a_0 b_0|^{\gamma+1} = |a_0|^{\gamma+1} |b_0|^{\gamma+1}$ and also

$$\limsup_{x \rightarrow \infty} \frac{\rho(x)}{x^2} = \limsup_{x \rightarrow \infty} x^{\gamma-1} = 0.$$

Proof. Let y be a solution to

$$\left\{ \begin{array}{l} \frac{1}{pq}(py)' = \lambda f(t, y) \text{ a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{array} \right. \tag{4.16}_\lambda$$

for $0 < \lambda < 1$. Multiply the differential equation by $-y$ and integrate from 0 to 1 to obtain

$$\begin{aligned} \|y'\|_0^2 &\leq -\lambda \int_0^1 pqyg(t, y) dt + \int_0^1 pq|yh(t, y)| dt \\ &\leq \lambda\mu_0 \|y\|^2 + \int_0^1 pq[\phi_1(t)|y(t)| + \phi_1(t)\rho(|y(t)|)] dt \end{aligned}$$

where for notational purposes $\|u\|^2 = \int_0^1 pq|u|^2 dt$ and $\|u\|_0^2 = \int_0^1 p|u|^2 dt$. Apply Theorem 4.2 if $0 \leq \mu_0 < \lambda_0$ to obtain

$$\eta_0 \|y'\|_0^2 \leq \int_0^1 pq\phi_1|y| dt + \int_0^1 pq\phi_2\rho(|y(t)|) dt \tag{4.17}$$

where η_0 is as described in (4.15). Also for $t \in (0, 1)$ we have from Hölder's inequality that

$$|y(t)| \leq \|y'\|_0 \left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} \tag{4.18}$$

and this together with (4.17) and the fact that ρ is nondecreasing yields

$$\eta_0 \|y'\|_0^2 \leq N_1 \|y'\|_0 + \int_0^1 p(t)q(t)\phi_2(t)\rho \left(\|y'\|_0 \left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} \right) dt$$

where $N_1 = \int_0^1 pq\phi_1 \left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} dt$. Using (4.14) we obtain

$$\eta_0 \|y'\|_0^2 \leq N_1 \|y'\|_0 + Q_1 N_2 \rho(\|y'\|_0) + Q_2 N_2$$

where $N_2 = \int_0^1 pq\phi_2\rho \left(\left(\int_t^1 \frac{ds}{p(s)} \right)^{1/2} \right) dt$. Consequently

$$\eta_0 \leq \frac{N_1 \|y'\|_0 + Q_2 N_2}{\|y'\|_0^2} + \frac{Q_1 N_2 \rho(\|y'\|_0)}{\|y'\|_0^2}.$$

Thus (as in Theorem 4.1) exists a constant M_1 , independent of λ , with $\|y'\|_0 \leq M_1$ for any solution y to (4.16)₂. This together with Theorem 4.2 yields

$$\int_0^1 pq|y|^2 dt \leq \frac{1}{\lambda_0} M_1^2$$

so the result follows from Theorem 3.4 (with $\mu=0$ and $\alpha=2$). □

Finally we examine the boundary value problem (4.9) where in the case $pqf : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ is an L^1 -Carathéodory function. By this we mean:

- (i) $pqf : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ is a Carathéodory function, and

(ii) for any $r > 0$ there exists $h_r \in L^1[0, 1]$ with $|p(t)q(t)f(t, u)| \leq h_r(t)$ for a.e. $t \in [0, 1]$ and for all $|u| \leq r$.

For the remainder of the paper assume (2.1), (2.2) and

$$\int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x) dx ds < \infty \tag{4.19}$$

and

$$\int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)h_r(x) dx ds < \infty \text{ for any } r > 0; \text{ here } h_r \text{ is as described above} \tag{4.20}$$

hold. In [8, 18] we proved the following existence principle.

Theorem 4.4. *Let $pqf : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ be a L^1 -Carathéodory function with (2.1), (2.2), (4.19) and (4.20) holding. In addition suppose there is a constant M_0 , independent of λ , with*

$$|y|_0 = \sup_{[0, 1]} |y(t)| \leq M_0$$

for any solution y (here $y \in C[0, 1] \cap C^1(0, 1)$ with $py' \in AC[0, 1]$) to

$$\begin{cases} \frac{1}{pq}(py')' = \lambda f(t, y) \text{ a.e. on } [0, 1] \\ \lim_{t \rightarrow 0^+} p(t)y'(t) = 0 \\ y(1) = 0 \end{cases} \tag{4.21}_\lambda$$

for each $0 < \lambda < 1$. Then (4.9) has at least one solution.

Theorem 4.5. *Let $pqf : [0, 1] \times \mathbf{R} \rightarrow \mathbf{R}$ be a L^1 -Carathéodory function with (2.1), (2.2) and (4.19) holding. In addition suppose*

$$\begin{cases} |f(t, u)| \leq \phi_1(t) + \phi_2(t)\Omega(|u|) \text{ a.e. on } [0, 1] \text{ where } \Omega : [0, \infty) \rightarrow [0, \infty) \\ \text{is a nondecreasing continuous function} \end{cases} \tag{4.22}$$

$$\int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_i(x) dx ds < \infty, i = 1, 2 \tag{4.23}$$

and

$$\left(\int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_2(x) dx ds \right) \limsup_{x \rightarrow \infty} \frac{\Omega(x)}{x} < 1 \tag{4.24}$$

are satisfied. Then (4.9) has at least one solution.

Proof. Let y be a solution to (4.21) $_{\lambda}$ for $0 < \lambda < 1$. Then for $t \in [0, 1]$ we have

$$y(t) = - \int_t^1 \frac{1}{p(s)} \int_0^s p(x)q(x)f(x, y(x)) dx ds$$

and so

$$|y(x)| \leq \int_t^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_1(x) dx ds + \int_t^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_2(x)\Omega(|y(x)|) dx ds.$$

Now $|y(x)| \leq \sup_{[0, 1]} |y(s)| \equiv |y|_0$ and this together with the fact that Ω is nondecreasing yields

$$|y(t)| \leq \int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_1(x) dx ds + \Omega(|y|_0) \int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_2(x) dx ds.$$

Let $K_i = \int_0^1 \frac{1}{p(s)} \int_0^s p(x)q(x)\phi_i(x) dx ds, i = 1, 2$ so

$$|y|_0 \leq K_1 + K_2\Omega(|y|_0)$$

and consequently

$$1 \leq \frac{K_1}{|y|_0} + \frac{K_2\Omega(|y|_0)}{|y|_0}.$$

Thus (as in Theorem 4.1) there exists a constant M_0 , independent of λ , with $|y|_0 \leq M_0$ for any solution y to (4.21) $_{\lambda}$. The result follows from Theorem 4.4. \square

REFERENCES

1. F. V. ATKINSON, *Discrete and continuous boundary problems* (Academic Press, New York, 1964).
2. L. E. BOBISUD and D. O'REGAN, Positive solutions for a class of nonlinear singular boundary value problems at resonance, *J. Math. Anal. Appl.* **184** (1994), 263–284.
3. P. L. CHAMBRE, On the solution of the Poisson–Boltzmann equation with application to the theory of thermal explosions, *J. Chem. Phys.* **20** (1952), 1795–1797.
4. M. M. CHAWLA and P. N. SHIVAKUMAR, On the existence of solutions of a class of singular nonlinear two point boundary value problems, *J. Comput. Appl. Math.* **19** (1987), 379–388.
5. D. R. DUNNINGER and J. C. KURTZ, A priori bounds and existence of positive solutions for singular nonlinear boundary value problems, *SIAM J. Math. Anal.* **17** (1986), 595–609.
6. M. A. EL-GEBEILY, A. BOUMENIR and A. B. M. ELGINDI, Existence and uniqueness of solutions

of a class of two-point singular nonlinear boundary value problems, *J. Comput. Appl. Math.*, **46** (1993), 345–355.

7. A. FONDA and J. MAWHIN, Quadratic forms, weighted eigenfunctions and boundary value problems for nonlinear second order differential equations, *Proc. Royal Soc. Edinburgh* **112A** (1989), 145–153.

8. M. FRIGON and D. O'REGAN, Some general existence principles for ordinary differential equations, *Topological Methods in Nonlinear Anal.* **2** (1993), 35–54.

9. A. GRANAS, R. B. GUENTHER and J. W. LEE, Some general existence principles in the Carathéodory theory of nonlinear differential systems, *J. Math. Pures Appl.* **70** (1991), 153–196.

10. R. IANNACCI and M. N. NKASHAMA, Nonlinear two-point boundary value problems at resonance without Landesman–Lazer conditions, *Proc. Amer. Math. Soc.* **106** (1989), 943–952.

11. J. B. KELLER, Electrodynamics I. The equilibrium of a charged gas in a container, *J. Rat. Mech. Anal.* **5** (1957), 715–724.

12. M. A. KRASNOSELSKII, *Topological methods in the theory of nonlinear integral equations* (MacMillan Co., New York, 1964).

13. J. MAWHIN, J. R. WARD and M. WILLEM, Necessary and sufficient conditions for the solvability of a nonlinear two point boundary value problem, *Proc. Amer. Math. Soc.* **93** (1985), 667–674.

14. J. MAWHIN and W. OMANO, Two point boundary value problems for nonlinear perturbations of some singular linear differential equations at resonance, *Comment. Math. Univ. Carolinae* **30** (1989), 537–550.

15. J. W. MOONEY, Numerical schemes for degenerate boundary value problems, *J. Phys. A.* **26** (1993), 413–421.

16. M. A. NAIMARK, *Linear differential operators* Part II (Ungar Publ. Co., London, 1968).

17. D. O'REGAN, Solvability of some two point boundary value problems of Dirichlet, Neumann or Periodic type, *Dynamic Systems and Appl.* **2** (1993), 163–182.

18. D. O'REGAN, Singular Sturm Liouville problems and existence of solutions to singular nonlinear boundary value problems, *Nonlinear Anal.* **20** (1993), 767–779.

19. D. O'REGAN, Existence theory for singular two point boundary value problems, in *Proc. Fourth Int. Coll. Diff. Eq.* (Plovdiv, Bulgaria, *Int. Science Publ.*, Utrecht, 1994), 215–228.

20. D. O'REGAN, Existence principles for second order nonresonant boundary value problems, *J. Appl. Math. Stoch. Anal.* **7** (1994), 487–507.

21. D. O'REGAN, Nonresonant and resonant singular boundary value problems, to appear.

22. D. POWERS, *Boundary value problems* (Harcourt Brace Jovanovich, San Diego, 1987).

23. L. SANCHEZ, Positive solutions for a class of semilinear two point boundary value problems, *Bull. Austral. Math. Soc.* **45** (1992), 439–451.

24. I. STAKGOLD, *Greens functions and boundary value problems* (John Wiley and Sons, New York, 1979).

DEPARTMENT OF MATHEMATICS
UNIVERSITY COLLEGE GALWAY
GALWAY
IRELAND